

# PARALLÉLISME EN TEMPS: Résultats de Convergence

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## Convergence for the Heat Equation

### Corollary

*The parareal algorithm applied to the heat equation  $u_t = \Delta u$  discretized with an  $L$ -stable method in time converges superlinearly on bounded time intervals,*

$$\max_{1 \leq n \leq N} \|u(t_n) - U_n^k\|_2 \leq \frac{\gamma_s^k}{k!} \prod_{j=1}^k (N - j) \max_{1 \leq n \leq N} \|u(t_n) - U_n^0\|_2,$$

*where the constant  $\gamma_s < 1$  is universal for each  $L$ -stable method. On unbounded time intervals the convergence is linear,*

$$\sup_{n>0} \|u(t_n) - U_n^k\|_2 \leq \gamma_l^k \sup_{n>0} \|u(t_n) - U_n^0\|_2,$$

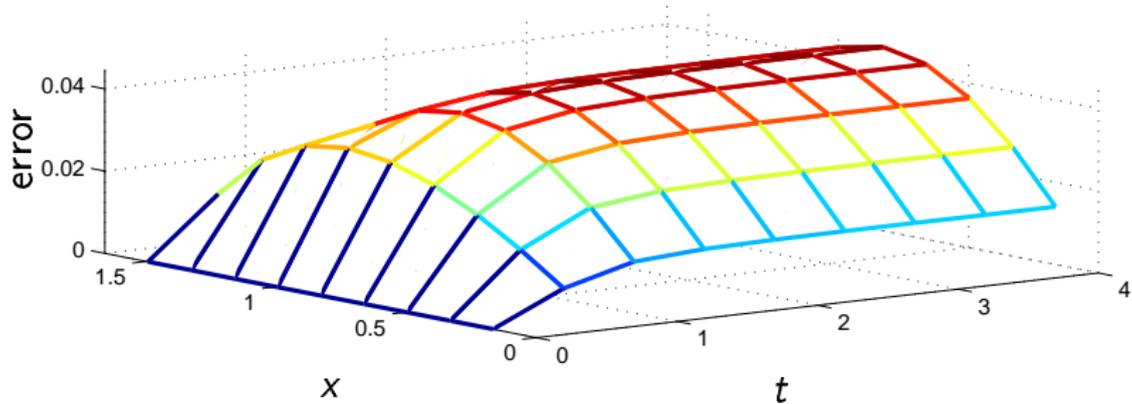
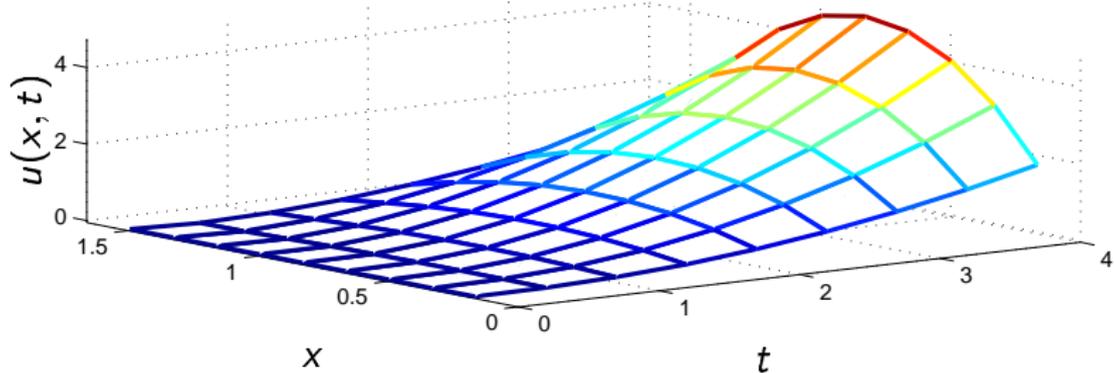
*where  $\gamma_l < 1$  is universal for each  $L$ -stable method.*

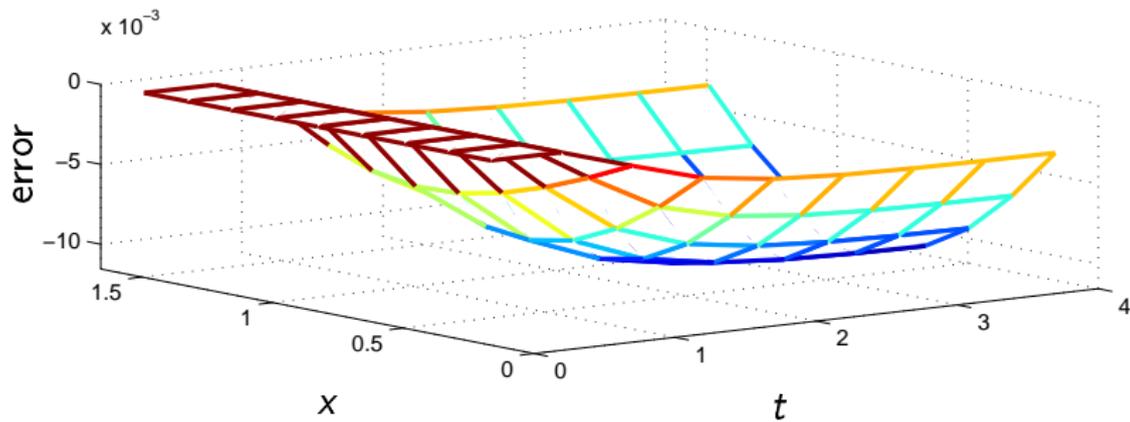
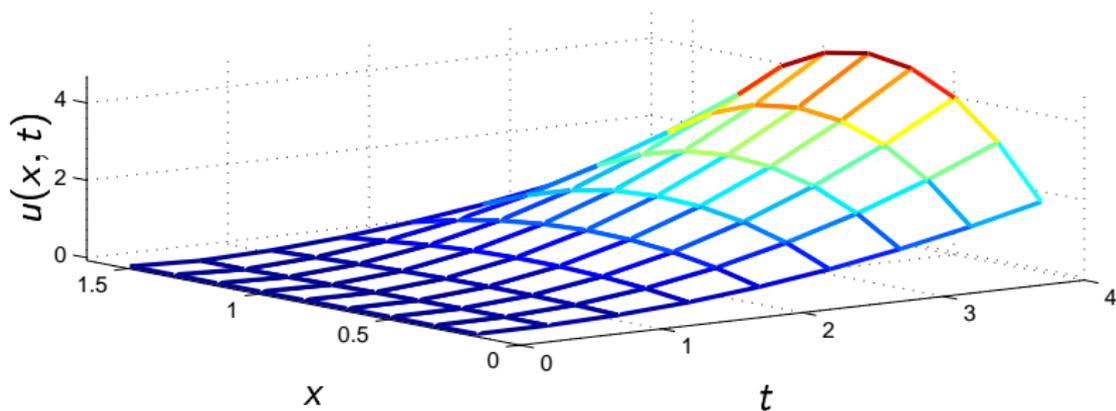
## Convergence Constants for the Heat Equation

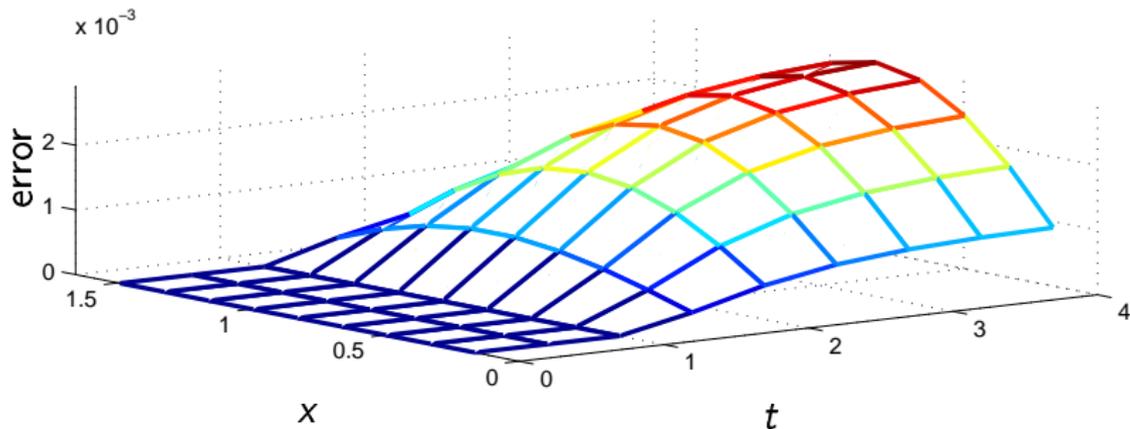
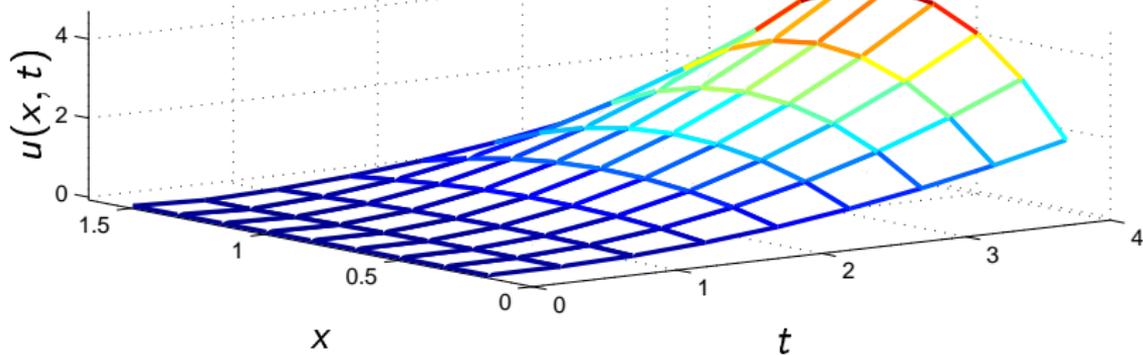
The convergence constants  $\gamma_s$  and  $\gamma_I$  can be computed for each L-stable method:

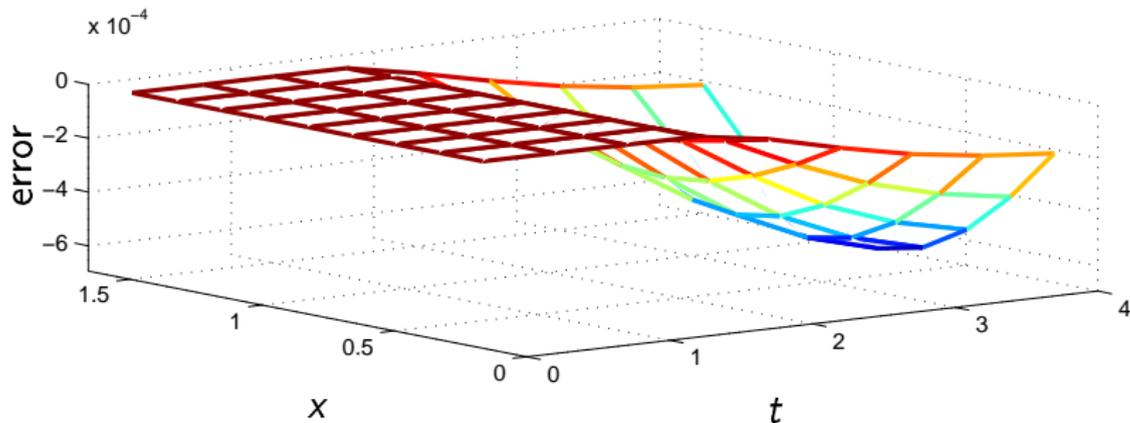
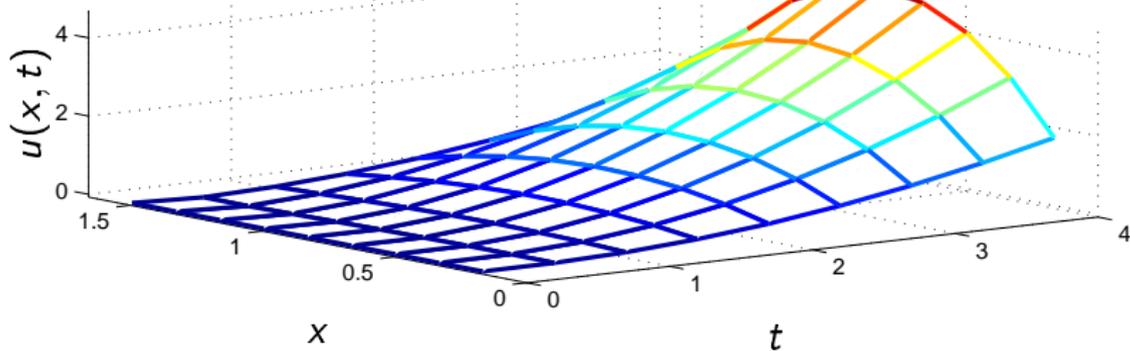
method	order	$\gamma_s$	$\gamma_I$
BE	1	0.2036321888	0.2984256075
SDIRK 3.1	3	0.1717941220	0.2338191487
SDIRK 3.2	3	0.2073822267	0.1718033767
Radau IIA	5	0.0634592650	0.0677592165

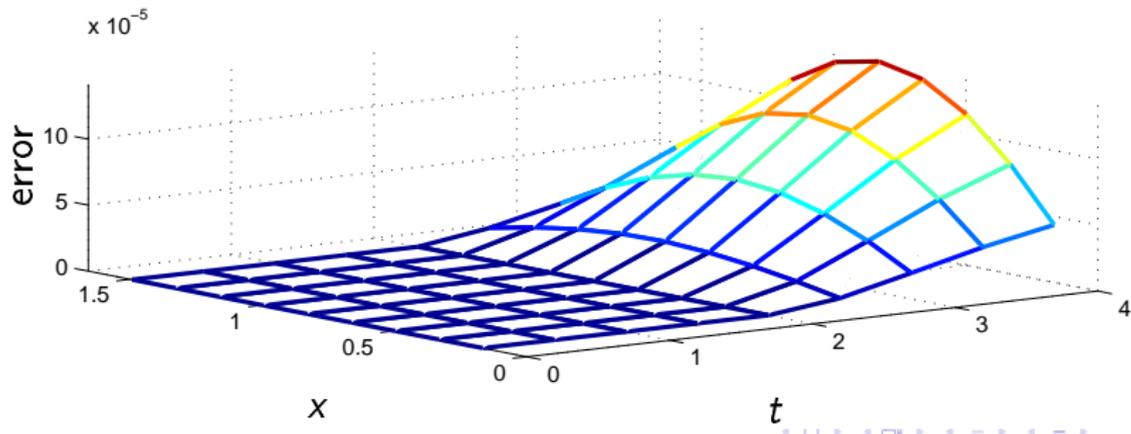
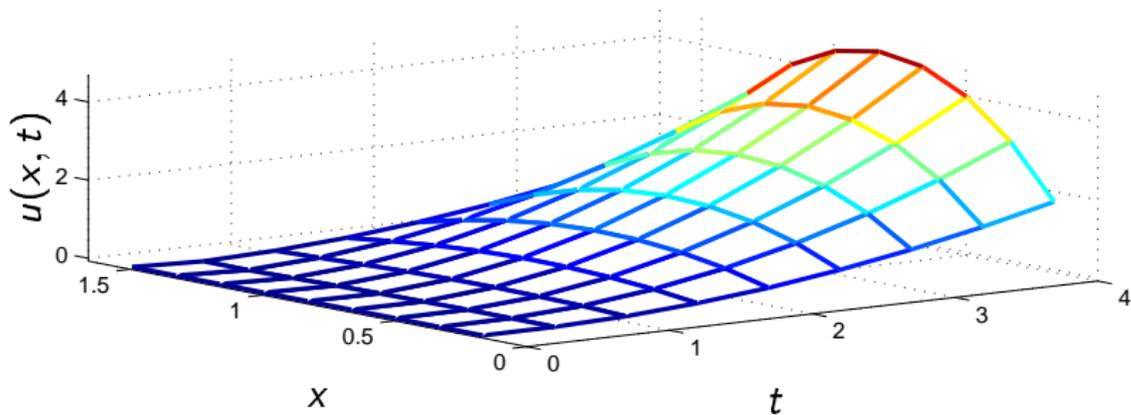
Note that higher order methods lead to faster convergence of the parareal algorithm.

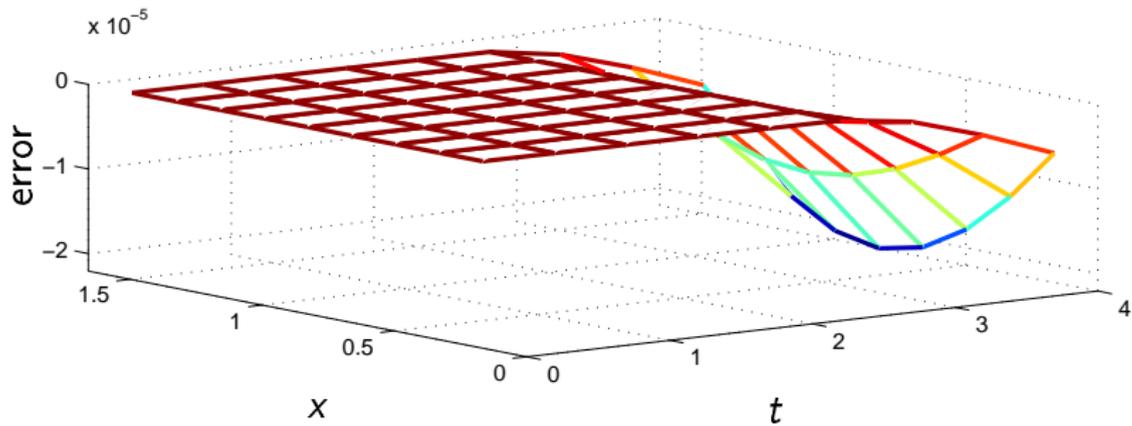
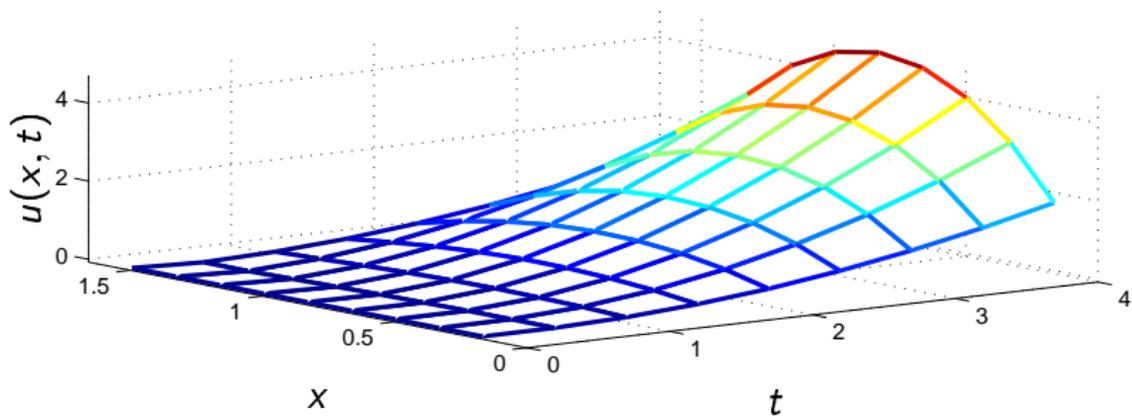


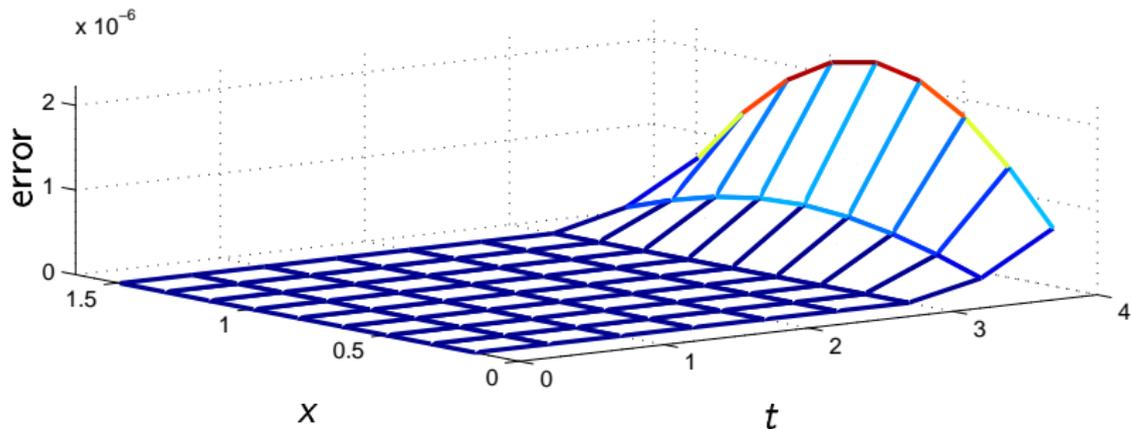
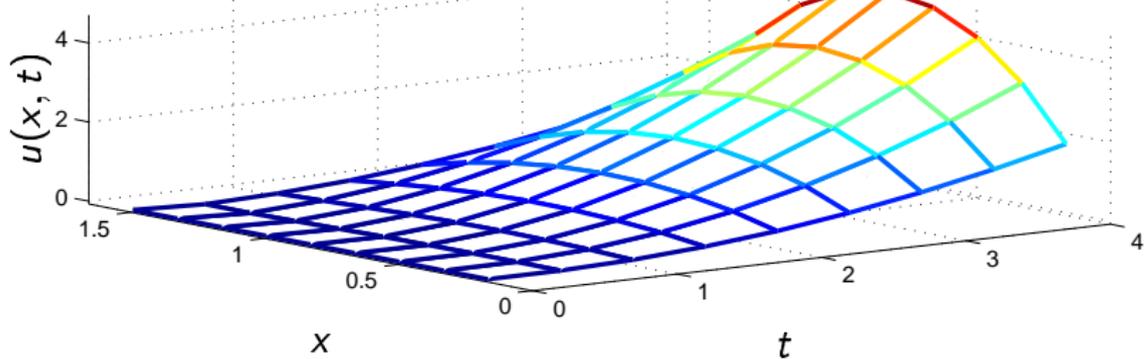


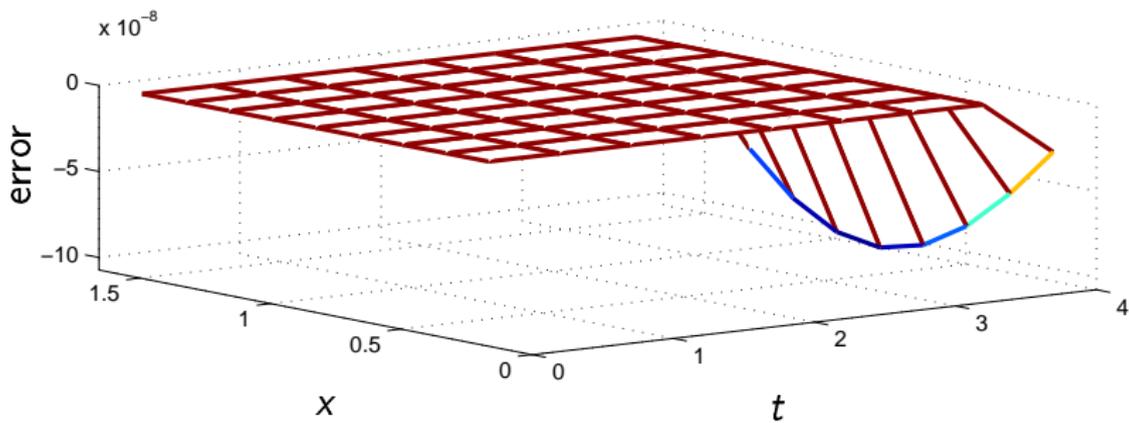
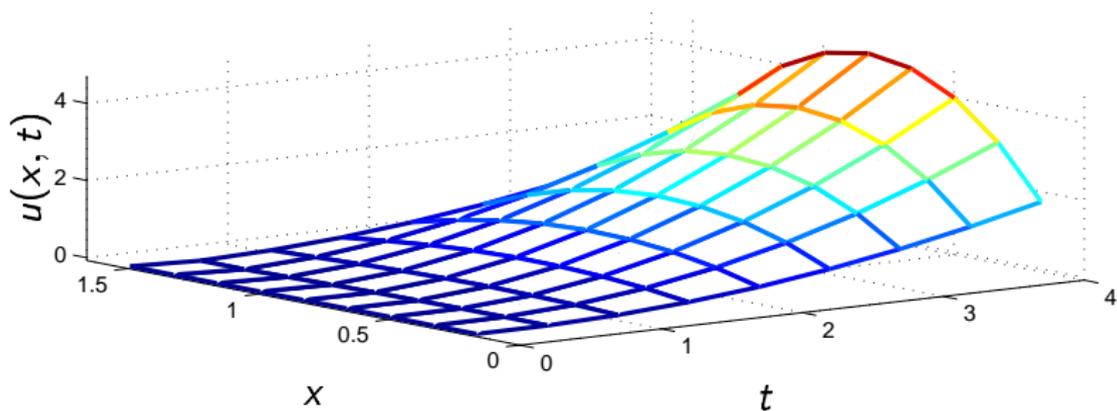


