Institute for Computational Mathematics



Model Order Reduction for Maxwell Equations based on Moment Matching

Matthias Bollhöfer (joint work with André Bodendiek) INRIA Sophia Antipolis , July 28, 2015

Outline

Maxwell's Equations

Model Order Reduction

- MOR for Maxwell Equations
- Model Order Reduction Based on Moment Matching
- Numerical Results

Modified Adaptive-Order Rational Arnoldi Method

- Savings for the AORA method
- QMR Method
- Simplified QMR
- QMR with Subspace Recycling
- Numerical Results
- Conclusions



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Model Problems



Left picture: branchline coupler on a substrate with PMC boundary conditions, two parallel microstriplines, coupled together in form of a transversal bridge, frequency range 1.0 to 10.0 GHz, N = 73'385.

Right picture: coplanar waveguide with a dielectric overlay, PEC boundary conditions, frequency range 0.6 to 3.0 GHz, N = 32'924.



Model Problems

Braunschweig



Maxwell's Equations

$$\frac{\partial(\varepsilon \mathbf{E})}{\partial t} = -\sigma \mathbf{E} + \nabla \times \mathbf{H}$$

$$\frac{\partial(\mu \mathbf{H})}{\partial t} = -\nabla \times \mathbf{E}$$

$$(\mathbf{0} = \nabla \cdot (\varepsilon \mathbf{E}), \ \mathbf{0} = \nabla \cdot (\mu \mathbf{H}))$$

 ϵ electric permittivity, μ magnetic permeability, σ electric conductivity.



Maxwell's Equations

$$\begin{aligned} \frac{\partial(\varepsilon \mathbf{E})}{\partial t} &= -\sigma \mathbf{E} + \nabla \times \mathbf{H} \\ \frac{\partial(\mu \mathbf{H})}{\partial t} &= -\nabla \times \mathbf{E} \\ (0 &= \nabla \cdot (\varepsilon \mathbf{E}), \ 0 = \nabla \cdot (\mu \mathbf{H}) \end{aligned}$$

 ϵ electric permittivity, μ magnetic permeability, σ electric conductivity. Discrete equations:

$$M_{\varepsilon}\dot{E} = -M_{\sigma}E + GH + B_{E}u + b.c.$$

$$M_{\mu}\dot{H} = -G^{T}E + B_{H}u + b.c.$$

$$(0 = D_{E}M_{\varepsilon}E, \quad 0 = D_{H}M_{\mu}H) + C_{E}E + C_{H}H$$

- *u* input, *y* output
- M_{ε} , M_{μ} are sym. pos. def., M_{σ} sym. pos. semidef. (mass matrices)
- G highly singular! (curl operator)



Technische Universität Braunschweig

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Model Order Reduction

$$\begin{aligned} \mathcal{M}\dot{x} &= \mathcal{A}x + \mathcal{B}u \\ (0 &= \mathcal{D}x) \\ y &= \mathcal{C}x \end{aligned}$$

$$\mathfrak{M} = \left(\begin{array}{cc} M_{\varepsilon} & 0\\ 0 & M_{\mu} \end{array}\right), \ \mathcal{A} = \left(\begin{array}{cc} -M_{\sigma} & G\\ -G^{T} & 0 \end{array}\right), \ \mathbf{x} = \left(\begin{array}{cc} E\\ H \end{array}\right).$$



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Model Order Reduction

$$\begin{aligned} \mathcal{M}\dot{x} &= \mathcal{A}x + \mathcal{B}u \\ (0 &= \mathcal{D}x) \\ y &= \mathcal{C}x \end{aligned}$$
$$\mathcal{M} = \begin{pmatrix} M_{\varepsilon} & 0 \\ 0 & M_{\mu} \end{pmatrix}, \ \mathcal{A} = \begin{pmatrix} -M_{\sigma} & G \\ -G^{T} & 0 \end{pmatrix}, \ x = \begin{pmatrix} E \\ H \end{pmatrix}. \end{aligned}$$