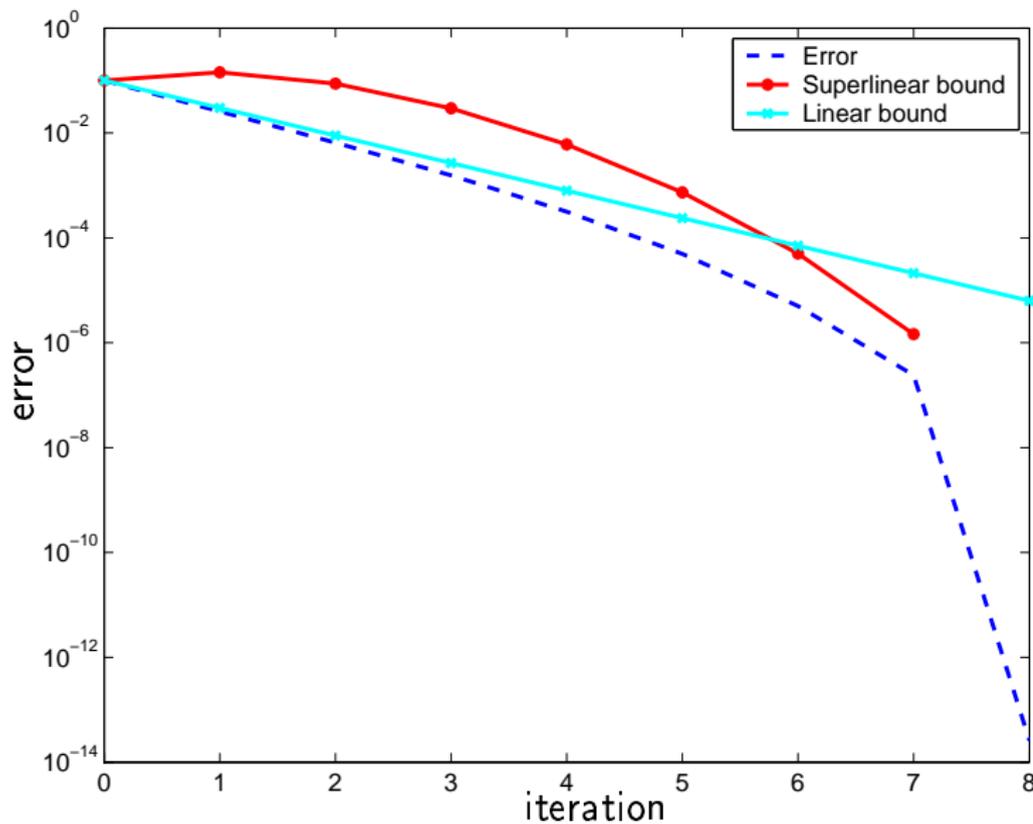


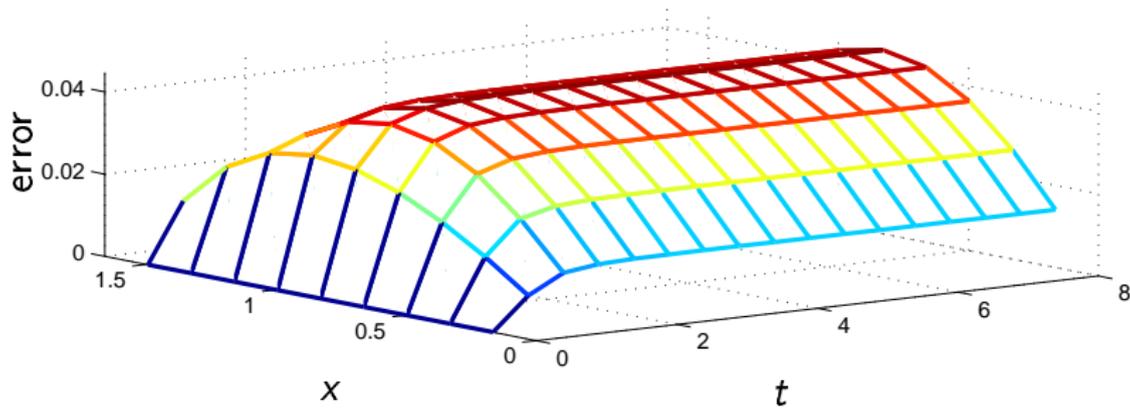
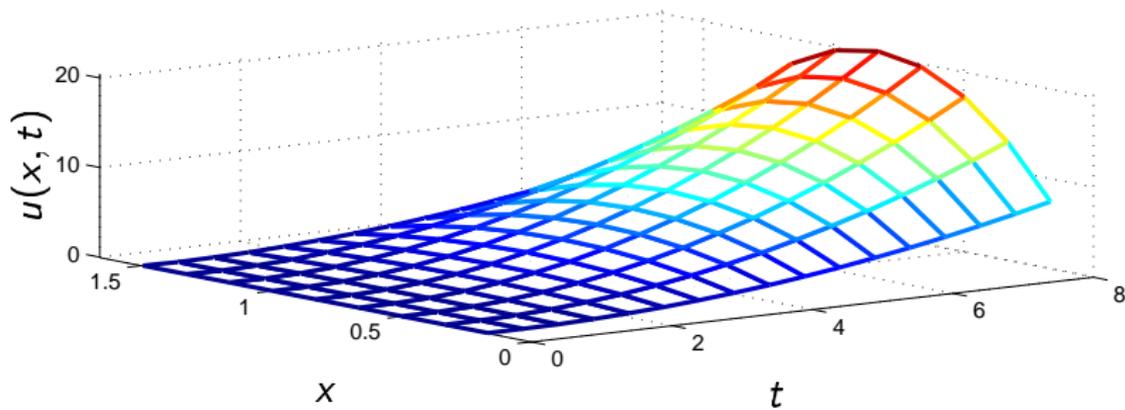


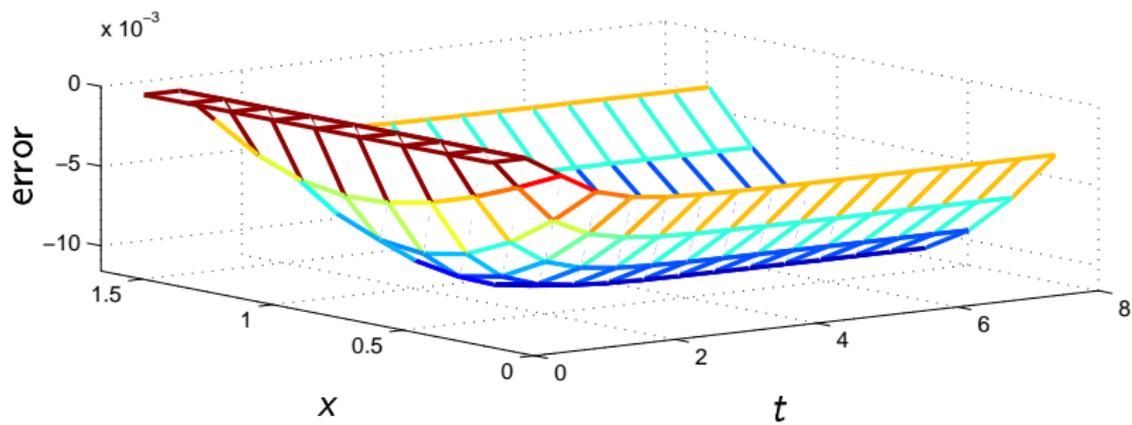
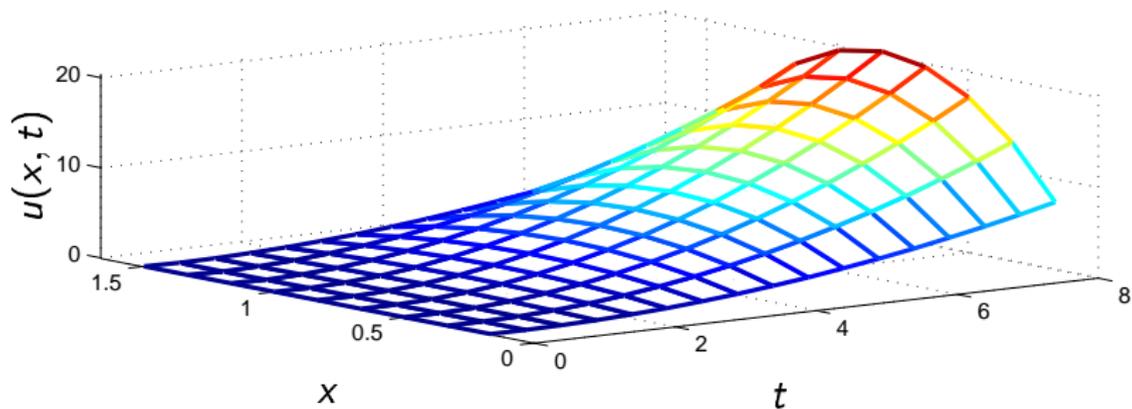


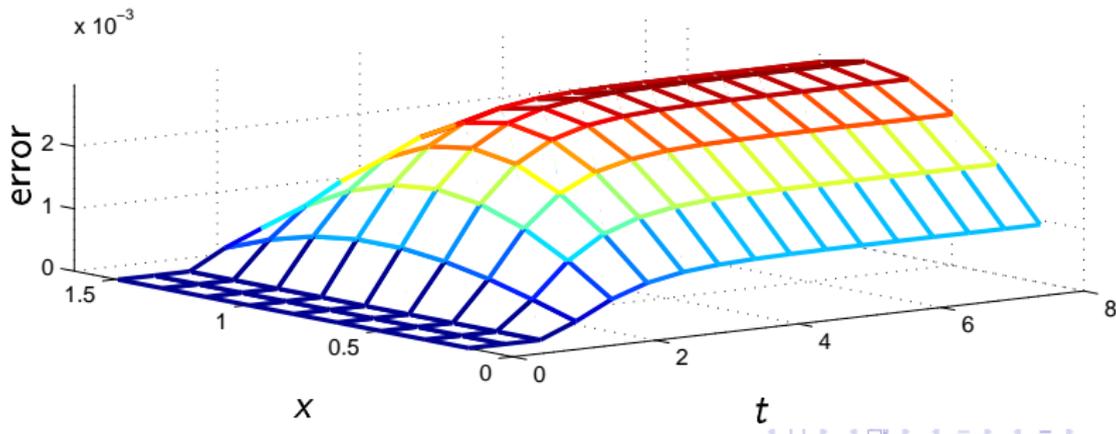
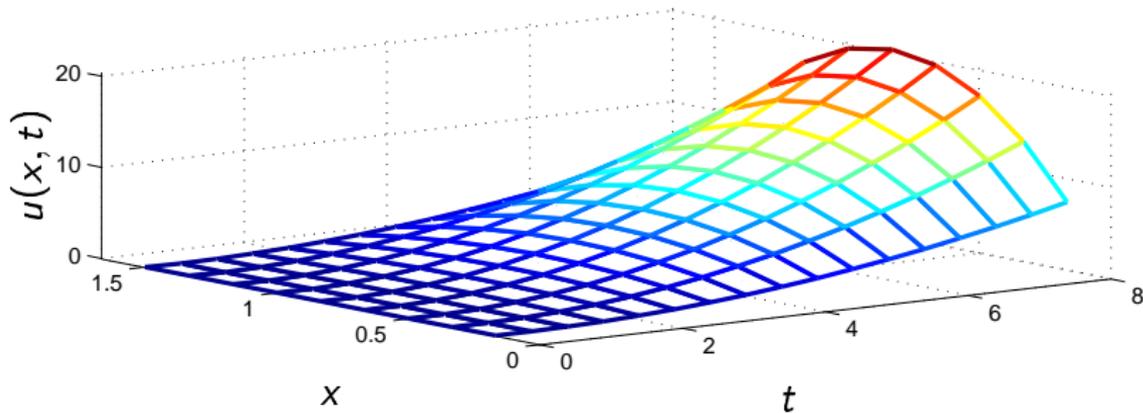


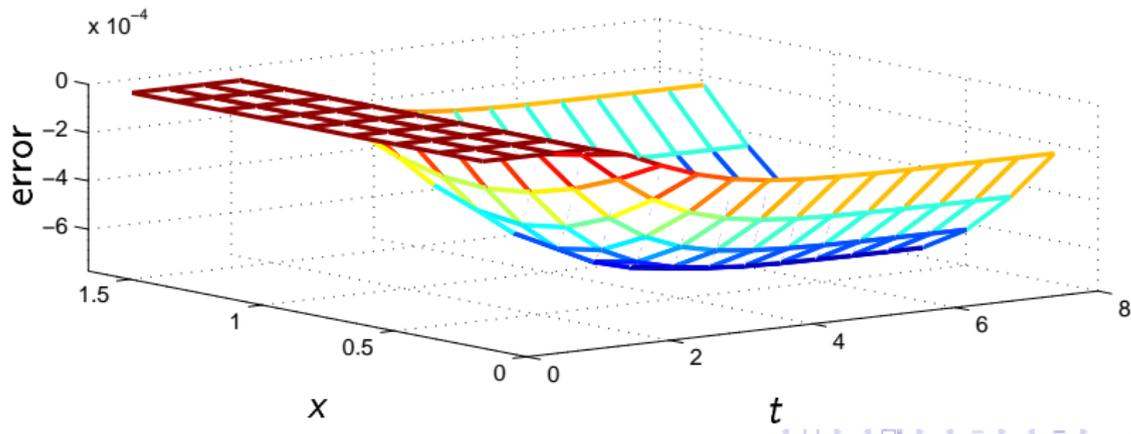
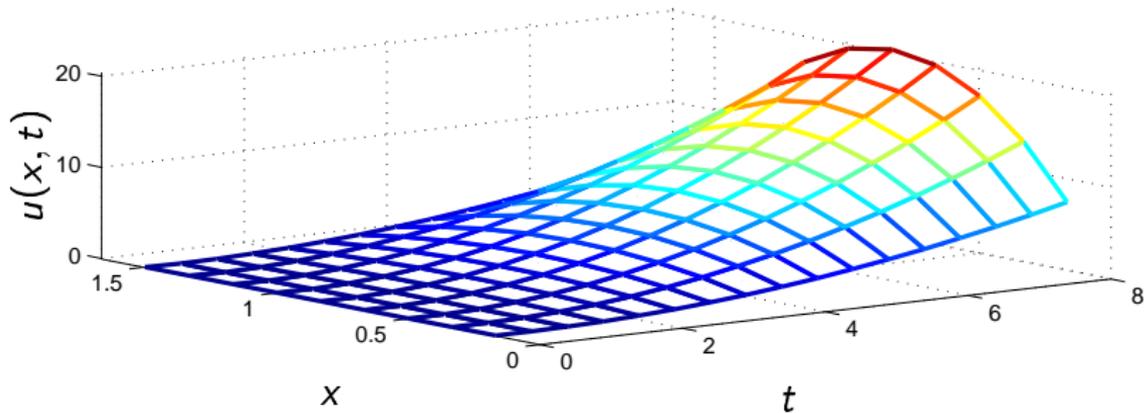


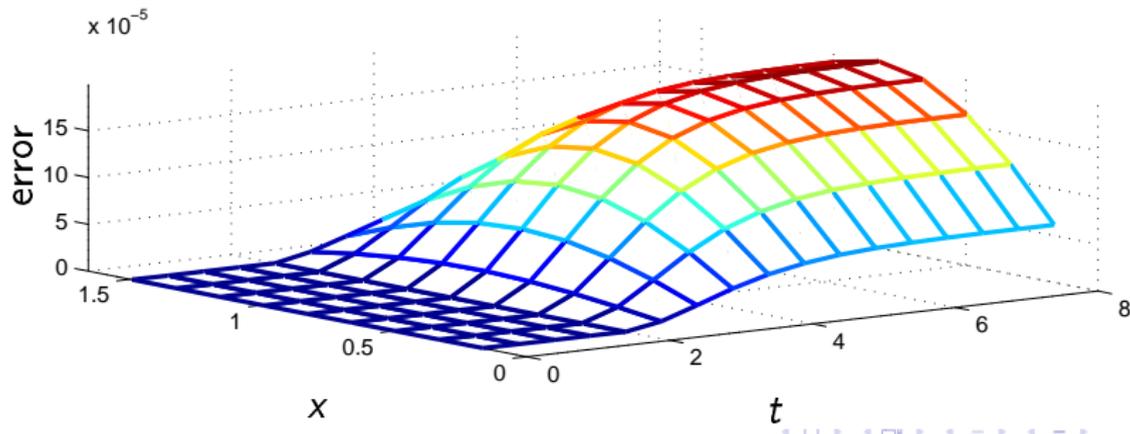
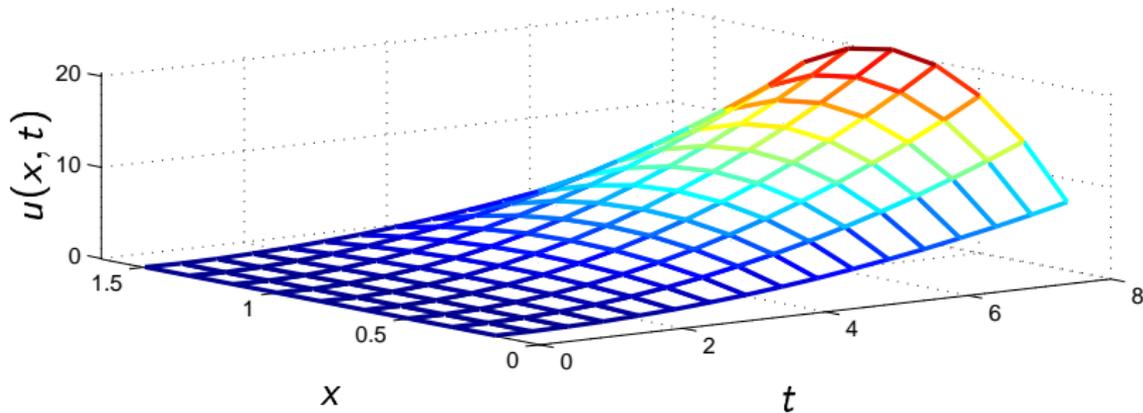


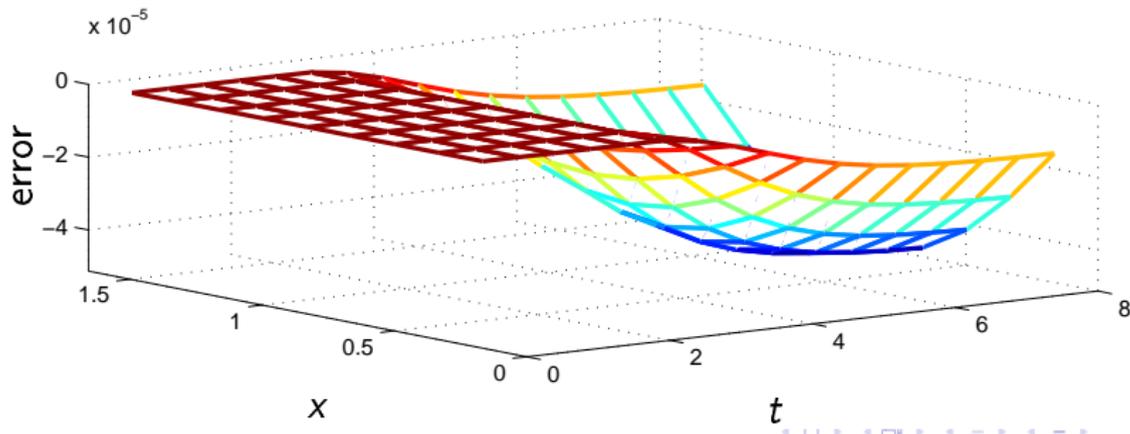
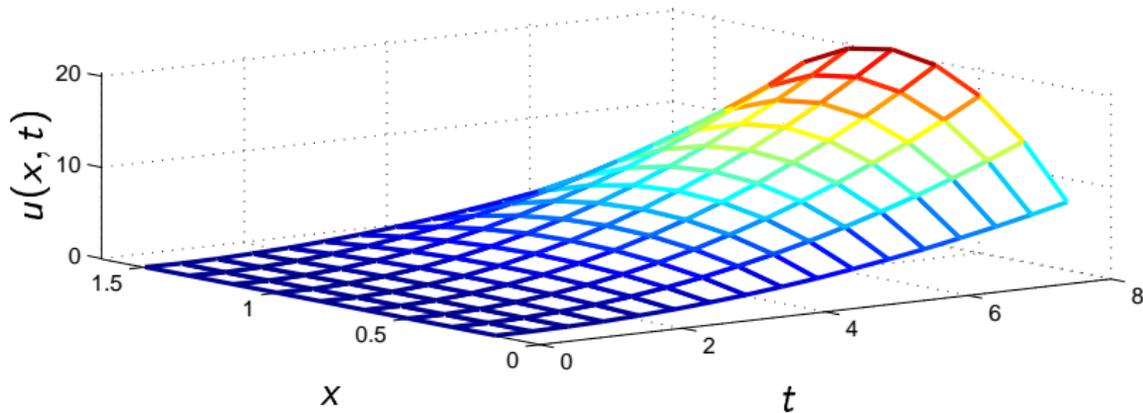


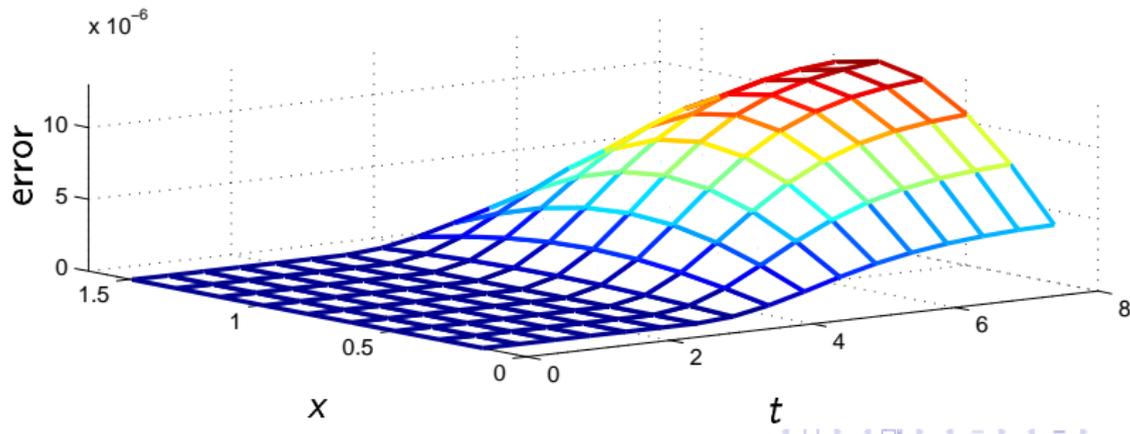
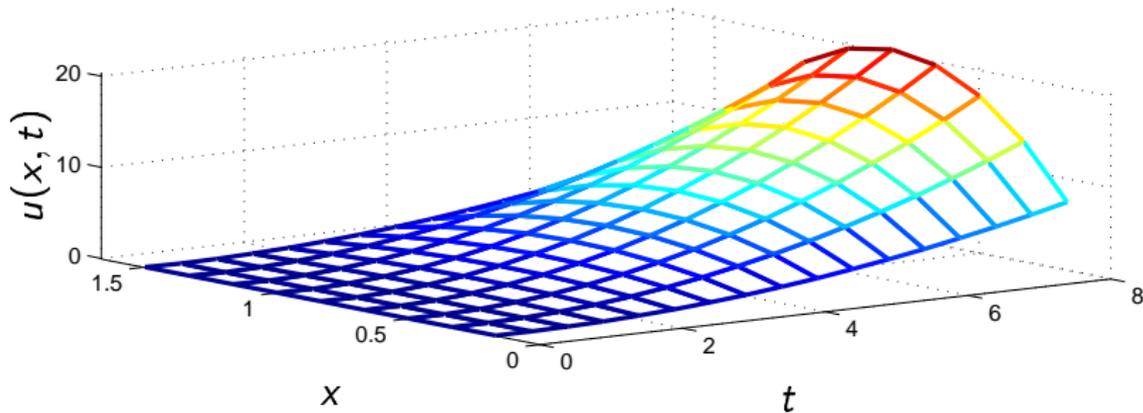


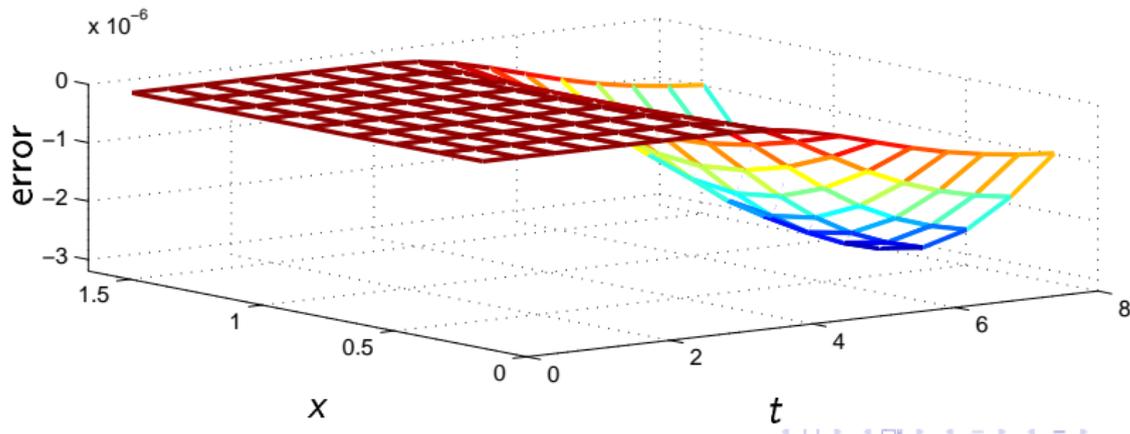
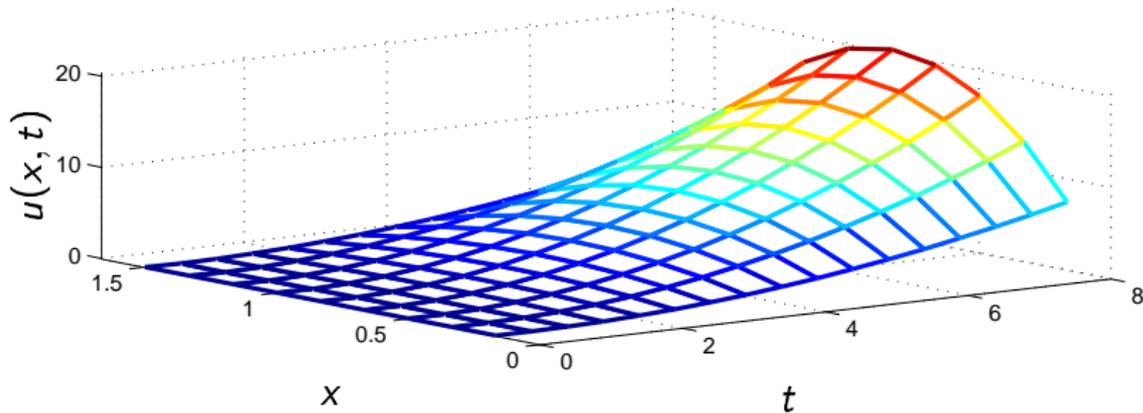


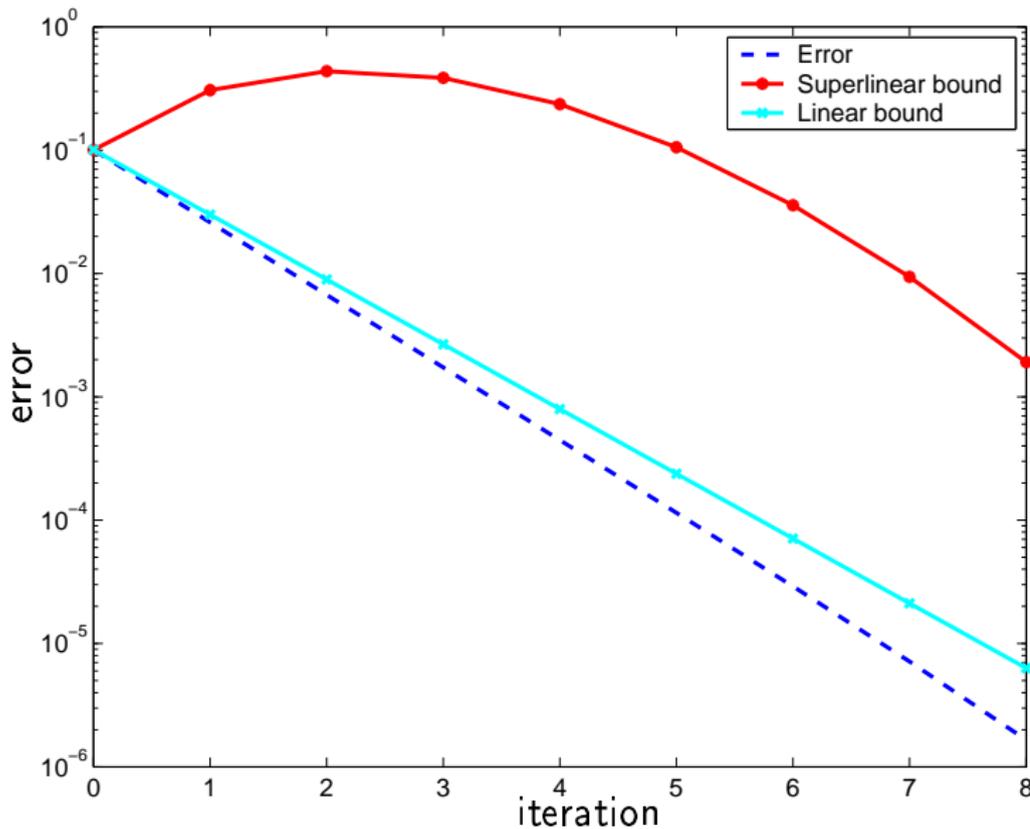




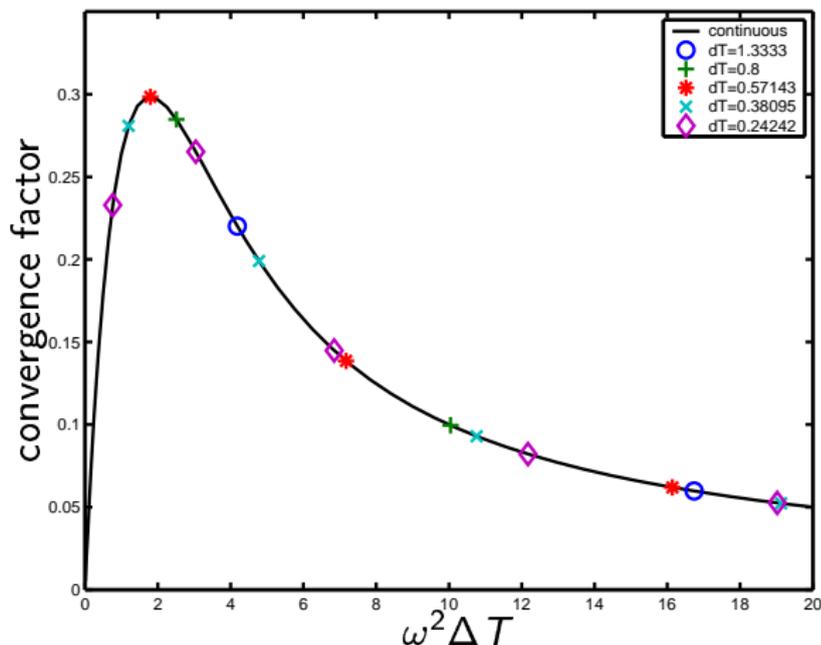






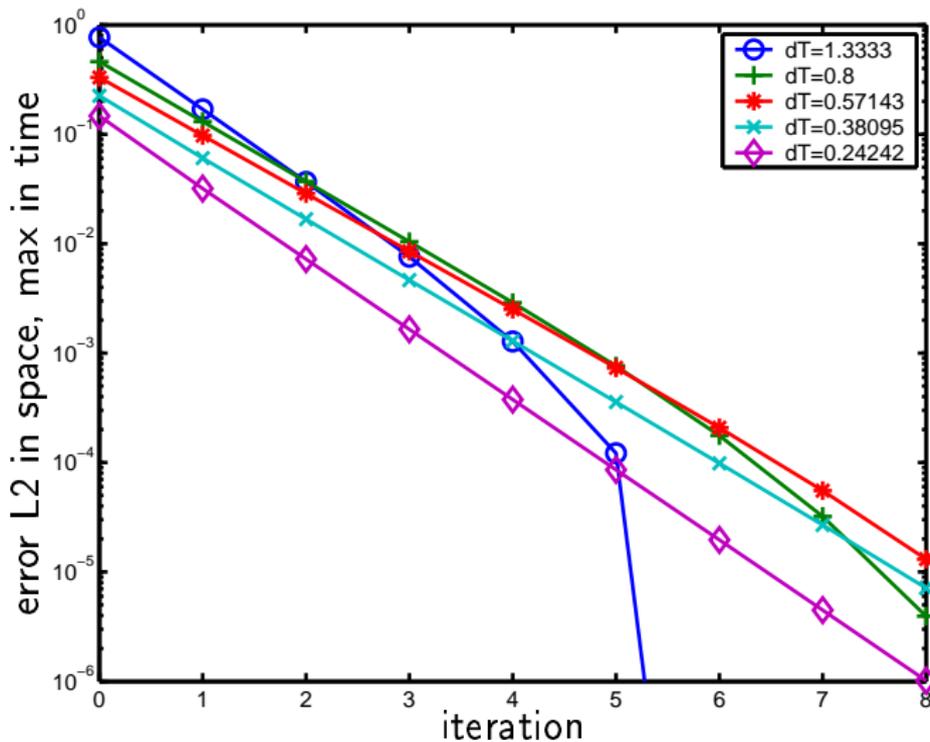


## The Results are Upper Bounds



$$\gamma_s := \sup_{\omega \in \mathbb{Z}} |e^{-\omega^2 \Delta T} - \beta(\omega)|, \quad \omega \text{ Fourier parameter in space}$$

# Corresponding Numerical Experiments



## Convergence for pure Advection Problems

### Corollary

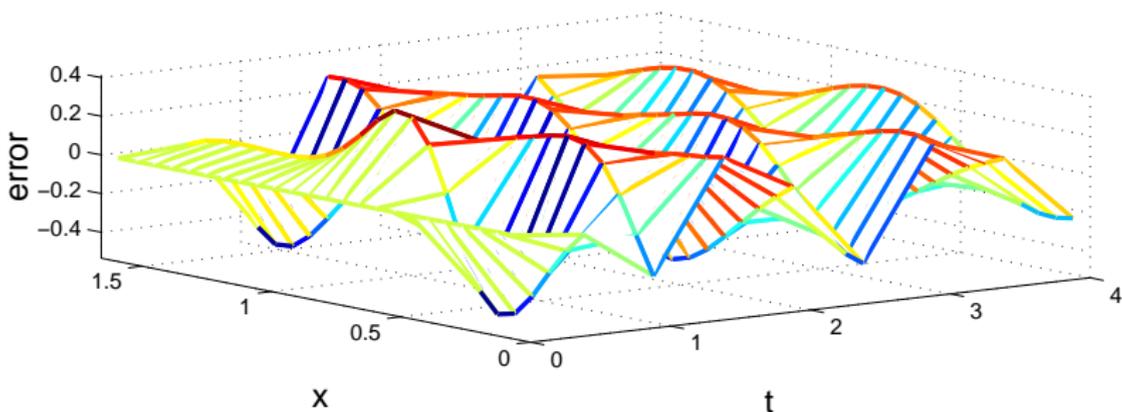
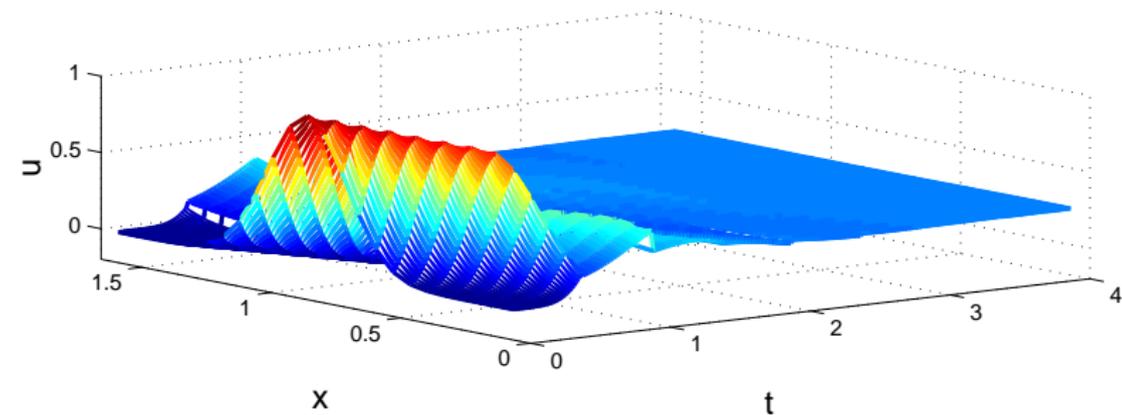
*The parareal algorithm applied to the advection equation  $u_t = u_x$  with backward Euler in time converges superlinearly on bounded time intervals,*

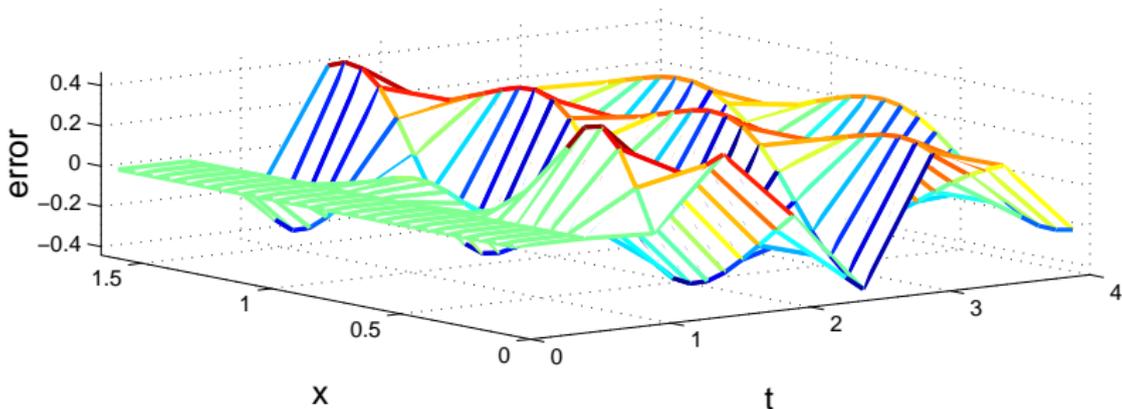
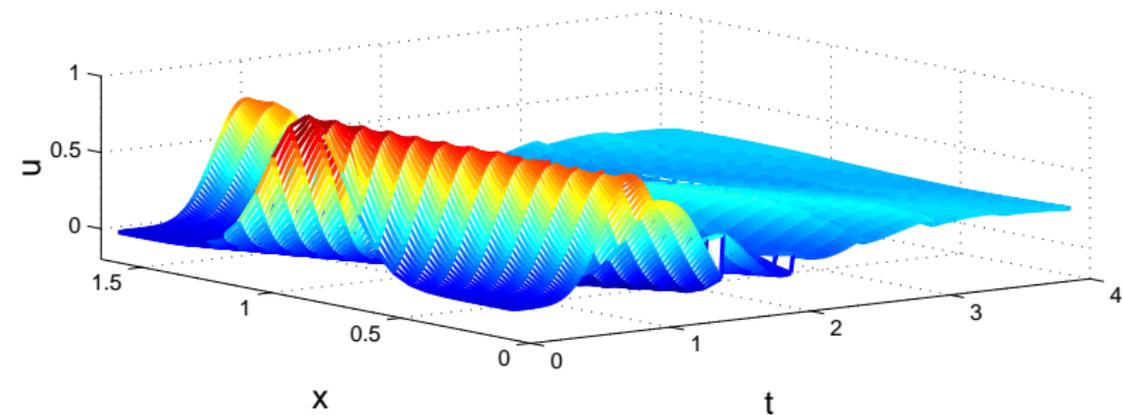
$$\max_{1 \leq n \leq N} \|u(t_n) - U_n^k\|_2 \leq \frac{\alpha_s^k}{k!} \prod_{j=1}^k (N - j) \max_{1 \leq n \leq N} \|u(t_n) - U_n^0\|_2,$$

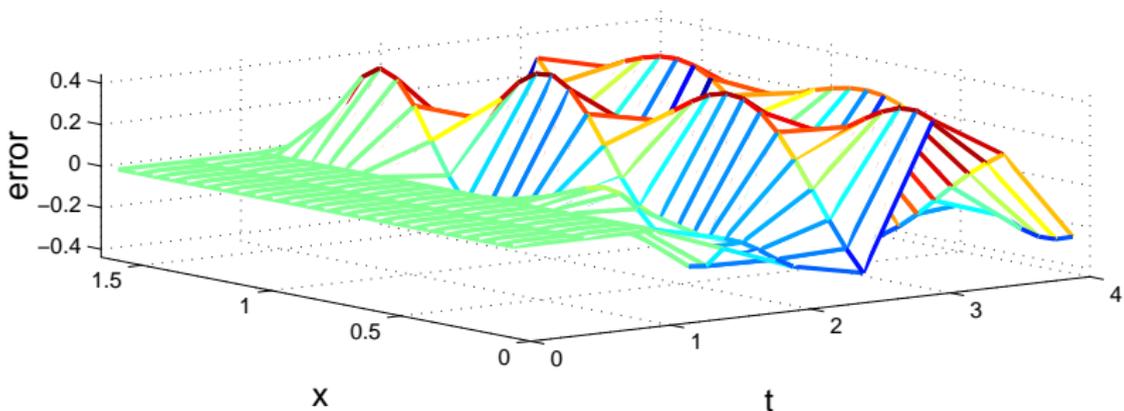
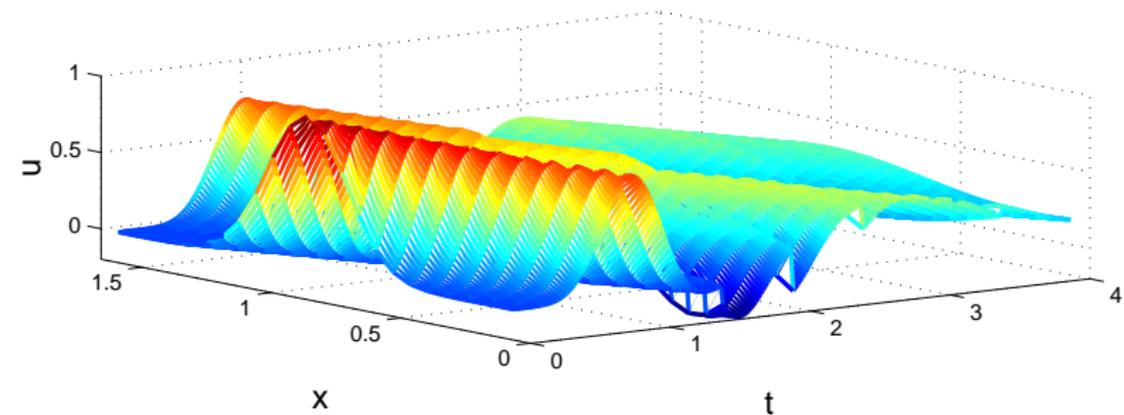
where the constant  $\alpha_s$  is universal,  $\alpha_s = 1.224353426$ .

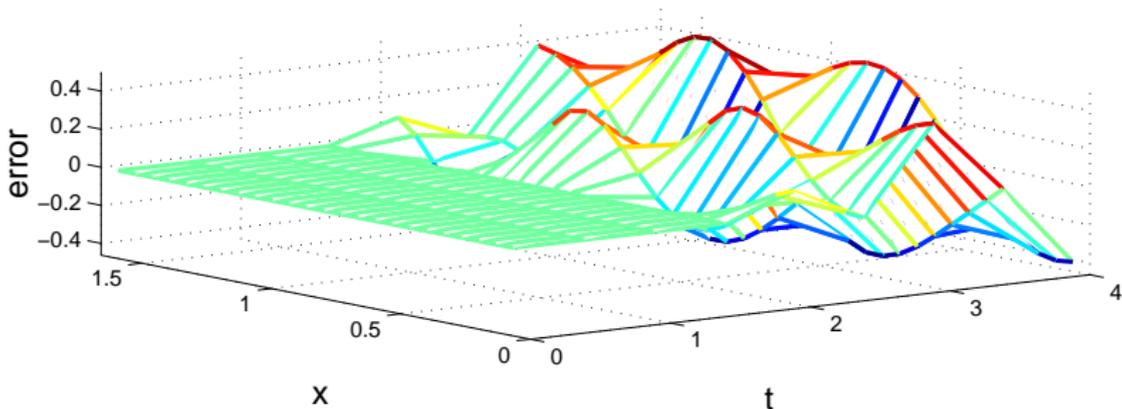
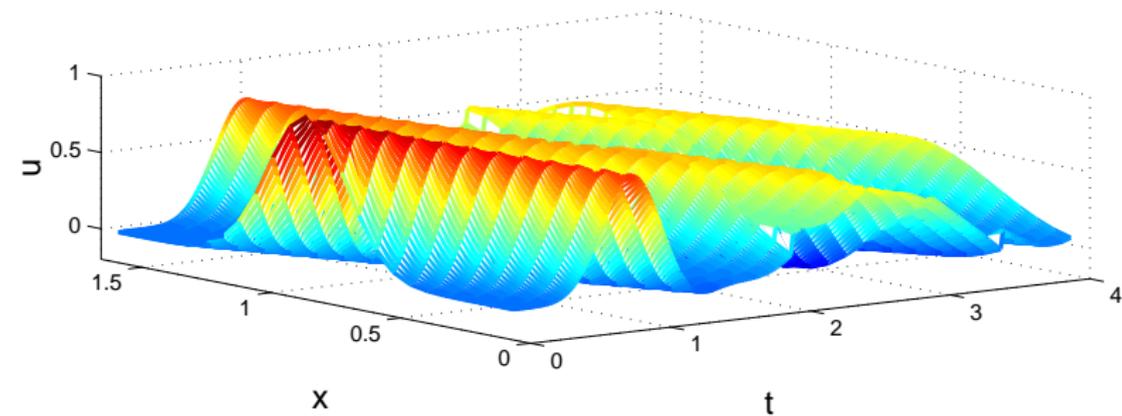
### Remarks:

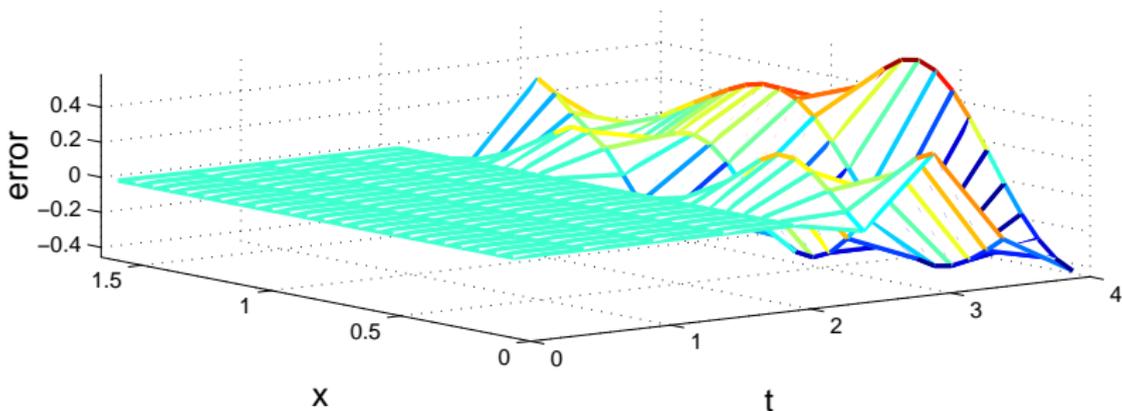
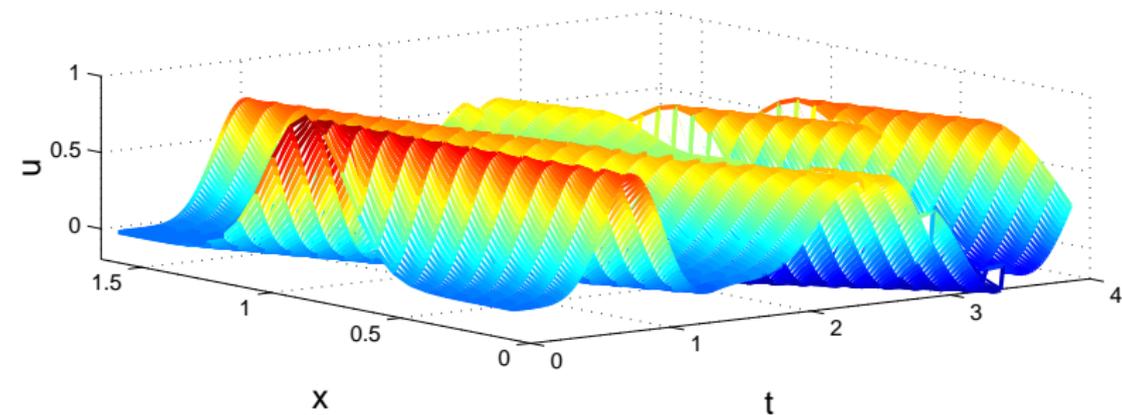
- ▶ No convergence result for unbounded time intervals.
- ▶ As soon as more than  $N$  iterations are needed, the method loses all interest.