Model Order Reduction: find full rank $S, T \in \mathbb{R}^{2n,r}$ such that $r \ll 2n$ and use instead

$$\begin{array}{rcl} \left(\boldsymbol{S}^{*}\boldsymbol{\mathcal{M}}\boldsymbol{T}\right) \dot{\boldsymbol{\tilde{x}}} &=& \left(\boldsymbol{S}^{*}\boldsymbol{\mathcal{A}}\boldsymbol{T}\right) \boldsymbol{\tilde{x}} + \left(\boldsymbol{S}^{*}\boldsymbol{\mathcal{B}}\right) \boldsymbol{u} \\ \left(\boldsymbol{0} &=& \left(\boldsymbol{\mathcal{D}}\boldsymbol{T}\right) \boldsymbol{\tilde{x}} \right) \\ \boldsymbol{\tilde{y}} &=& \left(\boldsymbol{\mathcal{C}}\boldsymbol{T}\right) \boldsymbol{\tilde{x}} \end{array}$$

$$\|y - \tilde{y}\|$$
 small



Model Order Reduction

$$\begin{aligned} \mathcal{M}\dot{x} &= \mathcal{A}x + \mathcal{B}u \\ (0 &= \mathcal{D}x) \\ y &= \mathcal{C}x \end{aligned}$$
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Model Order Reduction: find full rank $S, T \in \mathbb{R}^{2n,r}$ such that $r \ll 2n$ and use instead

$$\hat{\mathcal{M}}\dot{\tilde{x}} = \hat{\mathcal{A}}\tilde{x} + \hat{\mathcal{B}}u
(0 = \hat{\mathcal{D}}\tilde{x})
\tilde{y} = \hat{\mathbb{C}}\tilde{x}
\|y - \tilde{y}\| \text{ small}$$



Structure-Preserving MOR for Maxwell's Equations

Here use
$$S = T = \begin{pmatrix} V & 0 \\ 0 & W \end{pmatrix}$$

Structured MOR: find full rank V, $W \in \mathbb{R}^{n,r}$ such that $r \ll n$ and use instead

$$(V^* M_{\varepsilon} V) \dot{e} = -(V^* M_{\sigma} V) e + (V^* G W) h + (V^* B_E) u (W^* M_{\mu} W) \dot{h} = -(W^* G^T V) e + (W^* B_H) u (0 = (D_E M_{\varepsilon} V) e, 0 = (D_H M_{\mu} W) h) \tilde{y} = (C_E V) e + (C_H W) h$$

 $\|y - \tilde{y}\|$ small



Structure-Preserving MOR for Maxwell's Equations

Here use
$$S = T = \begin{pmatrix} V & 0 \\ 0 & W \end{pmatrix}$$

Structured MOR: find full rank $V, W \in \mathbb{R}^{n,r}$ such that $r \ll n$ and use instead

$$\begin{split} \tilde{M}_{\varepsilon} \dot{e} &= -\tilde{M}_{\sigma} e + \tilde{G} h + \tilde{B}_{E} u \\ \tilde{M}_{\mu} \dot{h} &= -\tilde{G}^{T} e + \tilde{B}_{H} u \\ (0 &= \tilde{D}_{E} e, 0 = \tilde{D}_{H} h) \\ \tilde{y} &= \tilde{C}_{E} e + \tilde{C}_{H} h \end{split}$$

 $\|y - \tilde{y}\|$ small



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Moment-Matching — Basic Idea

$$\mathcal{M} = \begin{pmatrix} M_{\varepsilon} & 0\\ 0 & M_{\mu} \end{pmatrix}, \ \mathcal{A} = \begin{pmatrix} -M_{\sigma} & G\\ -G^{T} & 0 \end{pmatrix}, \ \mathcal{B} = \begin{pmatrix} B_{E}\\ B_{H} \end{pmatrix}, \ \mathcal{C} = \begin{pmatrix} C_{E} & C_{H} \end{pmatrix}.$$



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 $\mathfrak{H}(\boldsymbol{s}) = \mathfrak{CP} \left(\boldsymbol{s} \mathfrak{M} - \mathcal{A} \right)^{-1} \mathfrak{B} \qquad (\mathfrak{P} \text{ projector to divergence-free part })$



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Taylor/Laurent expansion at some expansion point *s_j*:

$$\mathcal{A}_j := (\mathbf{s}_j \mathcal{M} - \mathcal{A})^{-1} \mathcal{M}, \ \mathcal{B}_j := (\mathbf{s}_j \mathcal{M} - \mathcal{A})^{-1} \mathcal{B}, \ \mathcal{C}_j := \mathcal{CP} (\mathbf{s}_j \mathcal{M} - \mathcal{A})^{-1}.$$