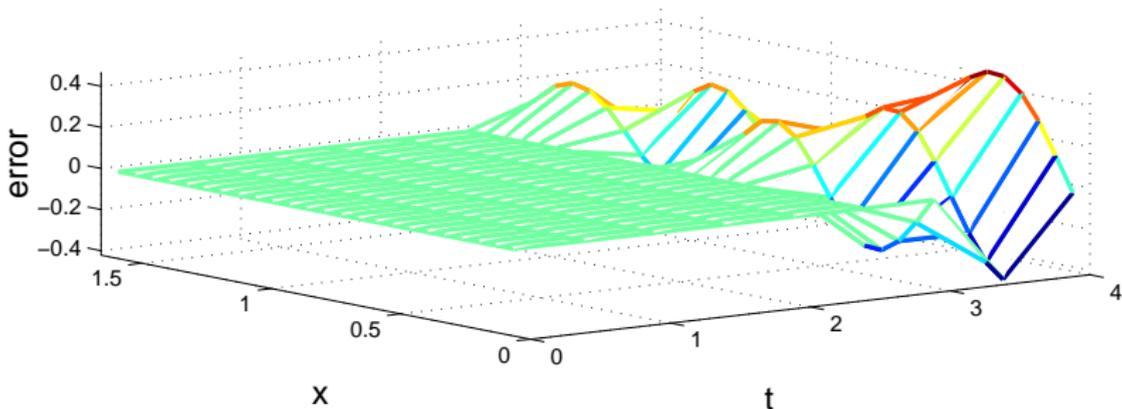
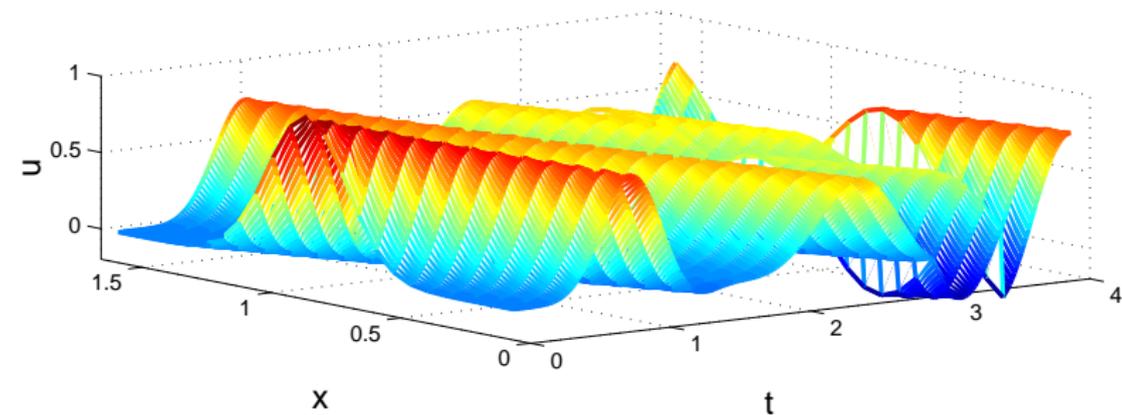


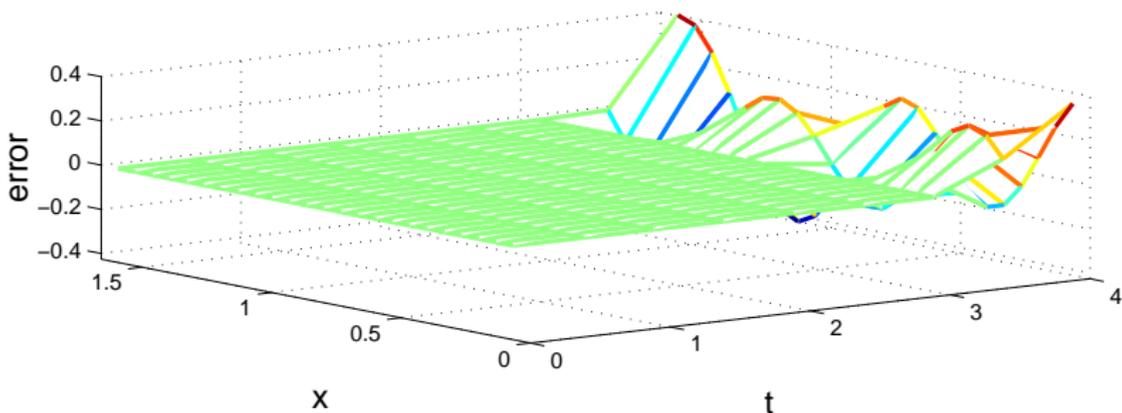
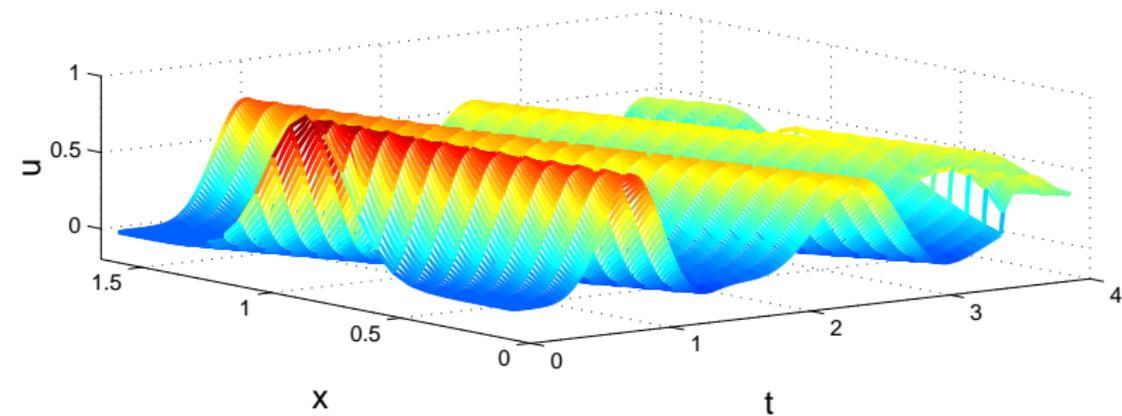


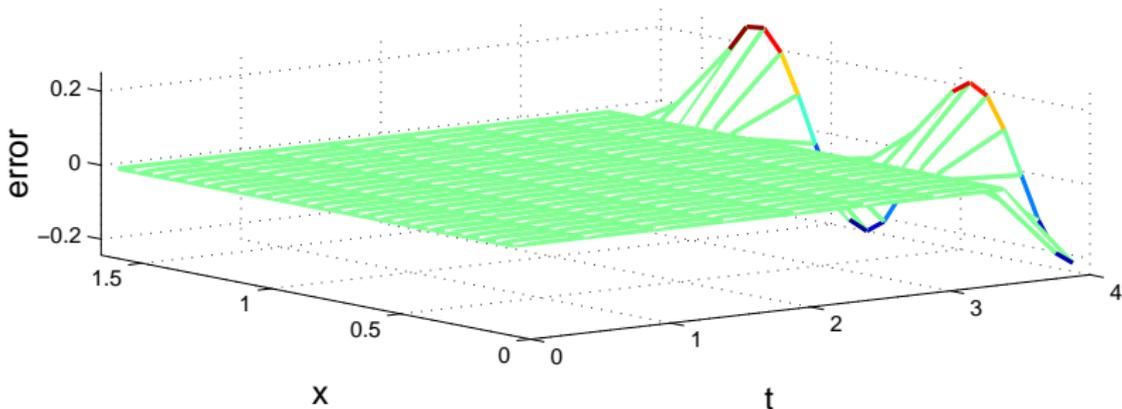
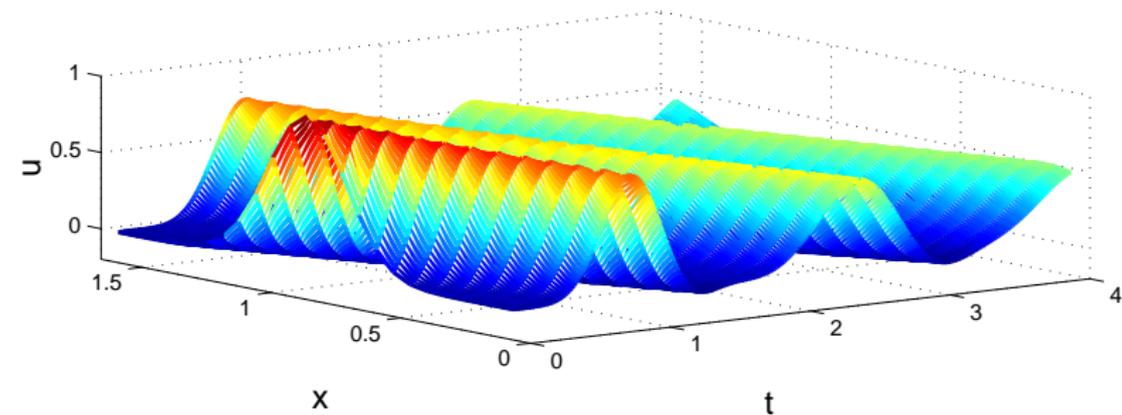


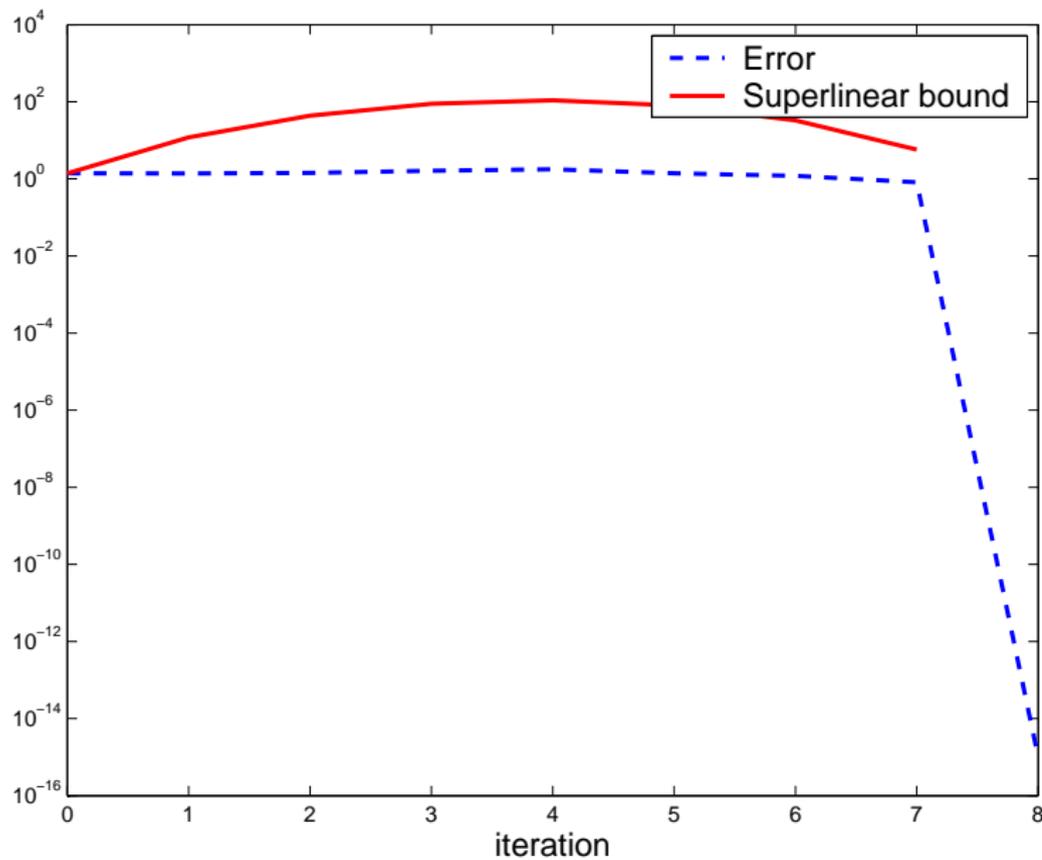












## Speedup

We define the speedup of the parareal algorithm by  $N/k$ , where  $N$  is the number of processors (=number of coarse intervals), and  $k$  is the number of iterations to achieve a given precision  $\varepsilon$ .

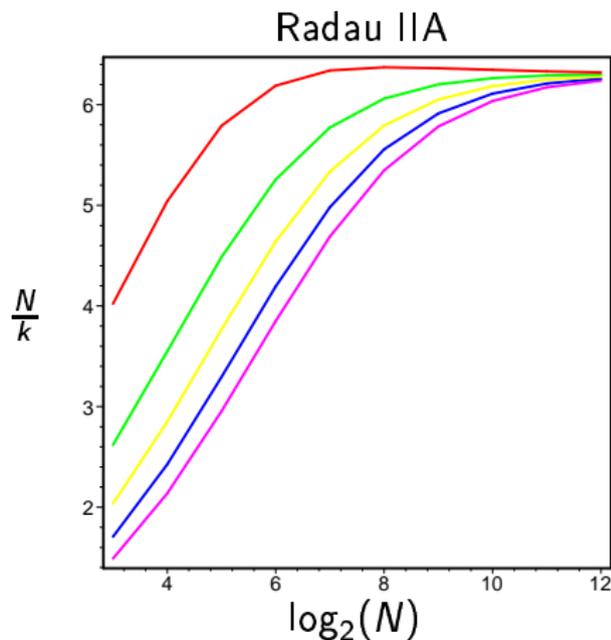
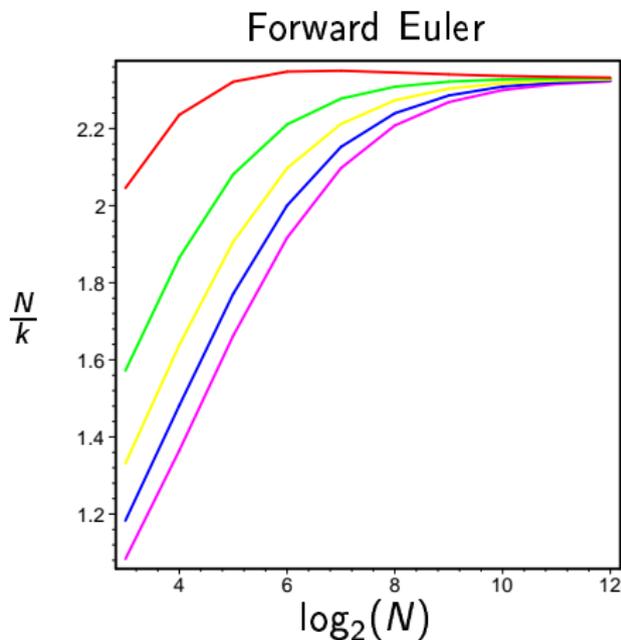
To quantify the speedup of the parareal algorithm, we need to study for  $k < N$  the function

$$f(\gamma, N, k) := \frac{\gamma^k}{k!} \prod_{j=1}^k (N - j) \leq \frac{\gamma^k}{k!} \frac{N^N}{e^k (N - k)^{(N-k)}}.$$

Goal: for a given  $\gamma$  from an L-stable method, and a desired precision  $\varepsilon$ , find  $N$ , such that the speedup  $N/k$  is maximized.

## Speedup for the Heat Equation

For given precision  $\varepsilon \in \{1/10, 1/100, \dots, 1/100000\}$ , the speedup  $N/k$  as a function of the number of coarse intervals  $N$ :



## Speedup for the Advection Equation ?

We have

$$\lim_{k \rightarrow N} f(\gamma, N, k) \leq \frac{\gamma^N N^N}{N! e^N} =: \bar{f}(\gamma, N),$$

and hence for a given  $N$ , speedup is possible if

$$\bar{f}(\gamma, N) < 1.$$

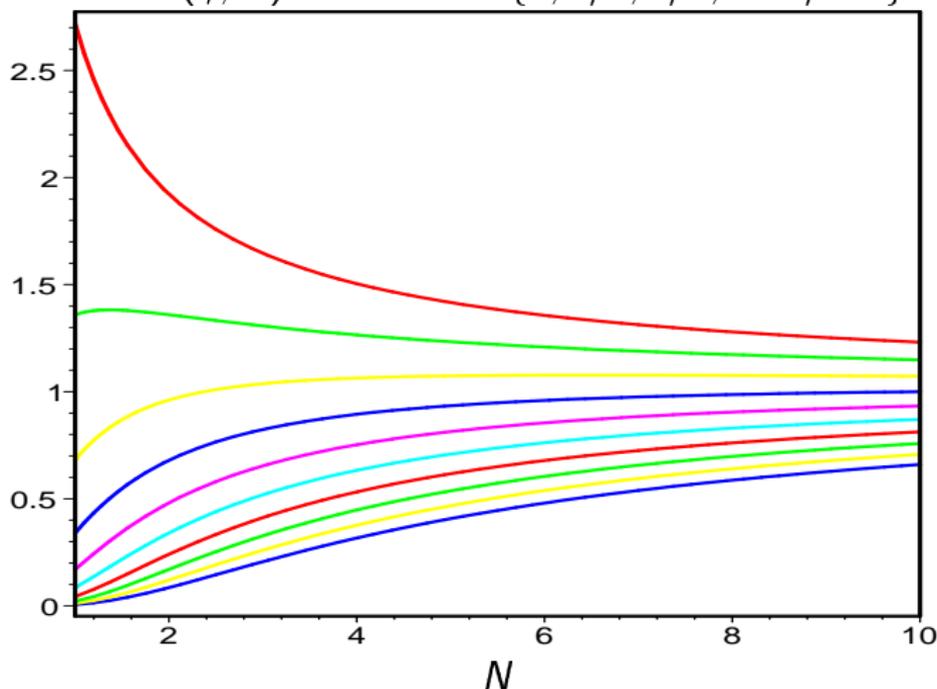
The limiting case  $\bar{f}(\gamma, N) = 1$  defines the function

$$\gamma = \gamma(N) = \frac{e(N!)^{\frac{1}{N}}}{N}.$$

and for  $\gamma$ -values above this curve, speedup is not possible.

## Possible Speedup Depending on $\gamma$

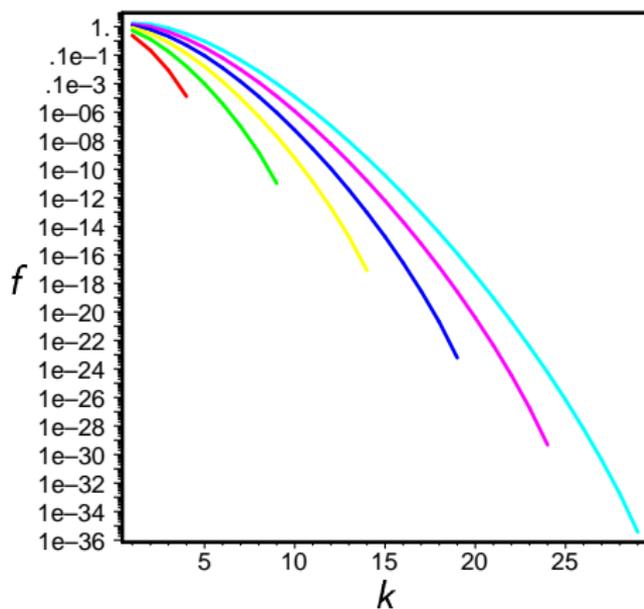
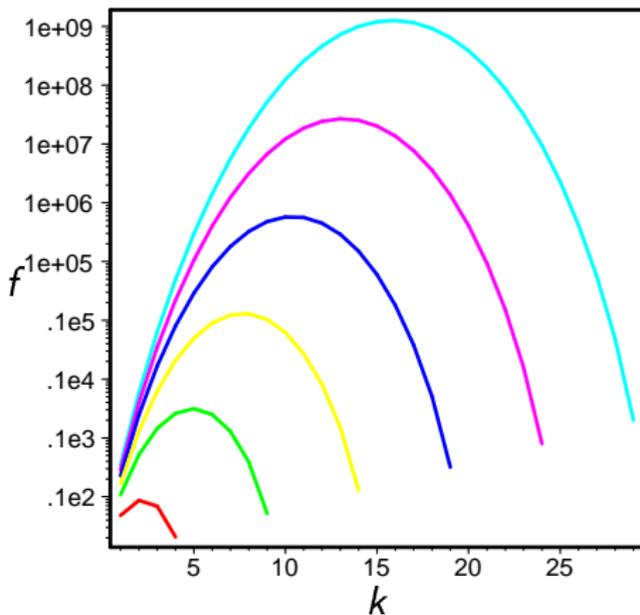
Curves  $\bar{f}(\gamma, N) = \varepsilon$  for  $\varepsilon = \{1, 1/2, 1/4, \dots, 1/512\}$ .



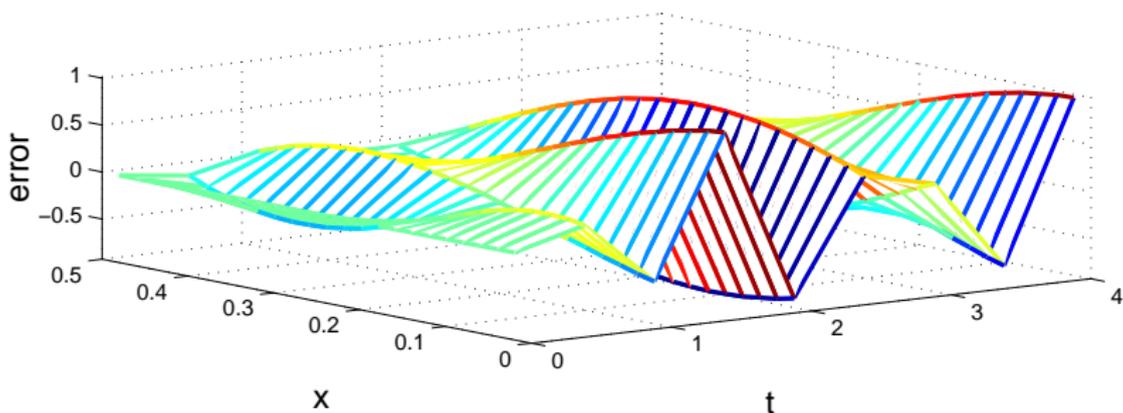
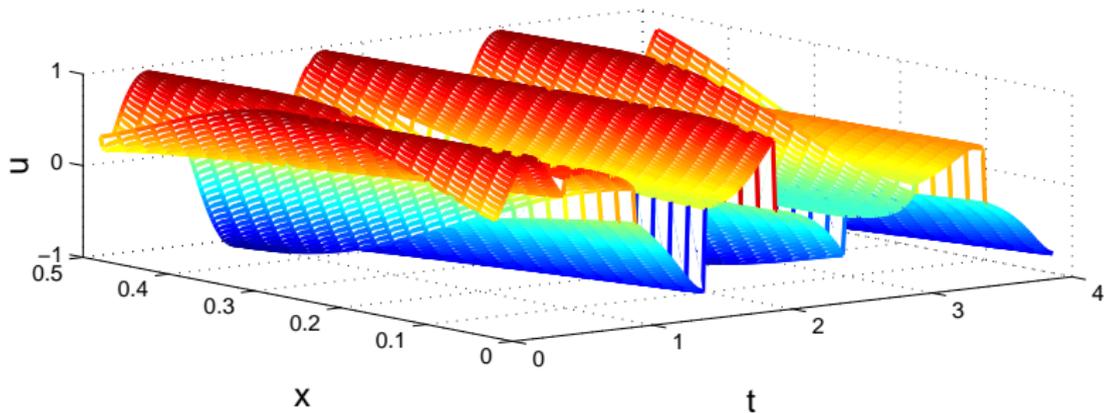
Note: the limit for  $N$  large is 1 for all  $\varepsilon$ .

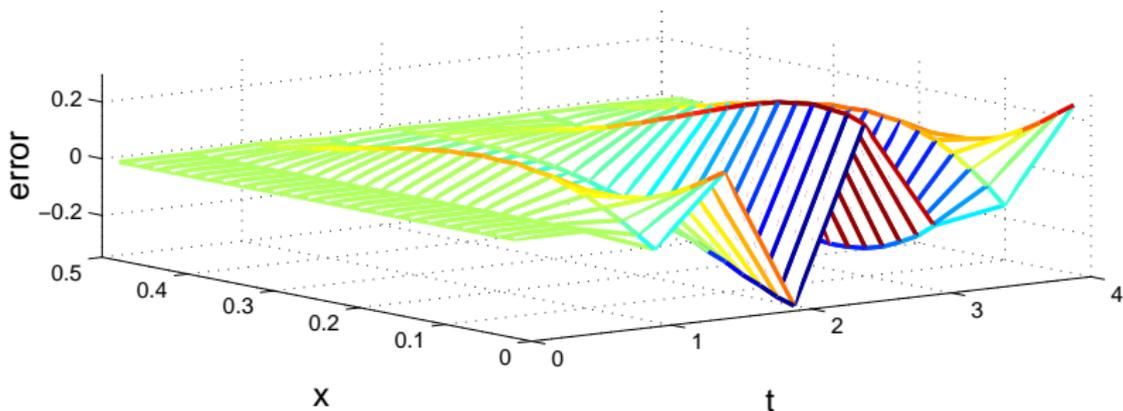
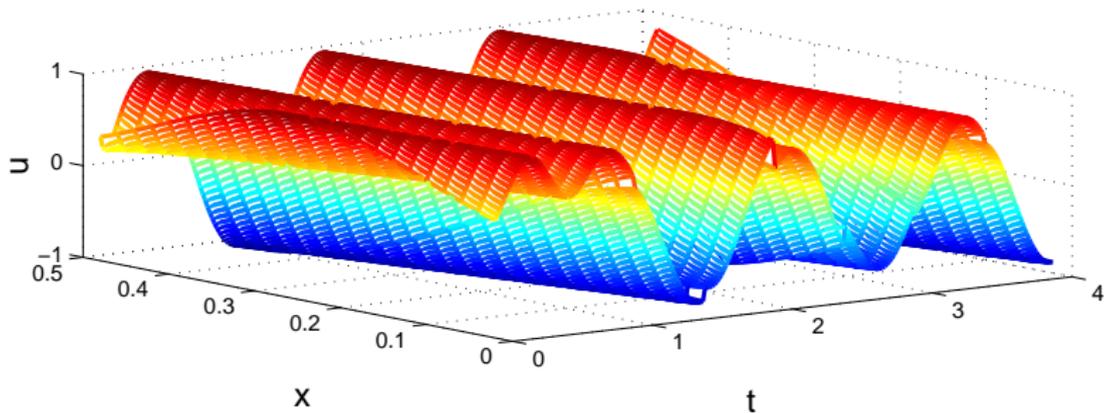
# Speedup Comparison between Advection and Diffusion

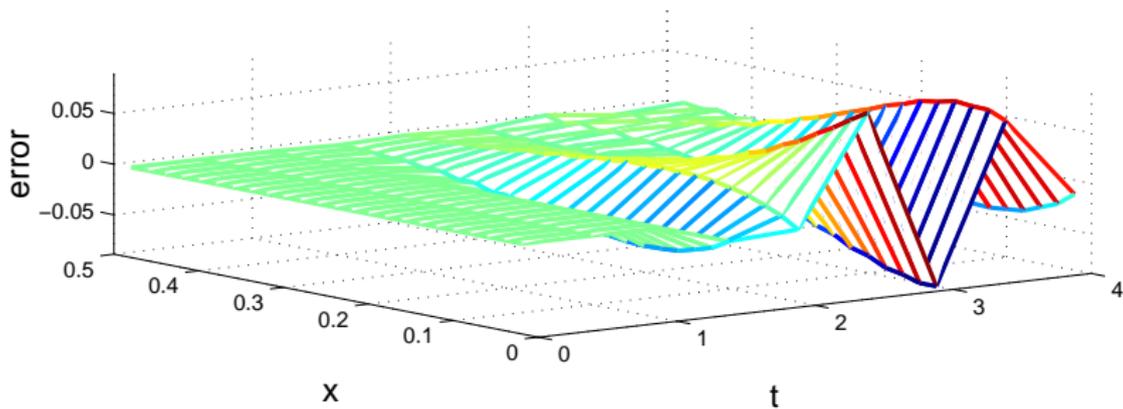
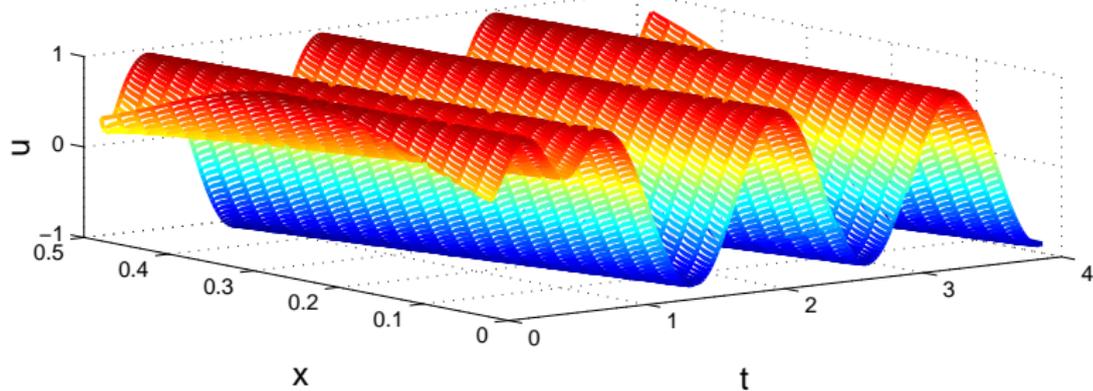
Convergence factors for  $N = 5, 10, 15, 20, 25, 30$ :

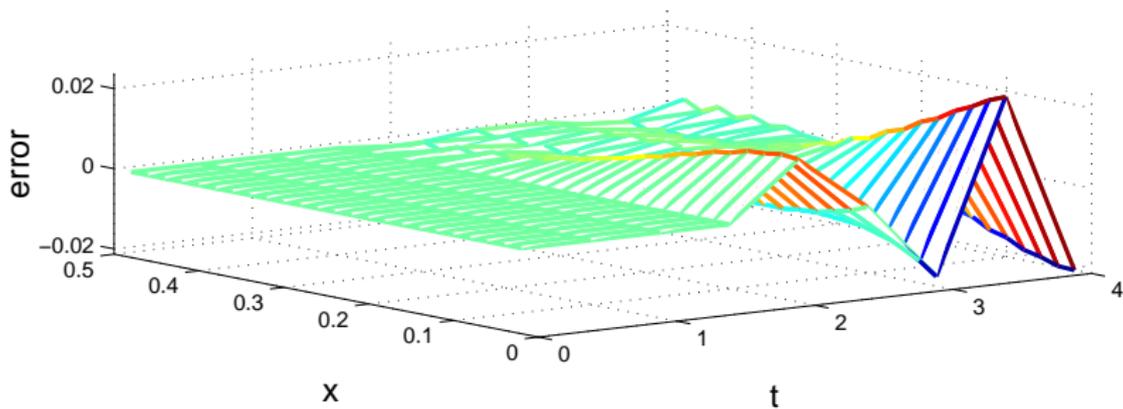
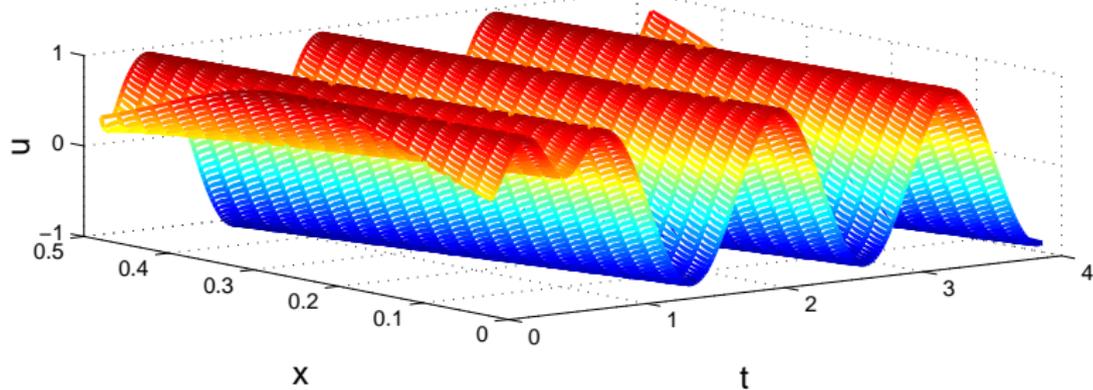


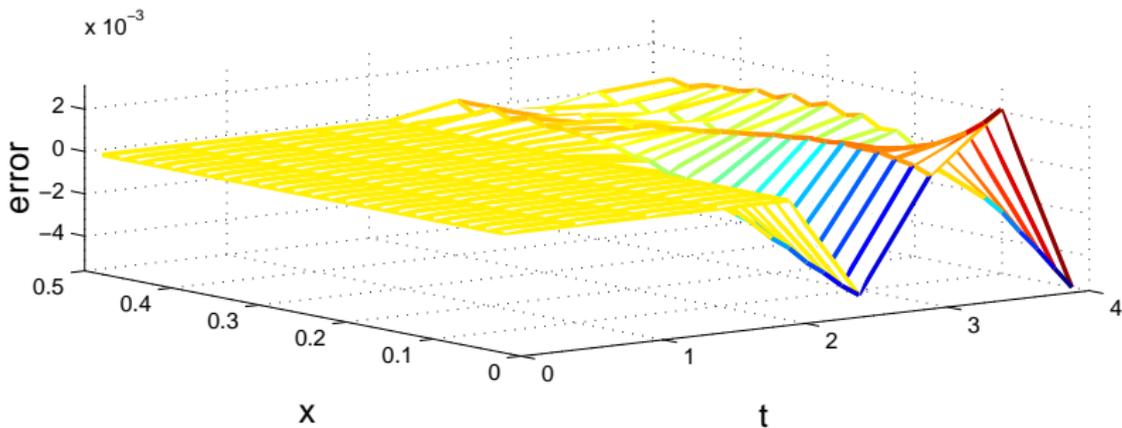
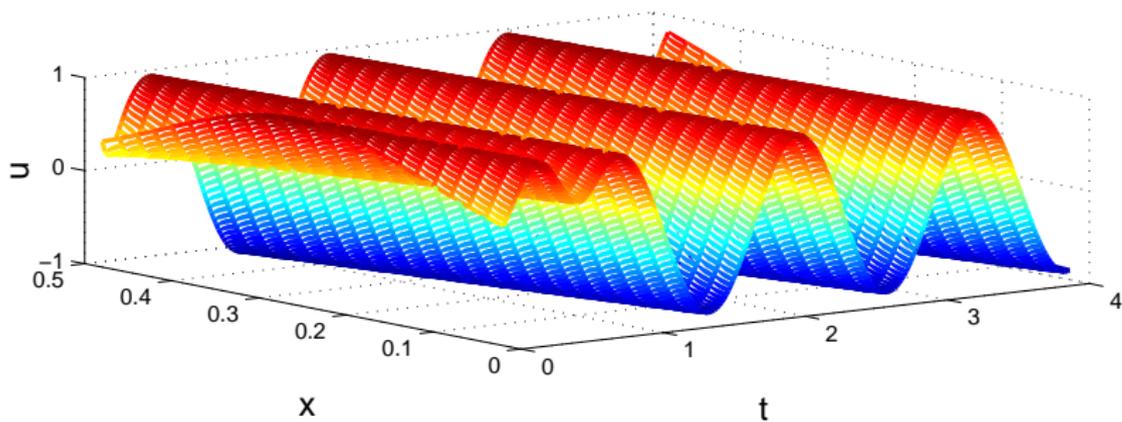
Advection case on the left, diffusion case on the right.

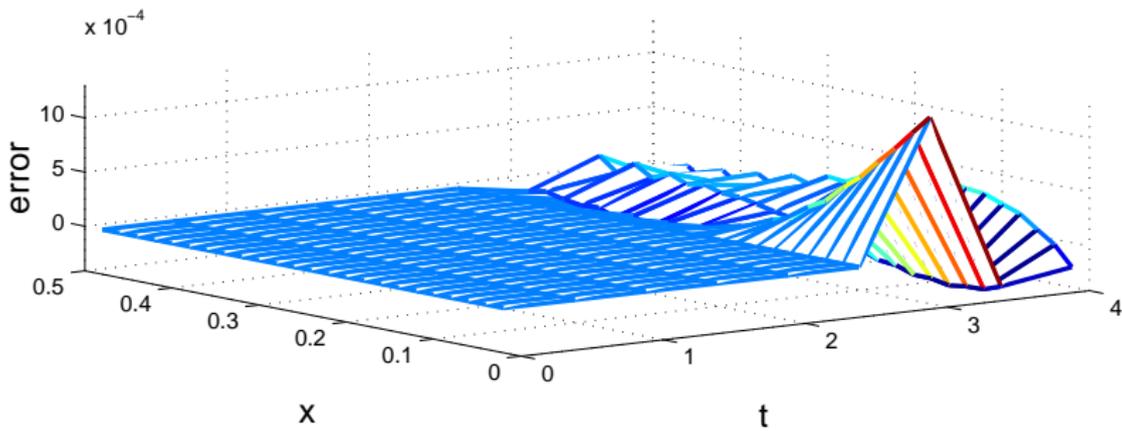
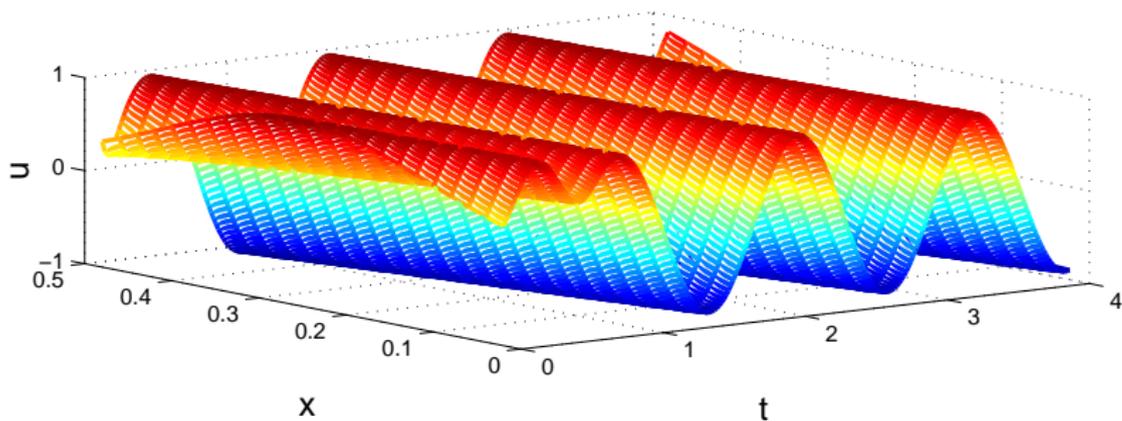


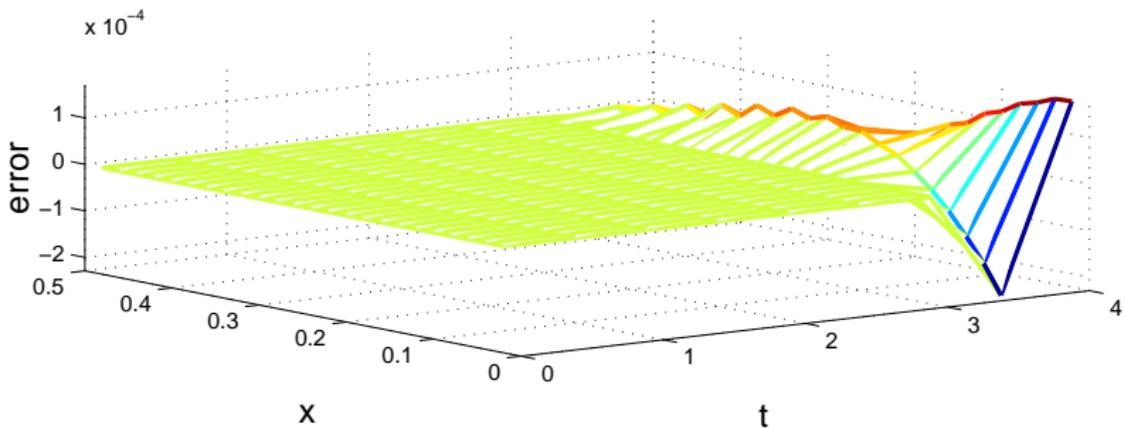
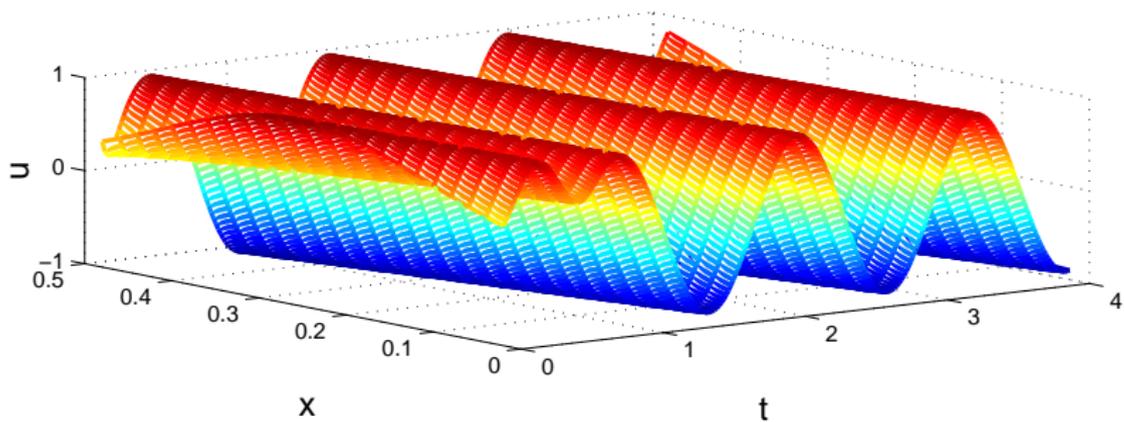


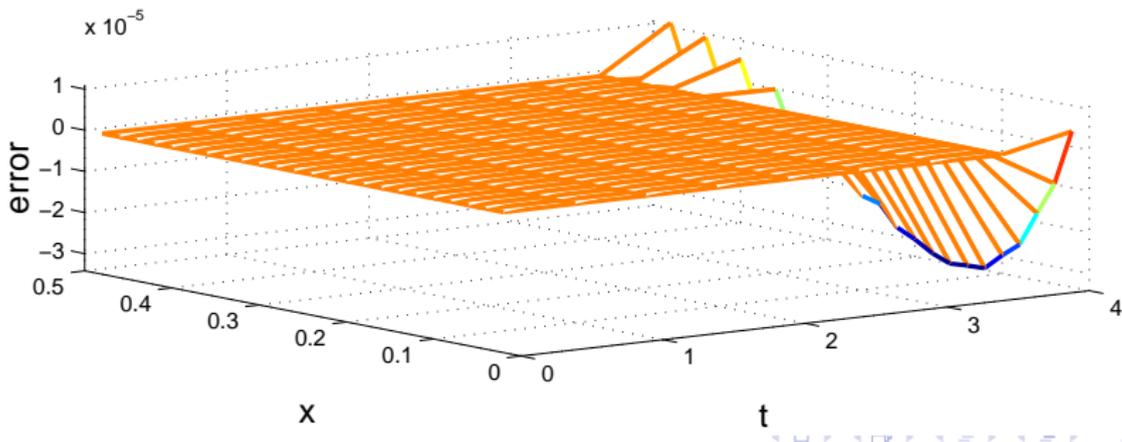
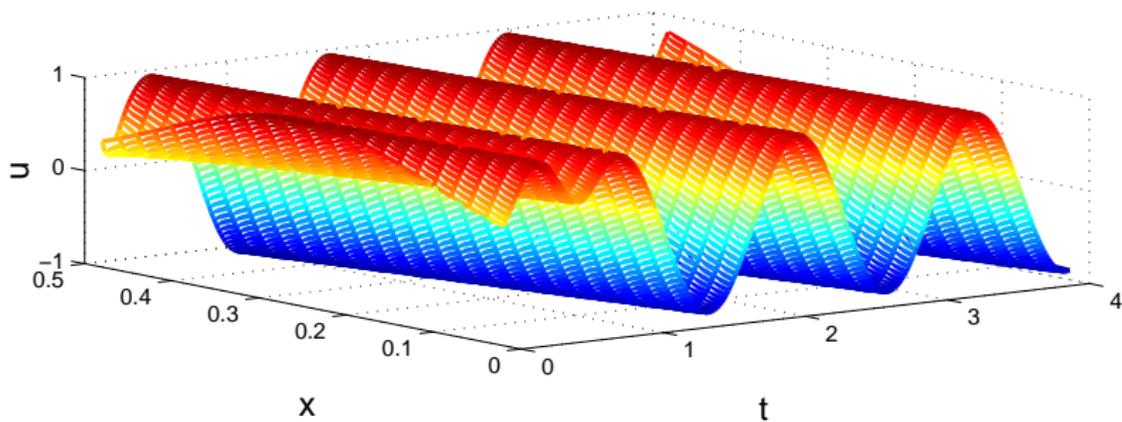


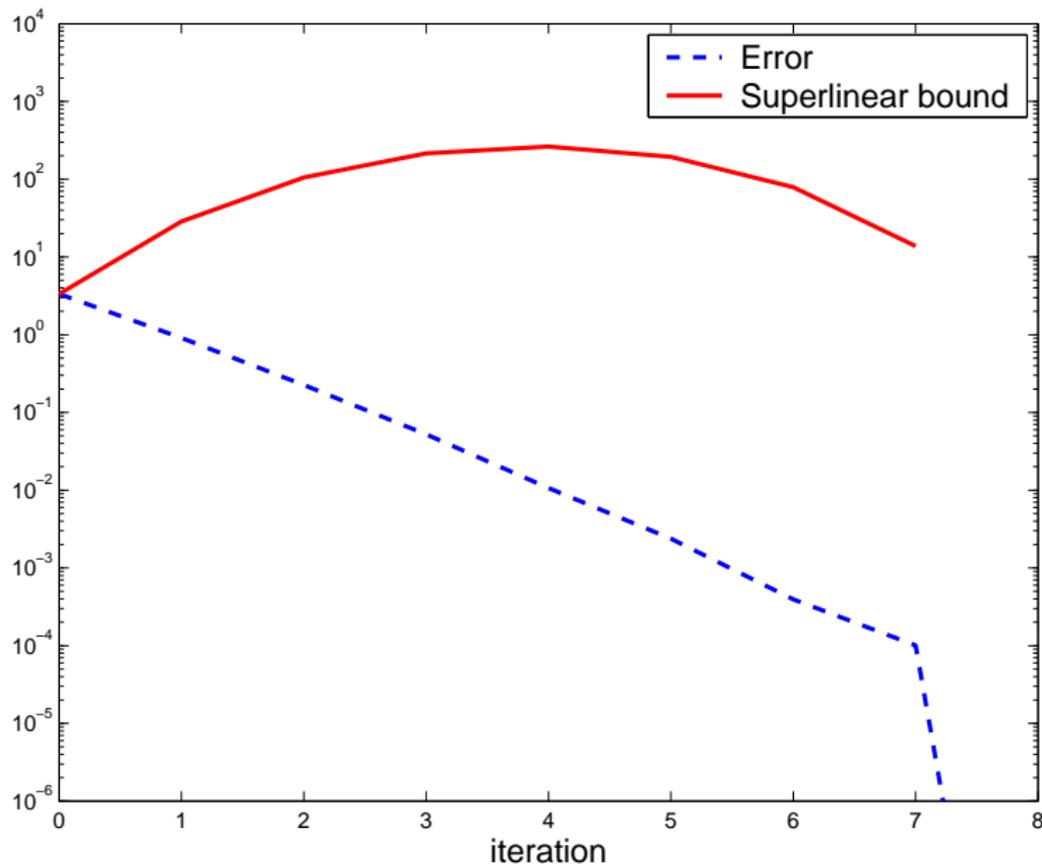












## A General Convergence Result

For the non-linear IVP  $u' = f(u)$ ,  $u(t_0) = u_0$ .