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 $\mathcal{H}(s) = \mathbb{CP} (s\mathcal{M} - \mathcal{A})^{-1} \mathcal{B}$ (\mathcal{P} projector to divergence-free part) Taylor/Laurent expansion at some expansion point s_i :

 $\mathcal{A}_j := (\mathbf{s}_j \mathcal{M} - \mathcal{A})^{-1} \mathcal{M}, \ \mathcal{B}_j := (\mathbf{s}_j \mathcal{M} - \mathcal{A})^{-1} \mathcal{B}, \ \mathcal{C}_j := \mathcal{CP} (\mathbf{s}_j \mathcal{M} - \mathcal{A})^{-1}.$

$$\Rightarrow \mathcal{H}(\boldsymbol{s}) = \sum_{\rho=0}^{\infty} \mathcal{CP}(\overline{-\mathcal{A}_j})^{\rho} \mathcal{B}_j(\boldsymbol{s}-\boldsymbol{s}_j)^{\rho} = \sum_{\rho=0}^{\infty} \underbrace{\mathcal{C}_j[-\mathcal{A}_j]^{\rho}}_{\boldsymbol{Y}_j^{(\rho)}} \mathcal{B}(\boldsymbol{s}-\boldsymbol{s}_j)^{\rho}$$



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 $X_j^{(p)}$ input moments, Taylor coefficients $Z_j^{(p)} = C \mathcal{P} X_j^{(p)} = Y_j^{(p)} \mathcal{B}$ output moments.



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Krylov subspace

$$\mathcal{K}_{p}(A, b) = \operatorname{span}\{b, Ab, \dots, A^{p-1}b\}$$



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input Krylov subpace

$$X_j^{(p)} \in \mathcal{K}_p(\mathcal{A}_j, \mathcal{B}_j).$$

output Krylov subspace

$$(Y_j^{(p)})^* \in \mathcal{K}_p(\mathcal{A}_j^*, \mathcal{C}_j^*).$$



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$$(\mathbf{Y}_{j}^{(p)})^{*} \in \mathcal{K}_{p}(\mathcal{A}_{j}^{*}, \mathcal{C}_{j}^{*}).$$

• Lanczos-type methods [PVL,Gragg'74,Gutknecht'92,Feldmann,Freund'94,...] generate dual basis $T \equiv T_r \in \mathbb{R}^{2n,r}$ and $S \equiv S_r \in \mathbb{R}^{2n,r}$ of

$$\mathcal{K}_r(\mathcal{A}_j, \mathcal{B}_j)$$
 and $\mathcal{K}_r(\mathcal{A}_i^*, \mathcal{C}_i^*)$

such that $S^*T = I \rightarrow$ matches 2*r* moments $Z_j^{(0)}, \ldots, Z_j^{(2r-1)}$



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such that $S^*T = I \rightarrow$ matches 2*r* moments $Z_j^{(0)}, \ldots, Z_j^{(2r-1)}$

• Arnoldi-type methods [PRIMA,Odabasioglu et al.'96,'97], [SPRIM,Freund'04,'08] compute one orthonormal basis $S = T = Q \equiv Q_r \in \mathbb{R}^{2n,r}$, say, from

$$\mathcal{K}_r(\mathcal{A}_j, \mathcal{B}_j)$$

using modified Gram-Schmidt \rightarrow matches *r* moments $Z_j^{(0)}, \ldots, Z_j^{(r-1)}$



$$\mathfrak{M} = \left(\begin{array}{cc} M_{\varepsilon} & 0 \\ 0 & M_{\mu} \end{array} \right), \ \mathcal{A} = \left(\begin{array}{cc} -M_{\sigma} & G \\ -G^{T} & 0 \end{array} \right).$$

$$\begin{aligned} \mathcal{M}\dot{x} &= \mathcal{A}x + \mathcal{B}u \\ y &= \mathcal{C}x \end{aligned}$$



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$$(S^* \mathcal{M}T) \dot{\tilde{x}} = (S^* \mathcal{A}T) \tilde{x} + (S^* \mathcal{B}) u$$
$$\tilde{y} = (CT) \tilde{x}$$



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$$\hat{\mathcal{M}}\dot{\tilde{x}} = \hat{\mathcal{A}}\tilde{x} + \hat{\mathcal{B}}u \\ \tilde{y} = \hat{\mathbb{C}}\tilde{x}$$



Structure Preservation for Maxwell Equations

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• Lanczos-type methods: NO! $S \neq T$!