### Theorem (G, Hairer 2006)

Let  $F(t_{n+1}, t_n, U_n^k)$  denote the exact solution at  $t_{n+1}$  and  $G(t_{n+1}, t_n, U_n^k)$  be a one step method with local truncation error bounded by  $C_1 \Delta T^{p+1}$ . If

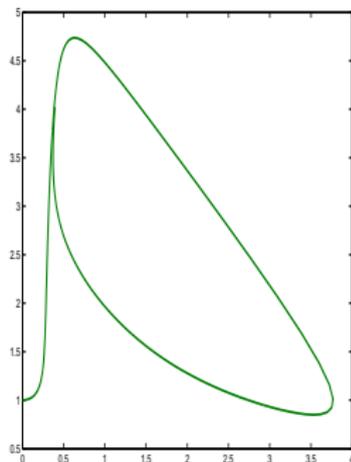
$$|G(t + \Delta T, t, x) - G(t + \Delta T, t, y)| \leq (1 + C_2 \Delta T) |x - y|,$$

then

$$\begin{aligned} \max_{1 \leq n \leq N} |u(t_n) - U_n^k| &\leq \frac{C_1 \Delta T^{k(p+1)}}{k!} (1 + C_2 \Delta T)^{N-1-k} \prod_{j=1}^k (N-j) \max_{1 \leq n \leq N} |u(t_n) - U_n^0| \\ &\leq \frac{(C_1 T)^k}{k!} e^{C_2(T-(k+1)\Delta T)} \Delta T^{pk} \max_{1 \leq n \leq N} |u(t_n) - U_n^0|. \end{aligned}$$

## Numerical experiments: Brusselator

$$\begin{aligned}\dot{x} &= A + x^2y - (B + 1)x \\ \dot{y} &= Bx - x^2y,\end{aligned}$$

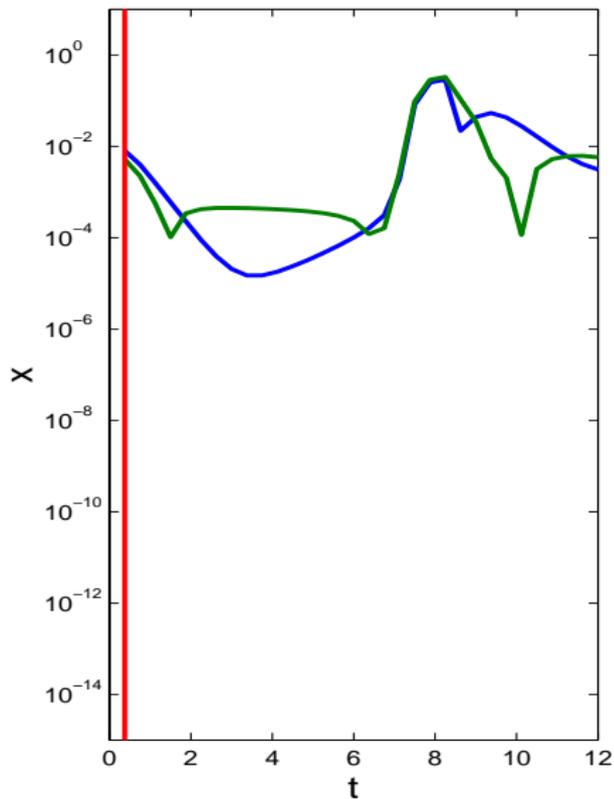
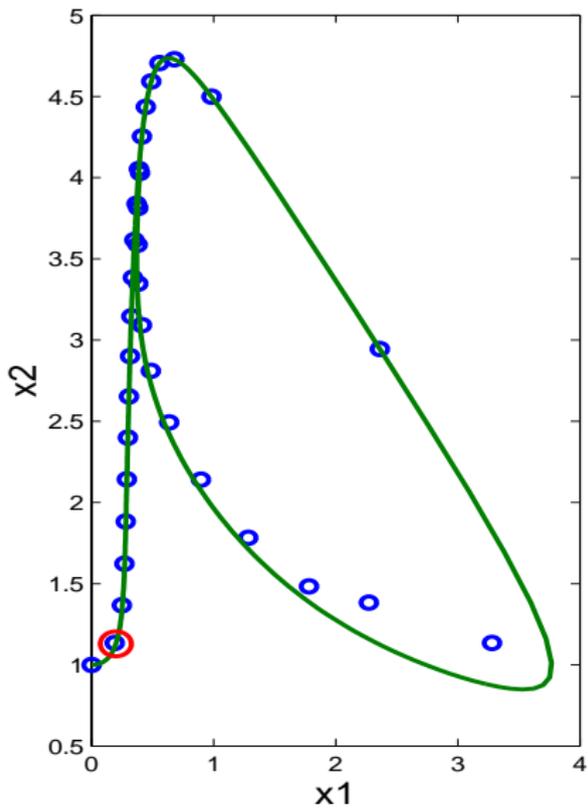


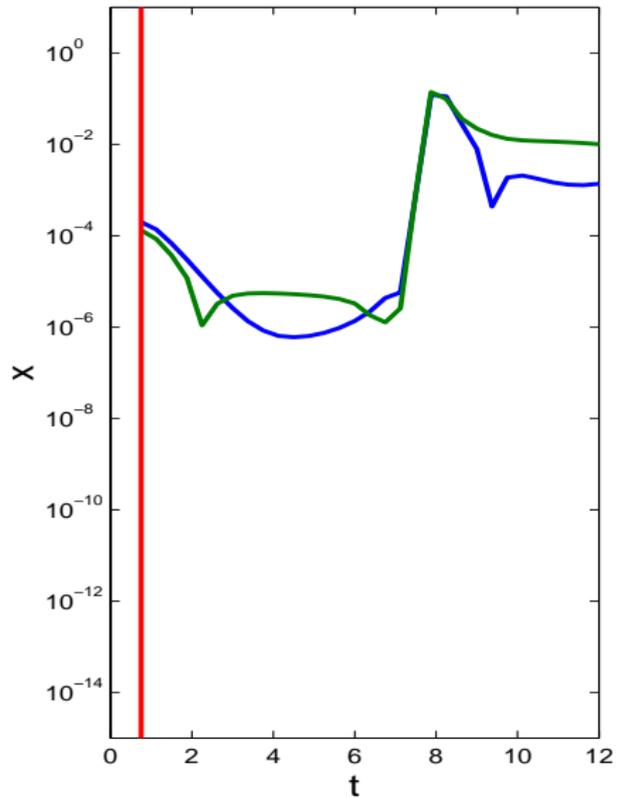
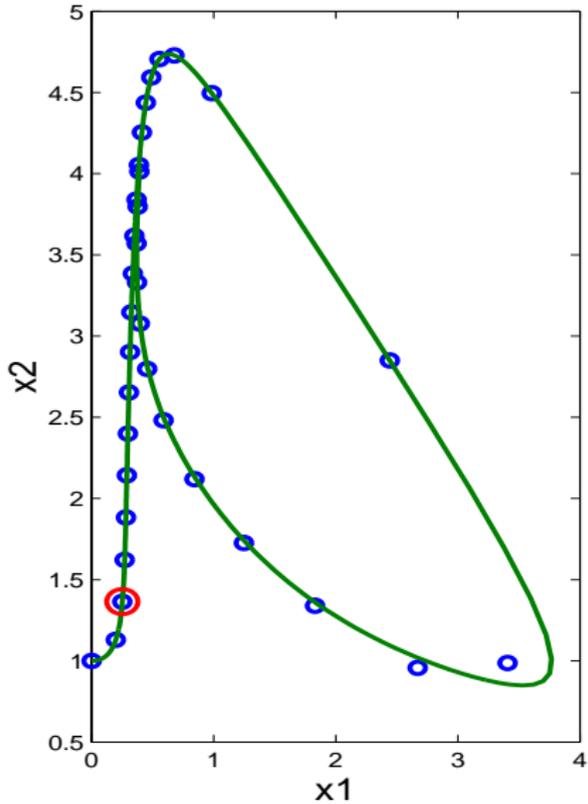
Parameters:  $A = 1$  and  $B = 3$ ,  $B > A^2 + 1 \implies$  limit cycle.

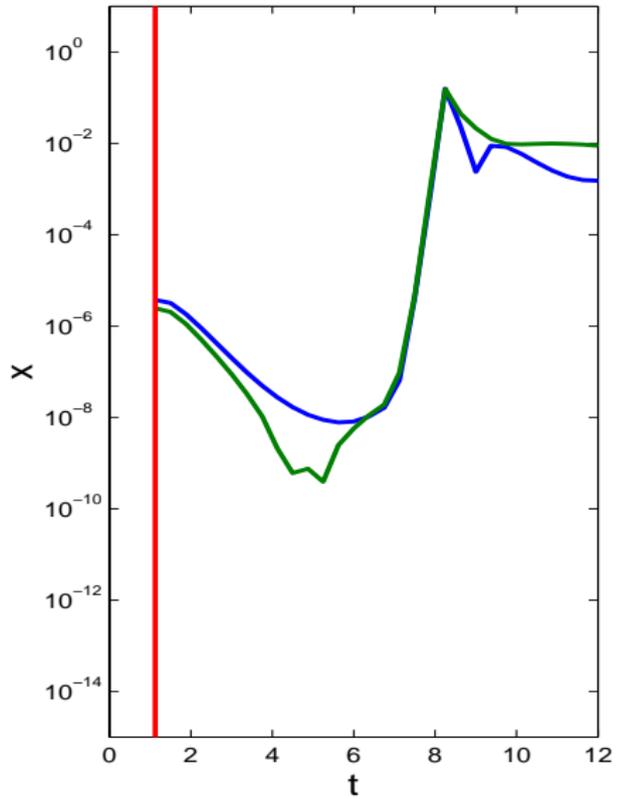
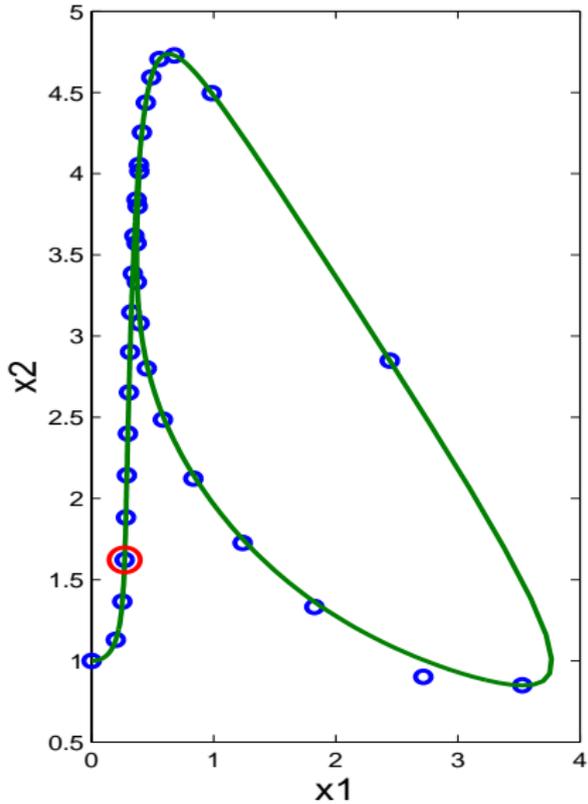
Initial conditions:  $x(0) = 0$ ,  $y(0) = 1$ .

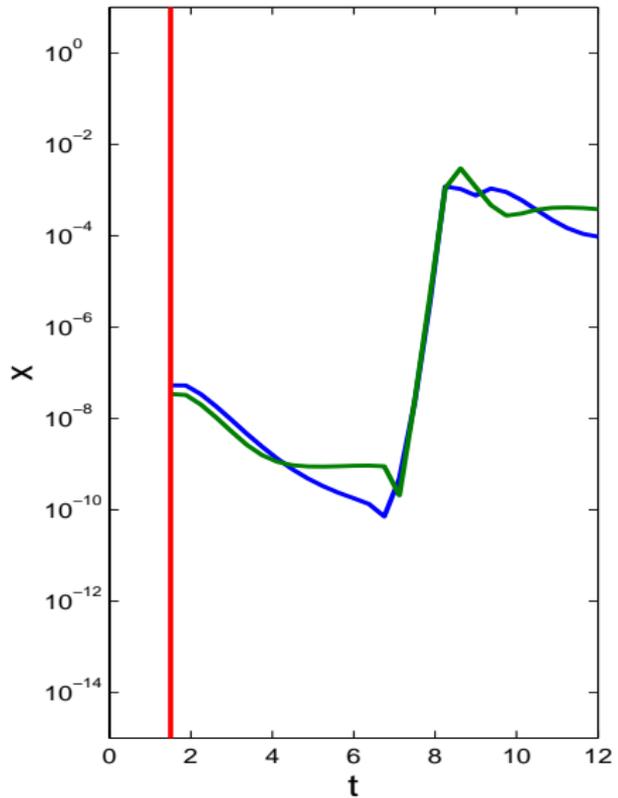
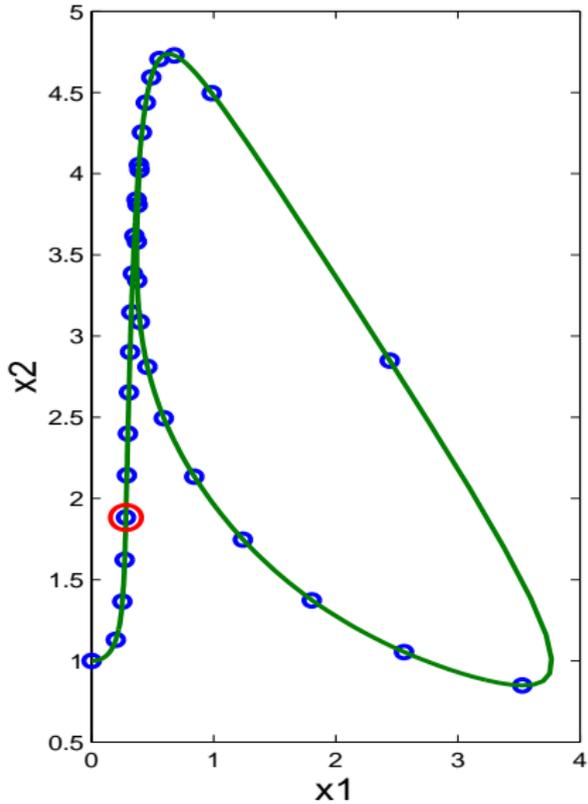
Simulation time:  $t \in [0, T = 12]$

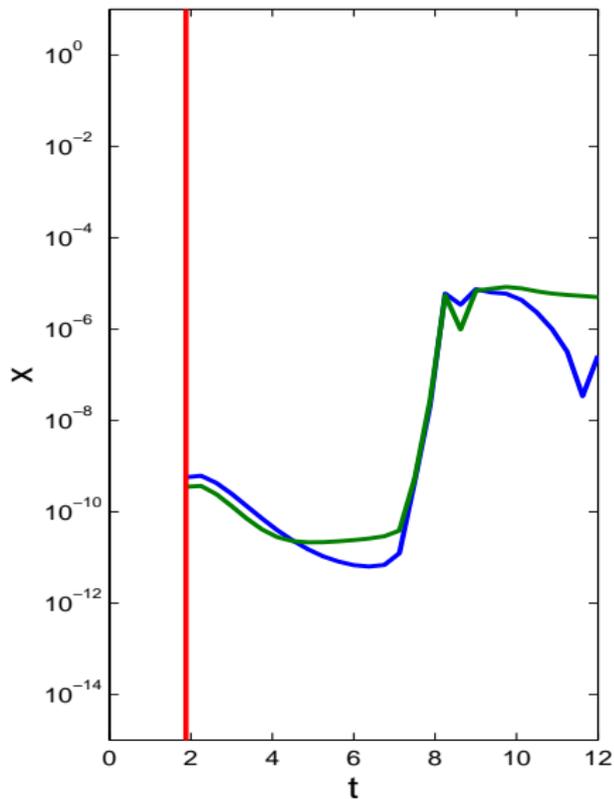
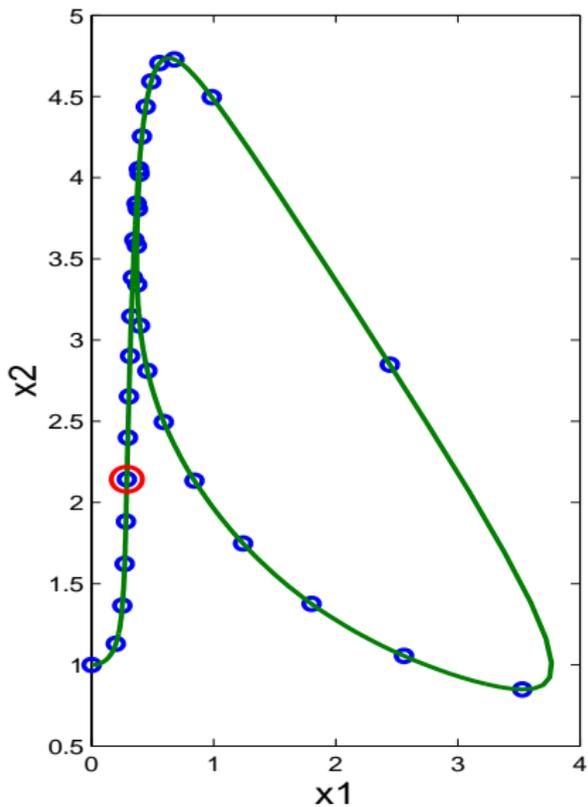
Discretization: Fourth order Runge Kutta,  $\Delta T = \frac{T}{32}$ ,  $\Delta t = \frac{T}{320}$ .

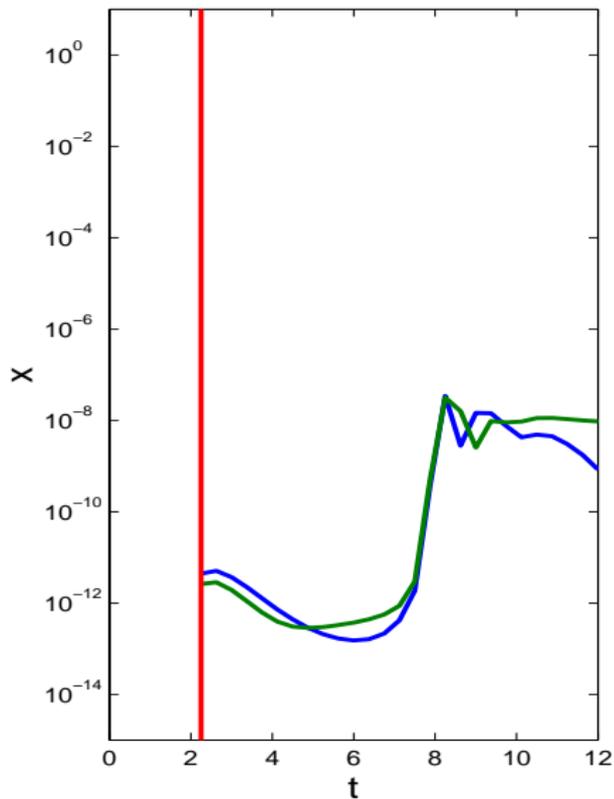
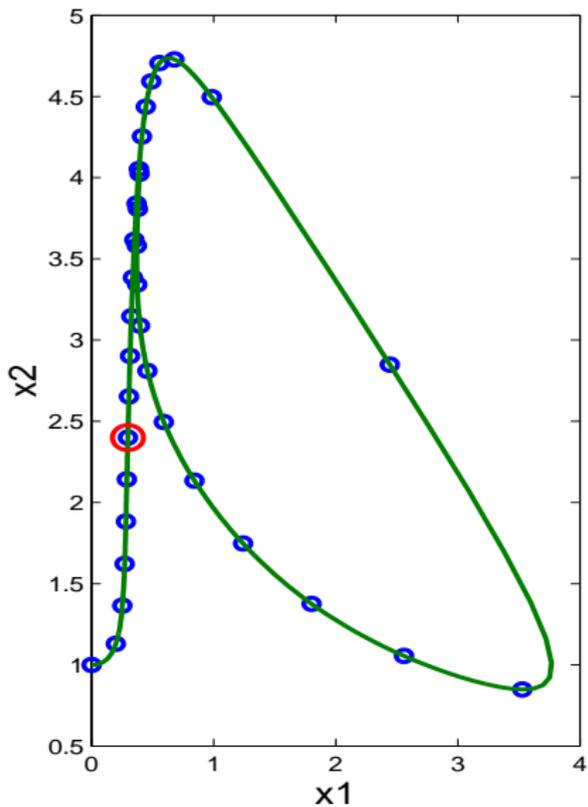


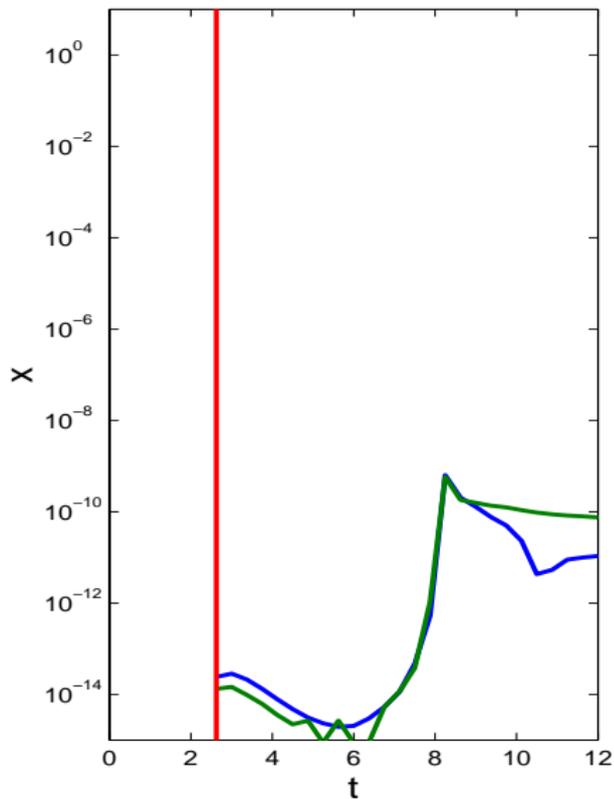
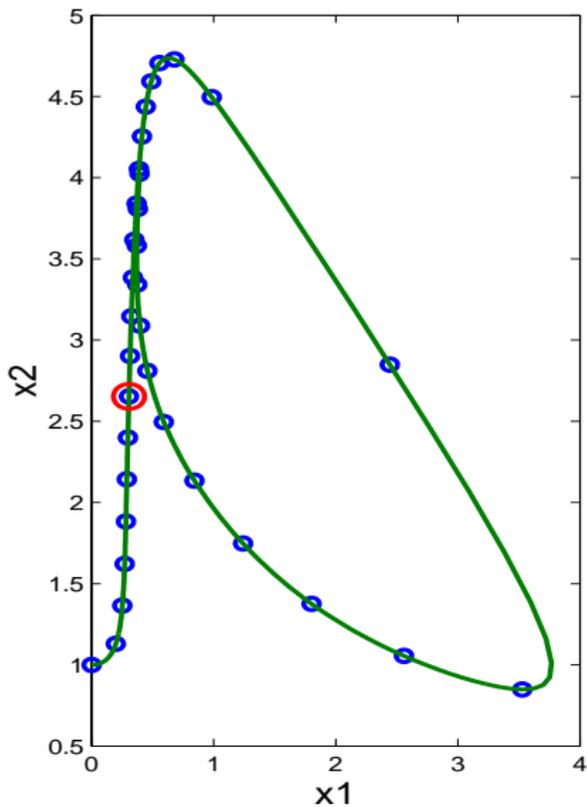