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- Lanczos-type methods: NO! $S \neq T$!
- Arnoldi-type methods: S = T = Q, but block structure is lost



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$$Q = \left[\begin{array}{c} V \\ W \end{array} \right] \rightarrow \left[\begin{array}{c} V & 0 \\ 0 & W \end{array} \right]$$



$$\mathfrak{M} = \left(\begin{array}{cc} M_{\varepsilon} & 0 \\ 0 & M_{\mu} \end{array} \right), \ \mathcal{A} = \left(\begin{array}{cc} -M_{\sigma} & G \\ -G^{T} & 0 \end{array} \right).$$

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$$Q = \left[\begin{array}{c} V \\ W \end{array} \right] \rightarrow \left[\begin{array}{c} V & 0 \\ 0 & W \end{array} \right]$$

- twice as big, but ...
- block-structure preserved, still r moments matched
- if we are lucky, up to 2r moments could be matched



Problems

- (How to) select $s_j \in [f_{\min}, f_{\max}]$
- (No) error bounds!? Choice of r, l, accuracy of the reduced model

$$\mathcal{H}(\boldsymbol{s}) = \mathcal{C} \left(\boldsymbol{s} \mathcal{M} - \mathcal{A} \right)^{-1} \mathcal{B}, \quad \mathcal{H}_r(\boldsymbol{s}) = \hat{\mathcal{C}} \left(\boldsymbol{s} \hat{\mathcal{M}} - \hat{\mathcal{A}} \right)^{-1} \hat{\mathcal{B}}$$
$$\|\mathcal{H}(i\boldsymbol{\omega}) - \mathcal{H}_r(i\boldsymbol{\omega})\| \leqslant \dots$$

- multiple expansion points $s_1, \ldots s_l \in [f_{\min}, f_{\max}]$
- Restarting Arnoldi and increasing r or / whenever the "error estimate" is not accurate enough



- multiple expansion points *s*₁,..., *s*_l
- multiple associated Taylor expansions

$$\mathfrak{H}(s) = \sum_{p=0}^{\infty} Z_j^{(p)} (s - s_j)^p, \ j = 1, \dots, I$$

• Rational Krylov method: Compute basis Q_r for the Krylov subspaces

$$\sum_{j=1}^{l} \mathcal{K}_{r_j}(\mathcal{A}_j, \mathcal{B}_j).$$



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Lemma (Key-Lemma, Grimme, Gallivan'98)

If $s_j \neq s_k$ then

$$\mathcal{A}_k \cdot \mathcal{A}_j^{p-1} \mathcal{B}_j \in \mathcal{K}_p(\mathcal{A}_j, \mathcal{B}_j) + \mathcal{K}_1(\mathcal{A}_k, \mathcal{B}_k)$$



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- multiple associated Taylor expansions

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 \Rightarrow mixing inverses with different shifts leads to a separate sum of Krylov subspaces, no "mixed powers of inverses"



 Traditional Rational Arnoldi methods build Q_r w.r.t. K_{rj}(A_j, B_j), j = 1, ..., l <u>one after another</u> using modified Gram-Schmidt.


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 $\mathfrak{K}_{\mathbf{r_1}}(\mathcal{A}_1, \mathfrak{B}_1)$



 Traditional Rational Arnoldi methods build Q_r w.r.t. K_{rj}(A_j, B_j), j = 1, ..., l <u>one after another</u> using modified Gram-Schmidt.

 $\mathfrak{K}_{r_1}(\mathcal{A}_1, \mathfrak{B}_1) + \mathfrak{K}_{r_2}(\mathcal{A}_2, \mathfrak{B}_2)$



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 $\mathcal{K}_{r_1}(\mathcal{A}_1, \mathcal{B}_1) + \mathcal{K}_{r_2}(\mathcal{A}_2, \mathcal{B}_2) + \mathcal{K}_{r_3}(\mathcal{A}_3, \mathcal{B}_3)$



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 $\mathfrak{K}_{r_1}(\mathcal{A}_1, \mathcal{B}_1) + \mathfrak{K}_{r_2}(\mathcal{A}_2, \mathcal{B}_2) + \mathfrak{K}_{r_3}(\mathcal{A}_3, \mathcal{B}_3) + \mathfrak{K}_{r_4}(\mathcal{A}_4, \mathcal{B}_4)$



 Traditional Rational Arnoldi methods build Q_r w.r.t. K_{rj}(A_j, B_j), j = 1, ..., l <u>one after another</u> using modified Gram-Schmidt.

$$\operatorname{span} \mathbf{Q}_{r} = \mathcal{K}_{r_{1}}(\mathcal{A}_{1}, \mathcal{B}_{1}) + \mathcal{K}_{r_{2}}(\mathcal{A}_{2}, \mathcal{B}_{2}) + \cdots + \mathcal{K}_{r_{l}}(\mathcal{A}_{l}, \mathcal{B}_{l})$$



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 [Ruhe'94], [Gallivan,Grimme,van Dooren'95], [Grimme'99], [Bai'02], [Gugercin,Antoulas'06], [Lee,Chu,Feng'06],... and many others



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- Adaptive-Order Rational Arnoldi (AORA) [Lee,Chu,Feng'06]
 Q_r is generated by **interchangeably** increasing the size *r_j* of *K_{r_i}(A_j, B_j)*.



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 Traditional Rational Arnoldi methods build Q_r w.r.t. K_{rj}(A_j, B_j), j = 1, ..., l one after another using modified Gram-Schmidt.

 $\operatorname{span} \boldsymbol{Q}_{r} = \mathcal{K}_{r_{1}}(\mathcal{A}_{1}, \mathcal{B}_{1}) + \mathcal{K}_{r_{2}}(\mathcal{A}_{2}, \mathcal{B}_{2}) + \cdots + \mathcal{K}_{r_{l}}(\mathcal{A}_{l}, \mathcal{B}_{l})$

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• At each step s_j is selected w.r.t. the largest output moment error $Z_j^{(r)} - \tilde{Z}_j^{(r)}$ which can be computed cheaply



• AORA called repeatedly *I* = 1, 2, 3, ... with increasing number of expansion points



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$$s_1 = 2\pi i f_{\min}, \ s_2 = 2\pi i \sqrt{f_{\min} f_{\max}}, \ s_3 = 2\pi i f_{\max}$$



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 S_1, S_2, S_3, S_4



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 S_1, S_2, S_3, S_4, S_5



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 $S_1, S_2, S_3, S_4, S_5, S_6$



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• relative error $\frac{|\mathcal{H}_{\ell}^{(I)}(s) - \mathcal{H}_{\ell}^{(I-1)}(s)|}{|\mathcal{H}_{\ell}^{(I)}(s)|}$ between two computed reduced-order transfer functions used as measure [Köhler et al.'10'12]



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Maxwell's Equations

Model Order Reduction

- MOR for Maxwell Equations
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Model Order Reduction — Numerical Results

- model problems have a frequency range in [f_{min}, f_{max}]
- computed reduced order models have fixed size n = 25(50)
- SPRIM uses expansion point $s_0 = \frac{f_{min} + f_{max}}{2}$
- expansion point selection based on relative error 10⁻⁹
- strategy finally leads to *I* = 8 expansion points *s*₁,..., *s*₈
- Rational Arnoldi (RA) and Adaptive Order Rational Arnoldi (AORA) repeated 5 times
- RA uses fixed sizes j = n/l for each Krylov subspace \mathcal{K}_j
- AORA adaptively increases each *K_{ji}*



Model Order Reduction — Numerical Results

Branchline Coupler



- size N = 73385, discretized using FIT
- frequency range $[f_{min}, f_{max}] = [10^9, 10^{10}]$



Model Order Reduction — Numerical Results Branchline Coupler

6

relative error

$$arepsilon_{rel}(f) = rac{|\mathcal{H}(oldsymbol{s}) - ilde{\mathcal{H}}(oldsymbol{s})|}{|\mathcal{H}(oldsymbol{s})|},$$

 $s = 2\pi i f$ with $f \in [10^9, 10^{10}]$ displayed as reference





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Model Order Reduction — Numerical Results

Coplanar Waveguide



- size N = 32924, discretized using FIT
- frequency range $[f_{min}, f_{max}] = [0.6 \cdot 10^9, 3.0 \cdot 10^9]$



Model Order Reduction — Numerical Results

Coplanar Waveguide

relative error

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Modified Adaptive-Order Rational Arnoldi Method

Repeatedly calling AORA method requires to recompute $span Q_r$ from scratch.