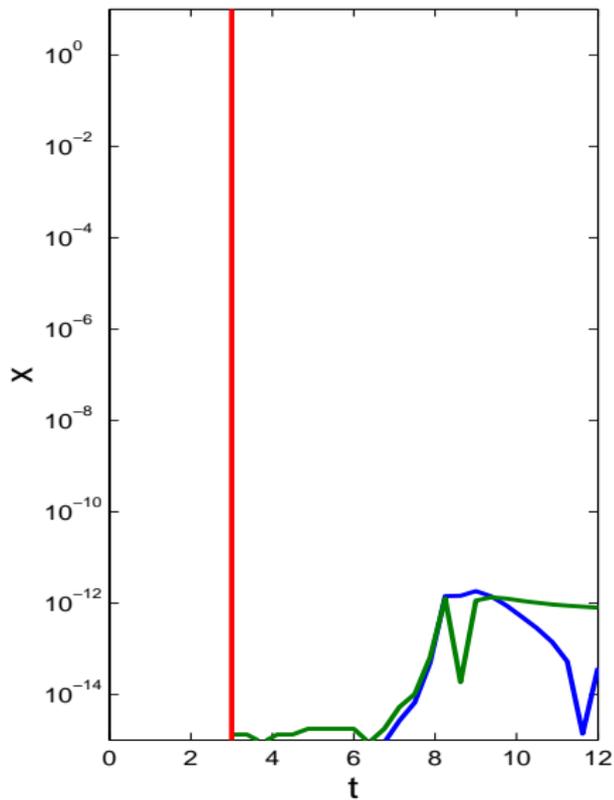
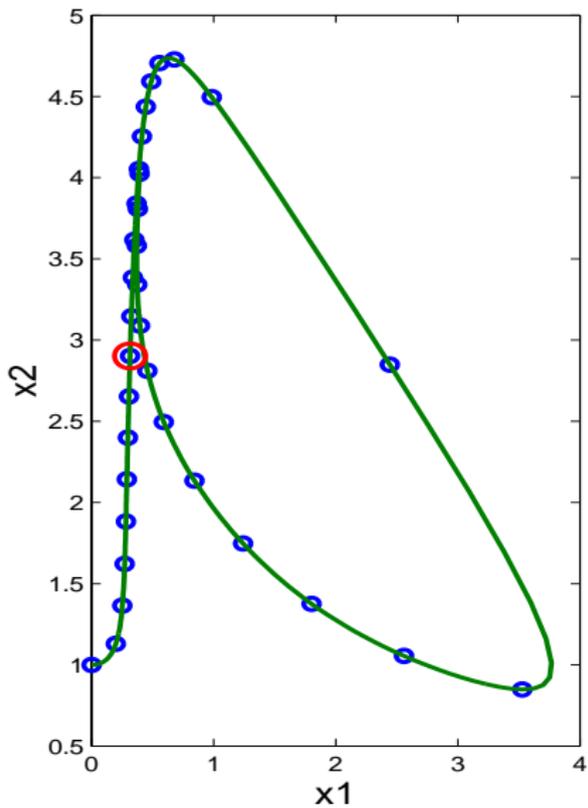










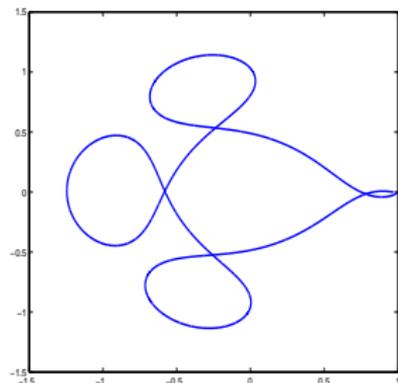


## Numerical experiments: Arenstorf orbit

$$\ddot{x} = x + 2\dot{y} - b \frac{x+a}{D_1} - a \frac{x-b}{D_2}$$

$$\ddot{y} = y - 2\dot{x} - b \frac{y}{D_1} - a \frac{y}{D_2},$$

$$D_1 = ((x+a)^2 + y^2)^{(3/2)}, \quad D_2 = ((x-b)^2 + y^2)^{(3/2)}$$



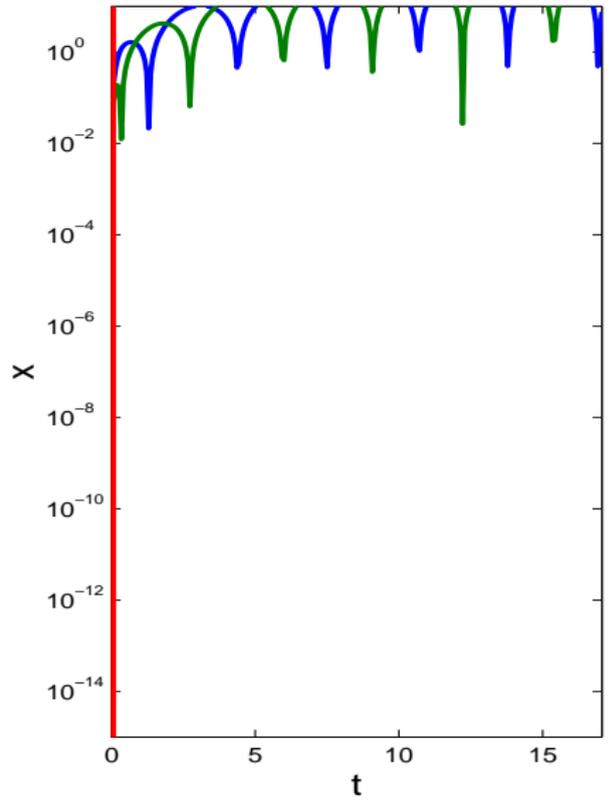
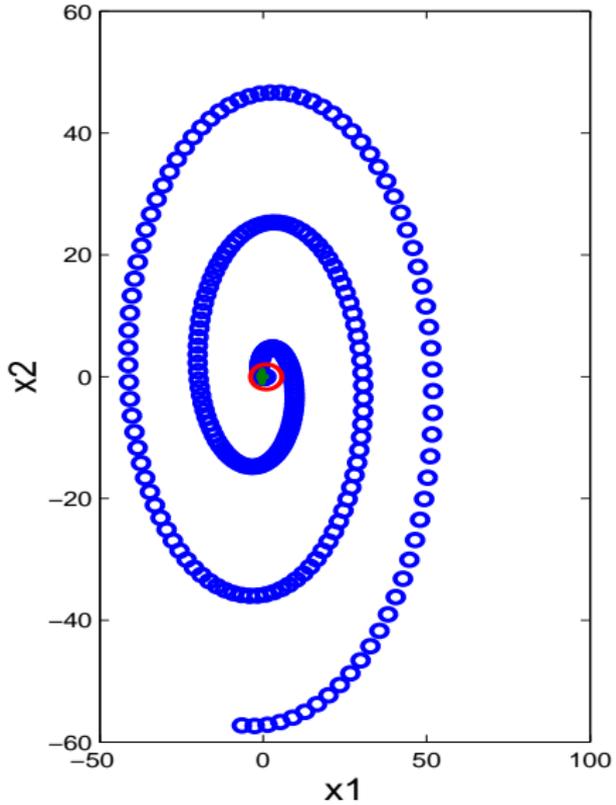
Parameters:  $a = 0.012277471$ ,  $b = 1 - a$ .

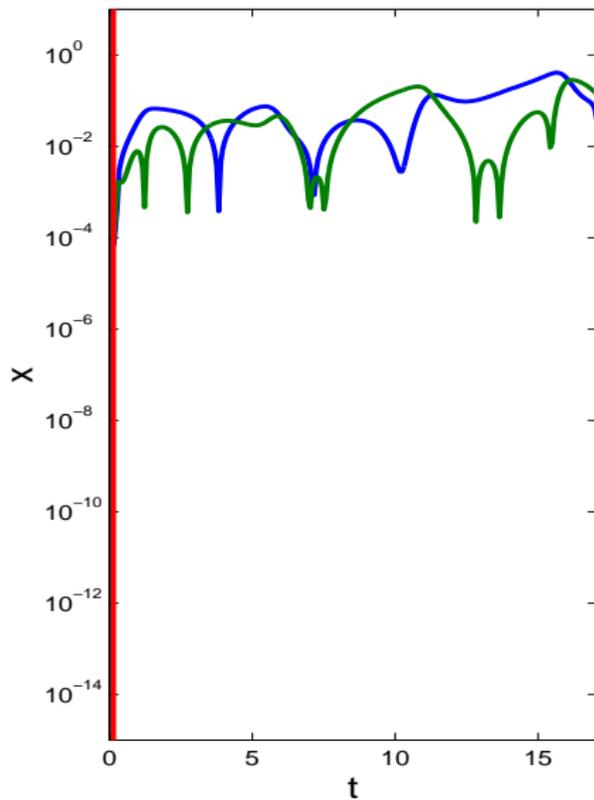
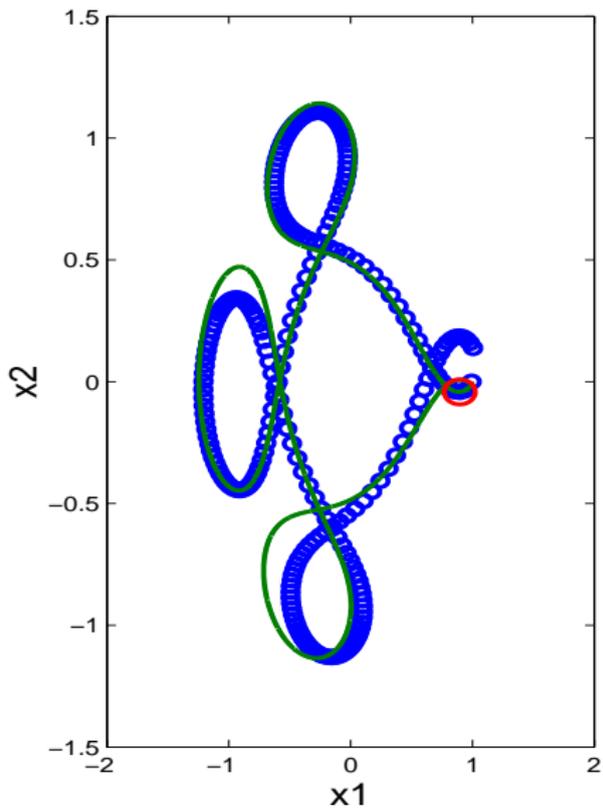
Initial conditions:  $x(0) = 0.994$ ,  $\dot{x} = 0$ ,  
 $y(0) = 0$ ,  $\dot{y}(0) = -2.00158510637908$

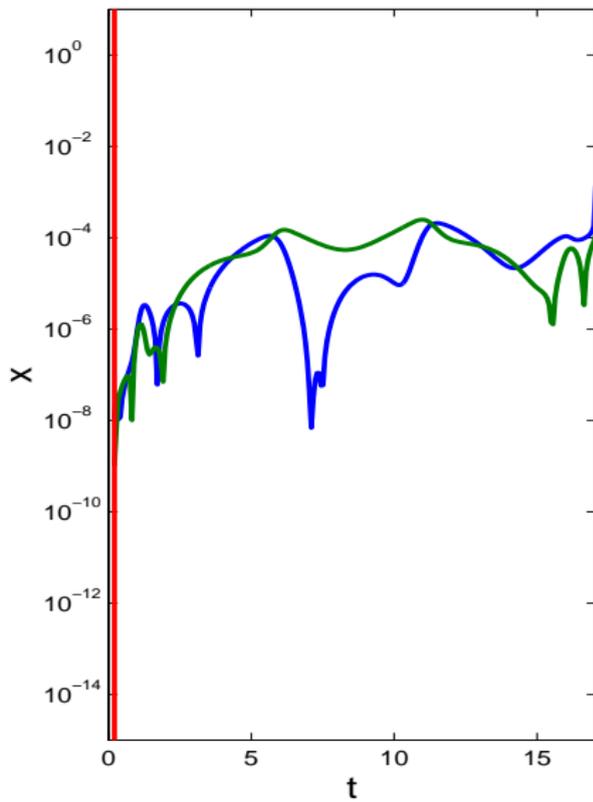
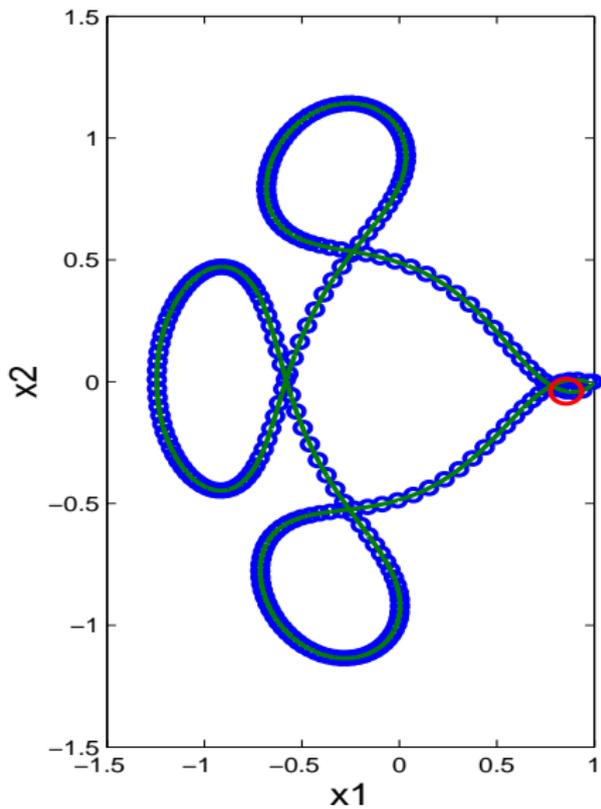
Simulation time:  $t \in [0, T = 17.06]$

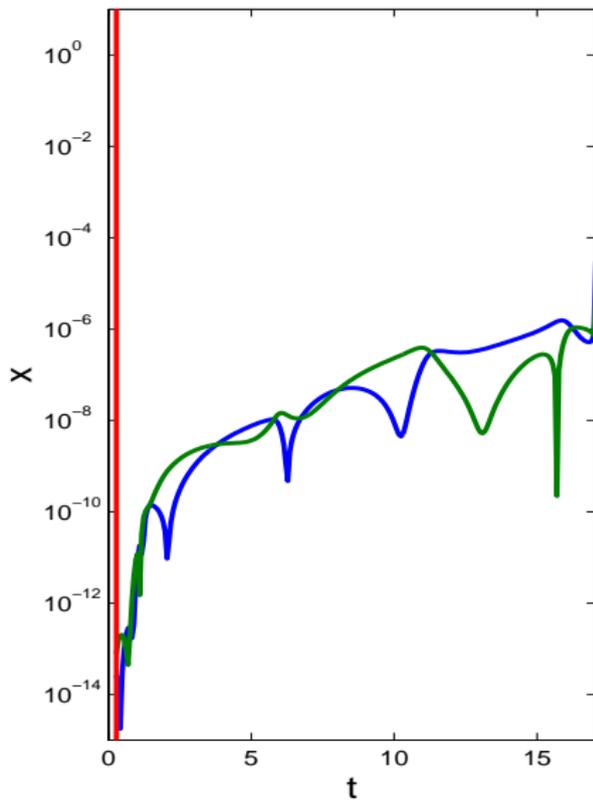
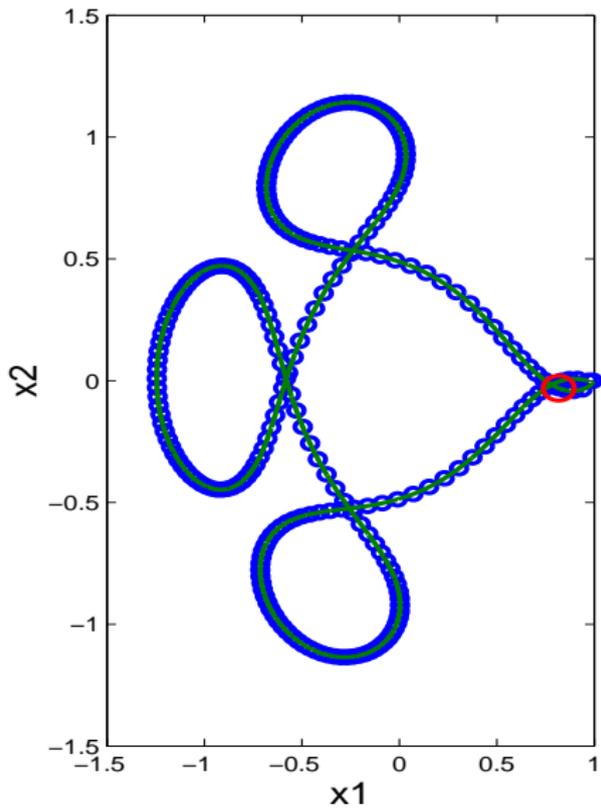
Discretization: Fourth order Runge Kutta,  $\Delta T = \frac{T}{250}$ ,  $\Delta t = \frac{T}{10000}$ .

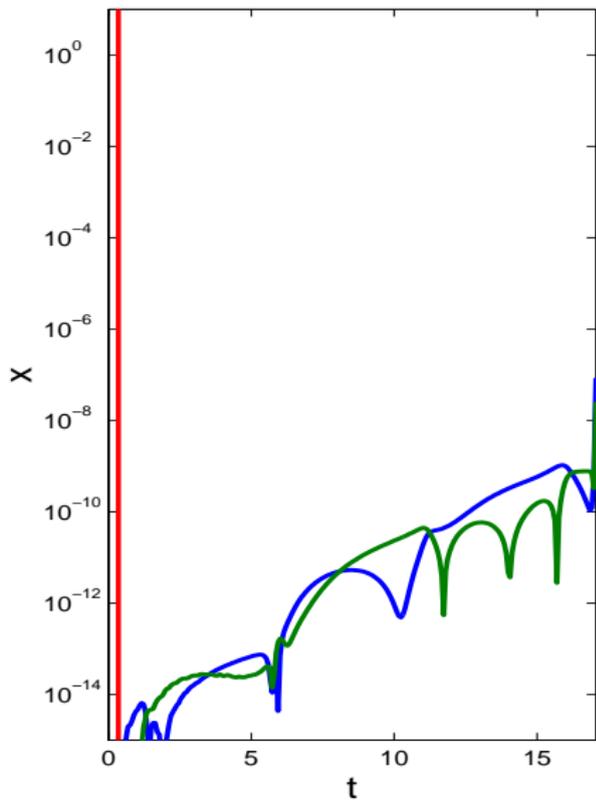
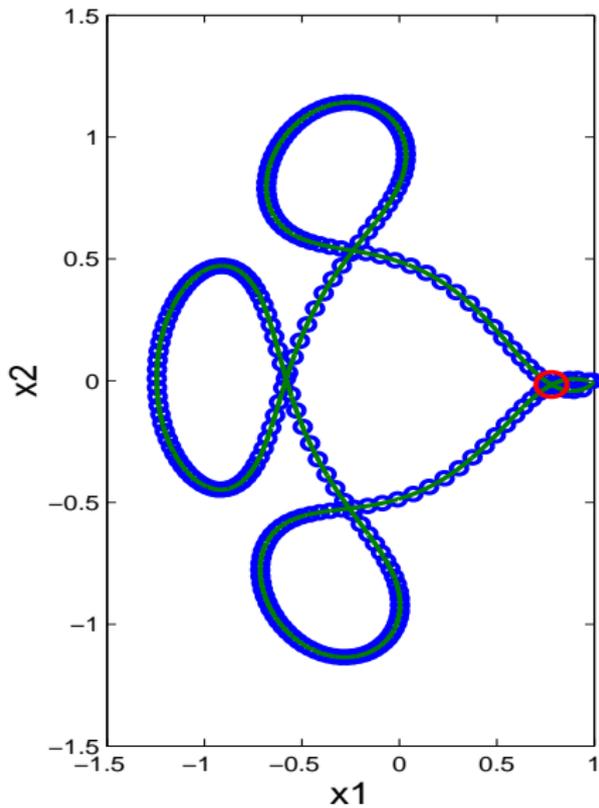
See also Saha, Stadel and Tremaine, a parallel integration method for solar system dynamics, 1997

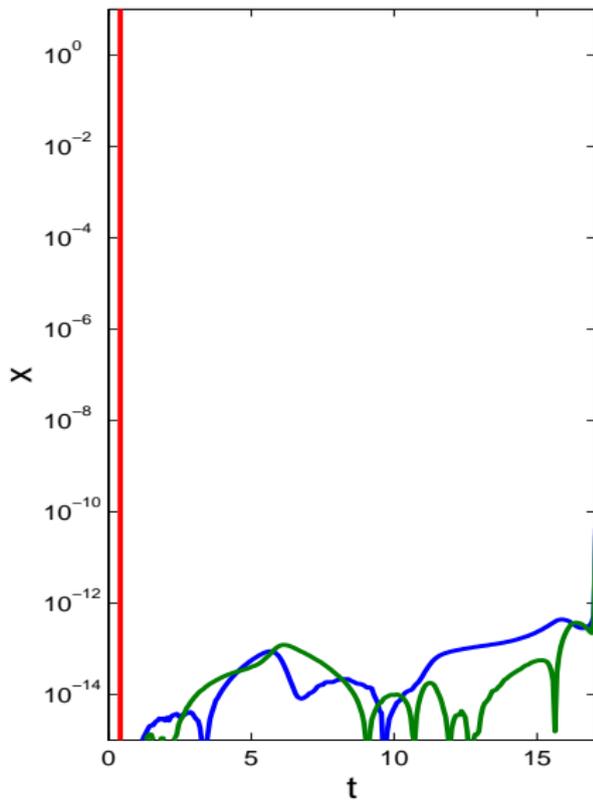
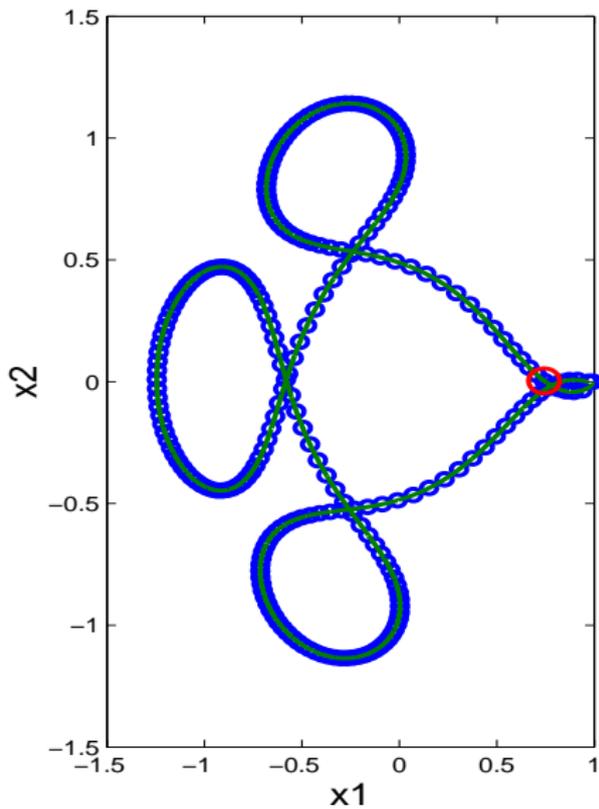