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Modified Adaptive-Order Rational Arnoldi Method

Repeatedly calling AORA method requires to recompute $\operatorname{span} Q_r$ from scratch.

Lemma

Suppose that

$$\operatorname{span} Q_{r}^{(l)} = \mathcal{K}_{r_{1}^{(l)}}(s_{1}) + \dots + \mathcal{K}_{r_{l}^{(l)}}(s_{l}),$$
$$\operatorname{span} Q_{r}^{(l+1)} = \mathcal{K}_{r_{1}^{(l+1)}}(s_{1}) + \dots + \mathcal{K}_{r_{l}^{(l+1)}}(s_{l}) + \mathcal{K}_{r_{l+1}^{(l+1)}}(s_{l+1})$$

such that

$$r_1^{(l+1)} \leqslant r_1^{(l)} \cdots r_l^{(l+1)} \leqslant r_l^{(l)}.$$

Then $\mathcal{K}_{r_{1}^{(l+1)}}(s_{1}) + \cdots + \mathcal{K}_{r_{l}^{(l+1)}}(s_{l})$ can be directly extracted from $\operatorname{span} Q_{r}^{(l)}$.


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- Lemma requires that shifts s_1, \ldots, s_l are selected in the same order
- new shift s_{l+1} can be injected at any time
- only new shift s_{l+1} requires to solve systems with $s_{l+1}\mathcal{M}-\mathcal{A}$
- $\blacksquare \longrightarrow \mathsf{mAORA} \text{ (modified AORA)}$



Solving Systems in the Modified AORA Method

Recall

- $\mathcal{J}(s\mathcal{M} \mathcal{A})$ is complex-symmetric, where $\mathcal{J} = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}$
- Schur complement $S(s) = s^2 M_{\epsilon} + s M_{\sigma} + G M_{\mu}^{-1} G^T$ is complex-symmetric



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$$A(s_j)x_{jp} = b_{jp}$$



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$$A(s_j)x_{jp} = b_{jp}$$

- Save memory and time: Only compute factorization of $A(s_*)$ for some characteristic shift $s_* = 2\pi i \sqrt{f_{\min} f_{\max}}$.
- For all $s_j \neq s_*$ use (recycling) Krylov subspace method as wrapper.



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QMR Method — Sketch

Two-sided Lanczos Method

$$AV_k = V_{k+1} \underline{T_k}, \ A^T \widetilde{V}_k = \widetilde{V}_{k+1} \underline{\widetilde{T}_k}$$



Two-sided Lanczos Method

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QMR for Ax = b based on two-sided Lanczos method:

- $r_0 = b Ax_0, v_1 = r_0 / ||r_0||$
- $x_k = x_0 + V_k y$
- quasi-minimize

$$\|\boldsymbol{b} - \boldsymbol{A}\boldsymbol{x}_k\| \leqslant \|\boldsymbol{V}_{k+1}\| \cdot \|\|\boldsymbol{r}_0\|\boldsymbol{e}_1 - \underline{T}_k\boldsymbol{y}\|$$

- $\longrightarrow y$ from least squares solution $\longrightarrow x$
- [Freund et al. '91,'93,'94]



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- [Freund et al. '94,'95]



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- [Freund et al. '94,'95]
- Maxwell equations: complex-symmetric system A with complex-symmetric preconditioner P satisfy J-symmetry!



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QMR with Subspace Recycling

Krylov subspace methods with subspace recycling

- GCROT, GCRO-DR [De Sturler et al.'99,'06]
- recycling BiCG, recycling BiCGStab [Ahuja et al. '12,'13]

Main idea: $U, \tilde{U} \in \mathbb{C}^{n,r}$, $C = AU, \tilde{C} = A^T \tilde{U}, r \ll n$ given subspaces, $D_c = \tilde{C}^T C$. Apply Lanczos (Arnoldi) to the systems

$$AV_k = V_{k+1} \underline{T_k}$$

$$A^{ au} ilde{V}_k = ilde{V}_{k+1} ilde{T}_k$$

with updated initial guess and projected initial residual

$$x_0^{(new)} = x_0 + U D_c^{-1} \tilde{C}^T r_0, \ r_0^{(new)} = (I - C D_c^{-1} \tilde{C}^T) r_0.$$



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$$(I - CD_c^{-1}\tilde{C}^T)AV_k = V_{k+1}\underline{T_k},$$

$$(I - \tilde{C} D_c^{-T} C^T) A^T \tilde{V}_k = \tilde{V}_{k+1} \underline{\tilde{T}_k}$$

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$$U, \tilde{U}, C = AU, \tilde{C} = A^T \tilde{U}, D_c = \tilde{C}^T C.$$

Changing the methods:

BiCG/QMR	Recycling BiCG/QMR
init x_0 init $r_0 = b - Ax_0$	$ \begin{aligned} \boldsymbol{x}_0^{(new)} &= \boldsymbol{x}_0 + \boldsymbol{U} \boldsymbol{D}_c^{-1} \tilde{\boldsymbol{C}}^{T} \boldsymbol{r}_0 \\ \boldsymbol{r}_0^{(new)} &= (\boldsymbol{I} - \boldsymbol{C} \boldsymbol{D}_c^{-1} \tilde{\boldsymbol{C}}^{T}) \boldsymbol{r}_0 \end{aligned} $
matvec Ax matvec $A^T x$	$(I - CD_c^{-1}\tilde{C}^T)Ax (I - \tilde{C}D_c^{-T}C^T)A^Tx$
solution update $x = x + \alpha p$	$x = x + \alpha (I - UD_c^{-1} \tilde{C}^T) \rho$



$$U, \tilde{U}, C = AU, \tilde{C} = A^T \tilde{U}, D_c = \tilde{C}^T C.$$

Changing the methods:

BiCG/QMR	Recycling BiCG/QMR
init x_0 init $r_0 = b - Ax_0$	$ \begin{aligned} x_0^{(new)} &= x_0 + U D_c^{-1} \tilde{C}^T r_0 \\ r_0^{(new)} &= (I - C D_c^{-1} \tilde{C}^T) r_0 \end{aligned} $
matvec Ax matvec $A^T x$	$(I - CD_c^{-1}\tilde{C}^T)Ax$ $(I - \tilde{C}D_c^{-T}C^T)A^Tx$
solution update $x = x + \alpha p$	$x = x + \alpha (I - UD_c^{-1} \tilde{C}^T)p$

Preconditioning and simplified QMR: a little bit more tricky but similar



Outline

Maxwell's Equations

- Model Order Reduction
 - MOR for Maxwell Equations
 - Model Order Reduction Based on Moment Matching
 - Numerical Results
- Modified Adaptive-Order Rational Arnoldi Method
 - Savings for the AORA method
 - QMR Method
 - Simplified QMR
 - QMR with Subspace Recycling
- Numerical Results
- Conclusions



Modified AORA versus AORA





Model Order Reduction

Numerical Results

Dielectric overlay



Sampling 20 frequencies $s_i = 2\pi i f_i$ $f_i \in [f_{\min}, f_{\max}] = [0.6 \cdot 10^9, 3 \cdot 10^9]$ Use complex-symmetric Schur complement system Preconditioner: LU-decomposition for $s_* = \sqrt{f_{\min} f_{\max}}$

Comparison

- simplified QMR without subspace recycling
- simplified QMR using subspace recycling



Model Order Reduction

Numerical Results

Preconditioner: LU decomposition for $S(s_*)$





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Model Order Reduction

Numerical Results

Preconditioner: LU decomposition for $S(s_*)$





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Conclusions

- (adaptive order) rational Arnoldi yields reduced order model for Maxwell equations
- main important structures can be preserved
- outer iteration: several calls of (modified) rational Arnoldi, inner iteration: preconditioned SQMR
- Recycling techniques \rightarrow modified AORA, recycling SQMR

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