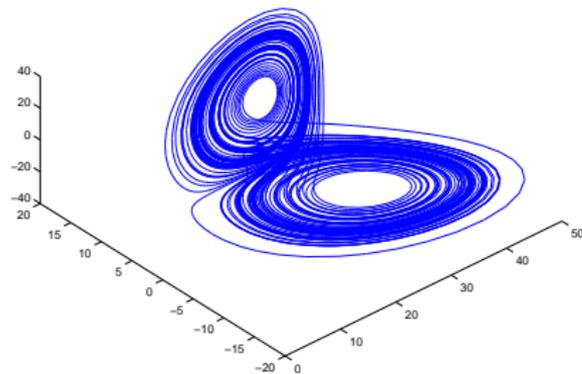






## Results for the Lorenz Equations

$$\begin{aligned}\dot{x} &= -\sigma x + \sigma y \\ \dot{y} &= -xz + rx - y \\ \dot{z} &= xy - bz\end{aligned}$$

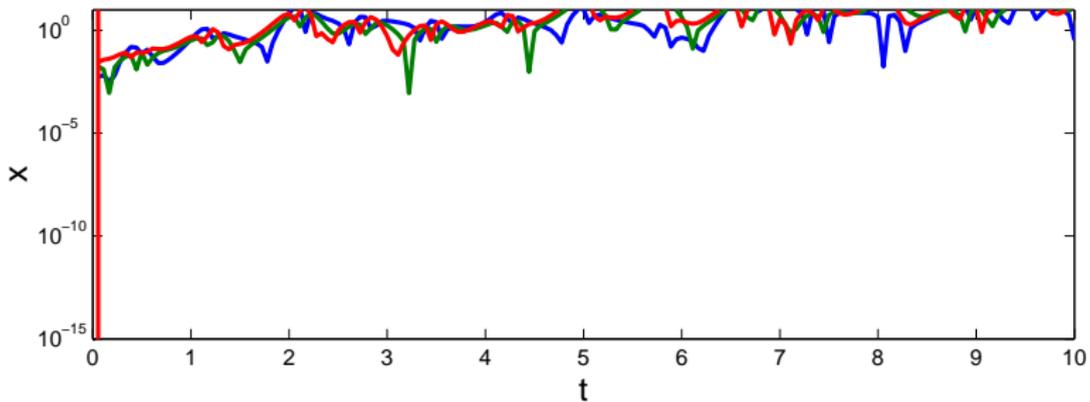
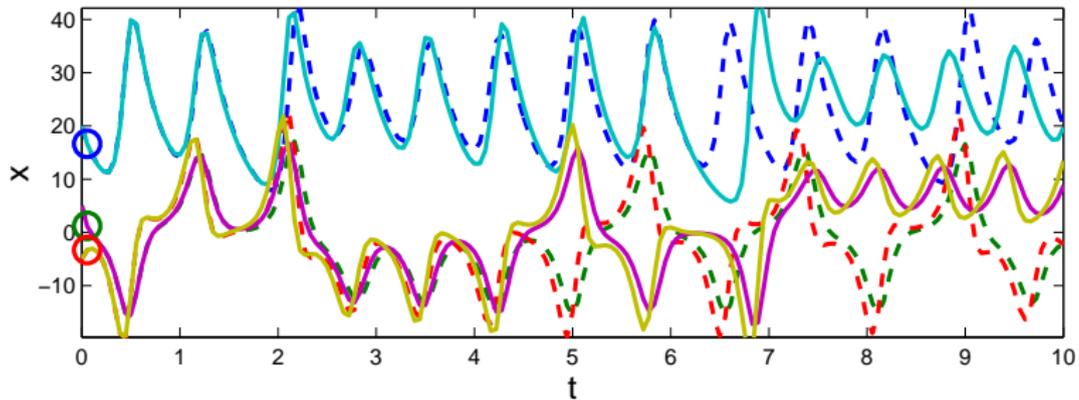


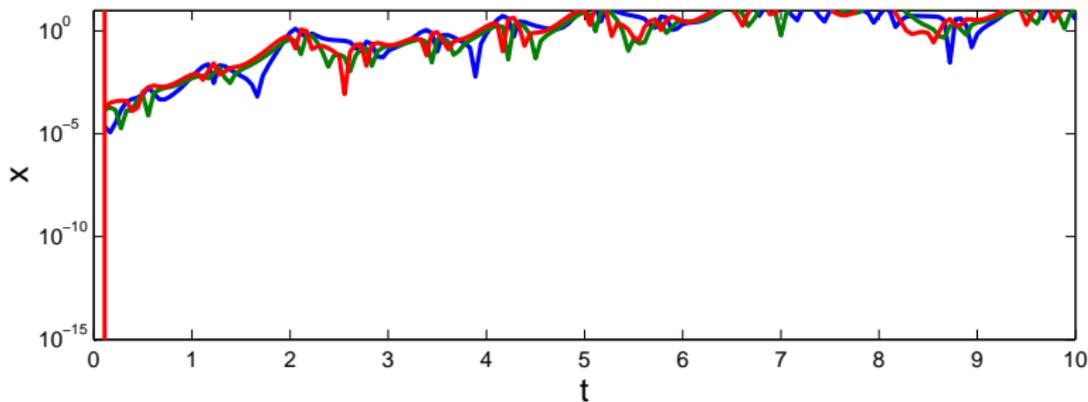
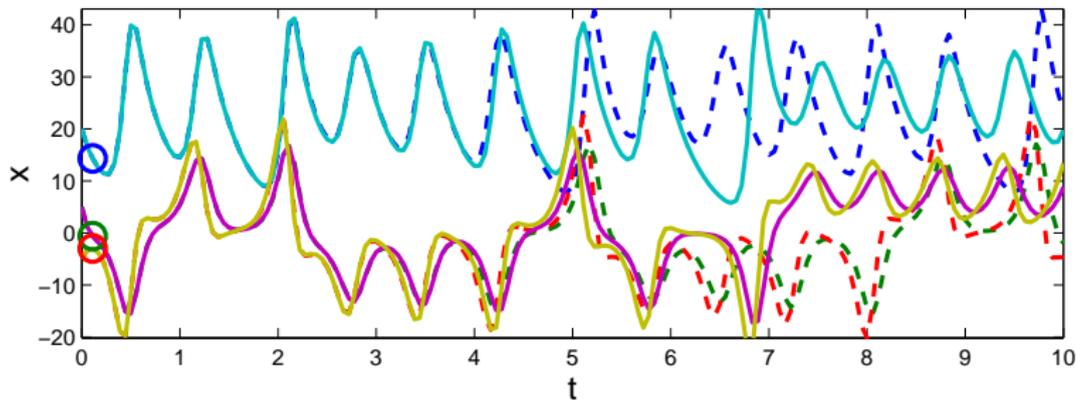
Parameters:  $\sigma = 10$ ,  $r = 28$  and  $b = \frac{8}{3} \implies$  chaotic regime.

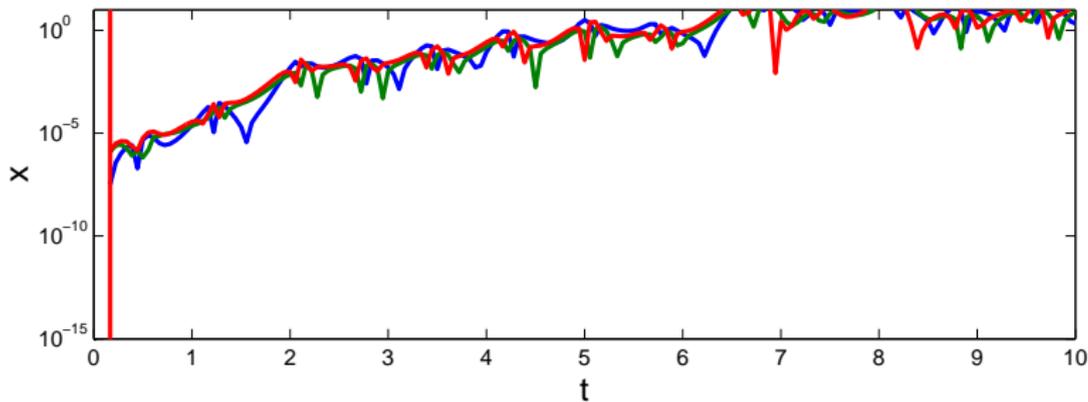
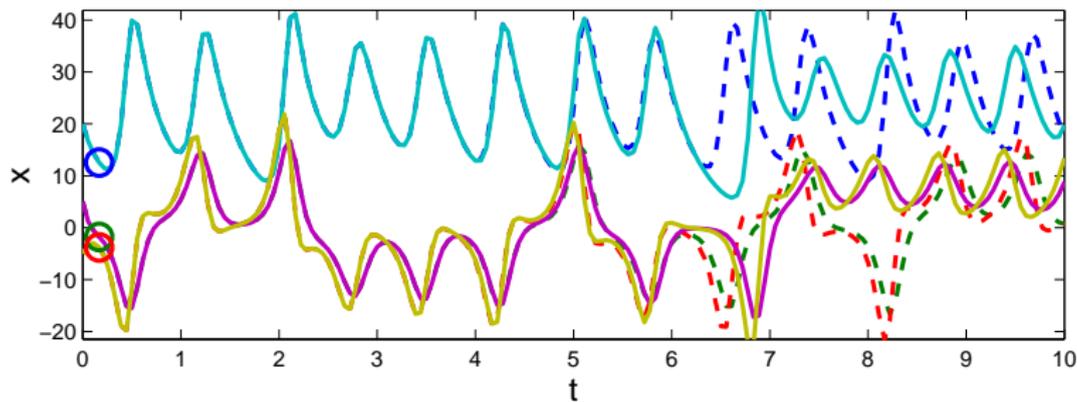
Initial conditions:  $(x, y, z)(0) = (20, 5, -5)$

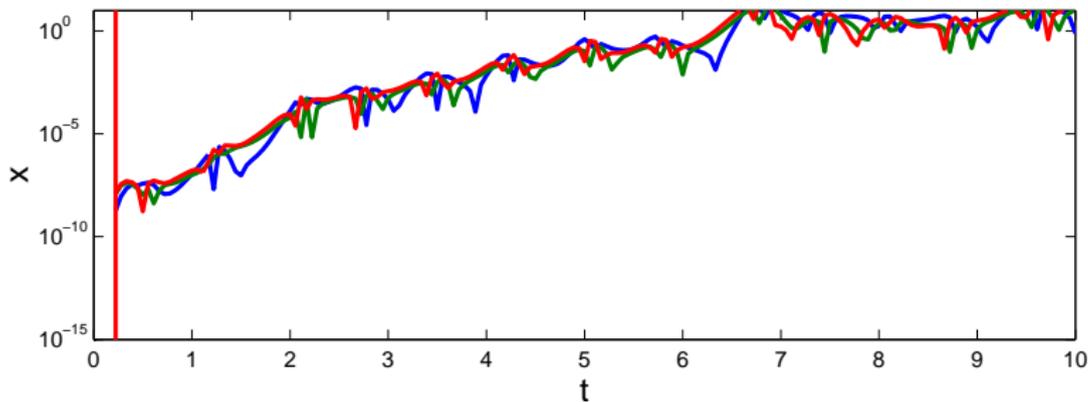
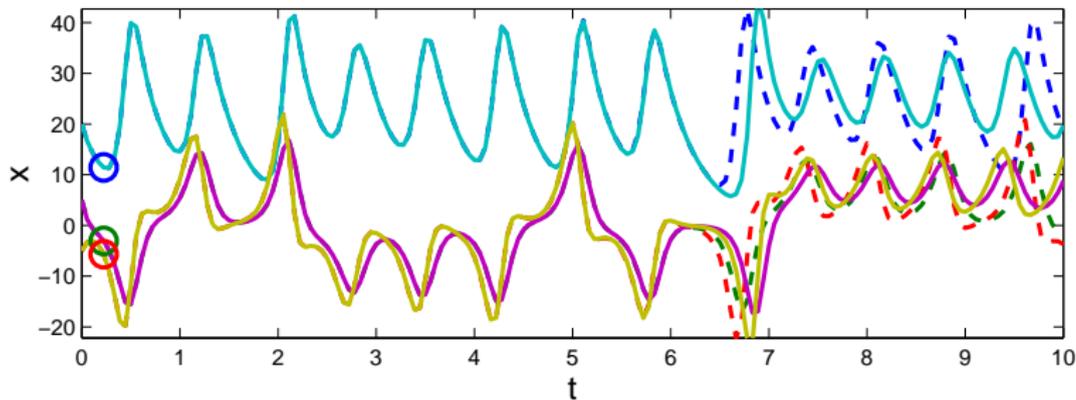
Simulation time:  $t \in [0, T = 10]$

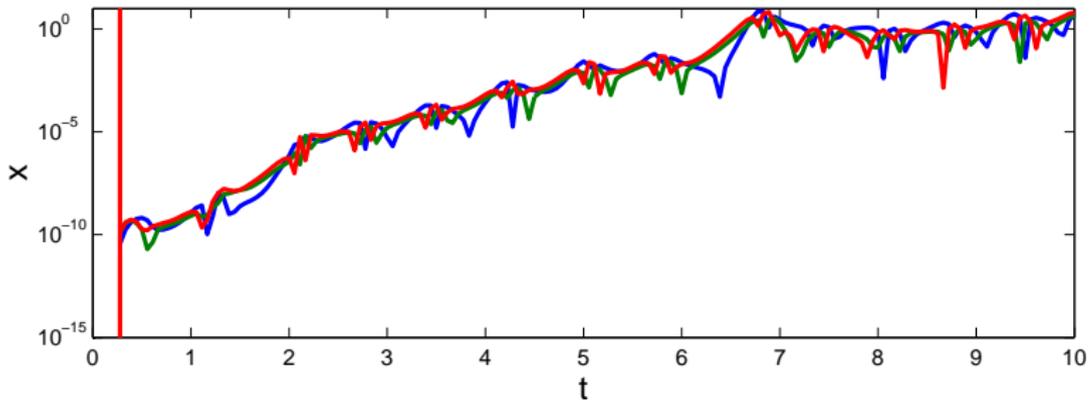
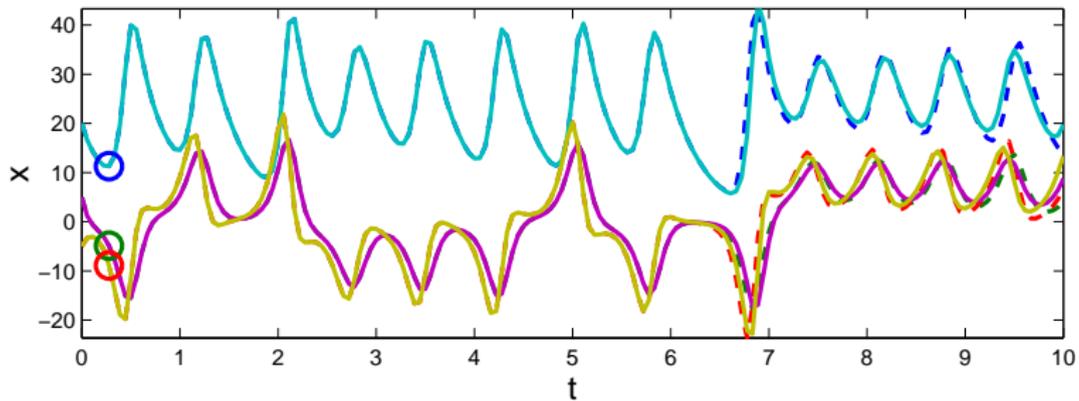
Discretization: Fourth order Runge Kutta,  $\Delta T = \frac{T}{180}$ ,  $\Delta t = \frac{T}{1800}$ .

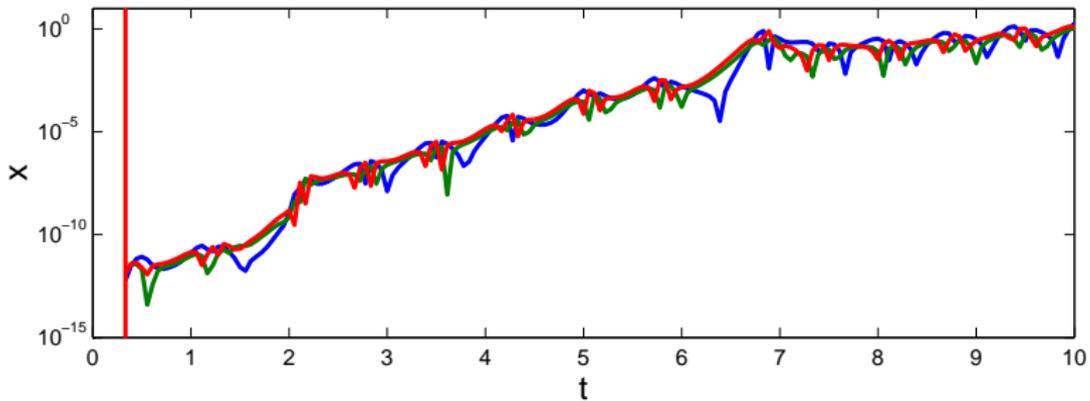
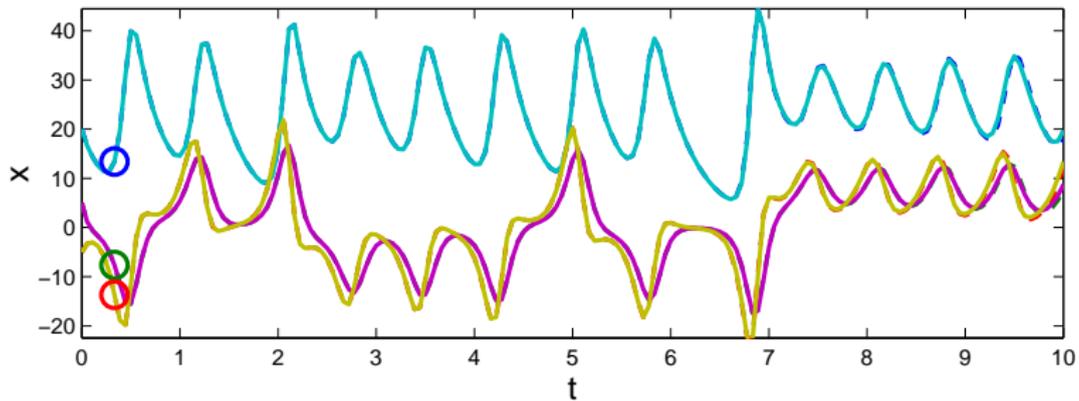


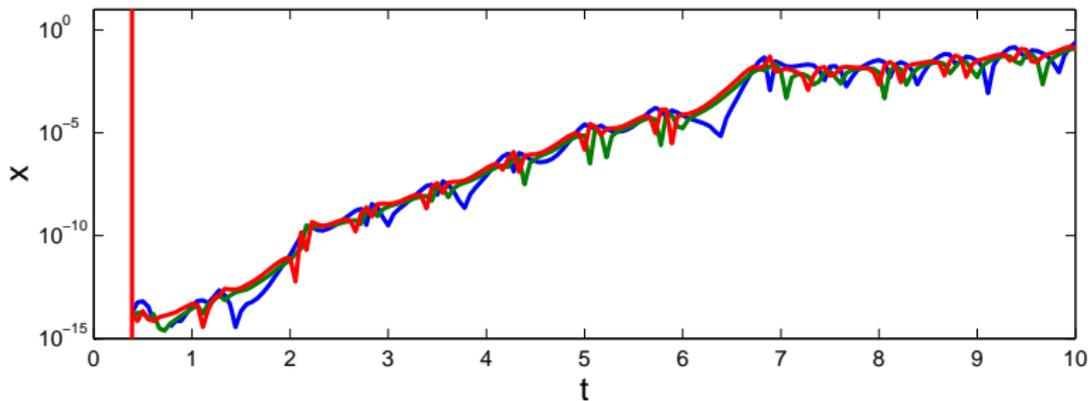
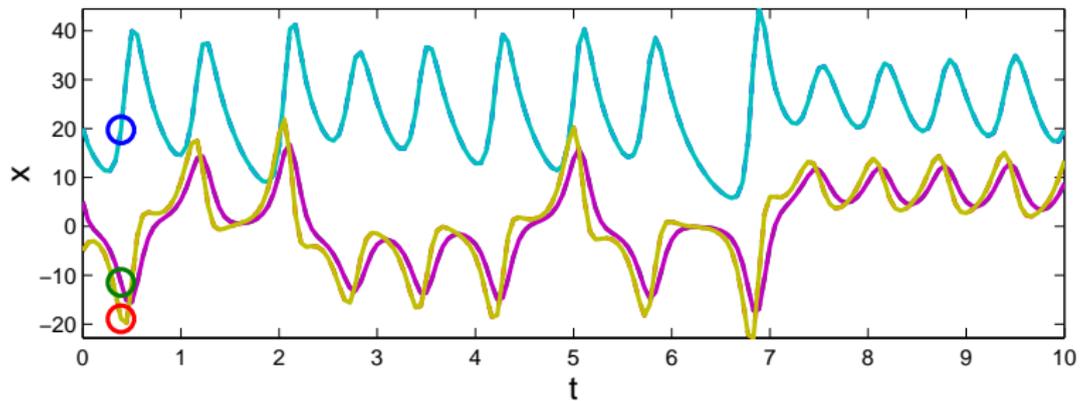


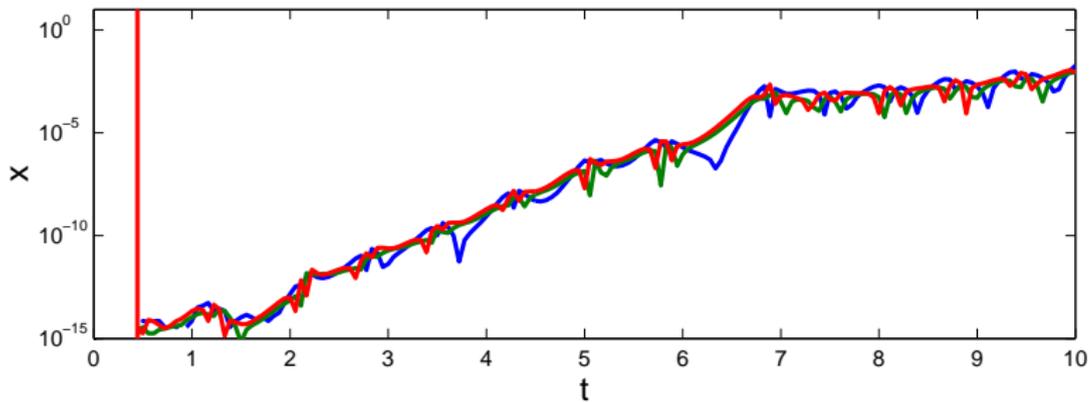
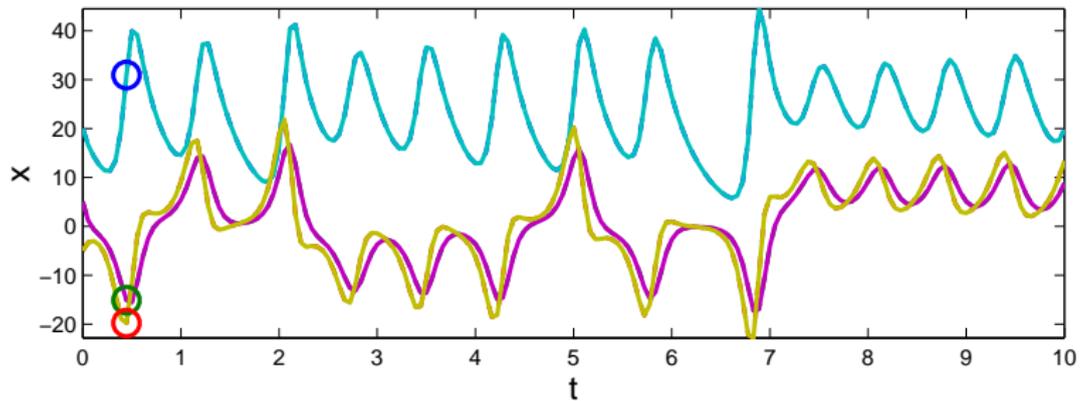


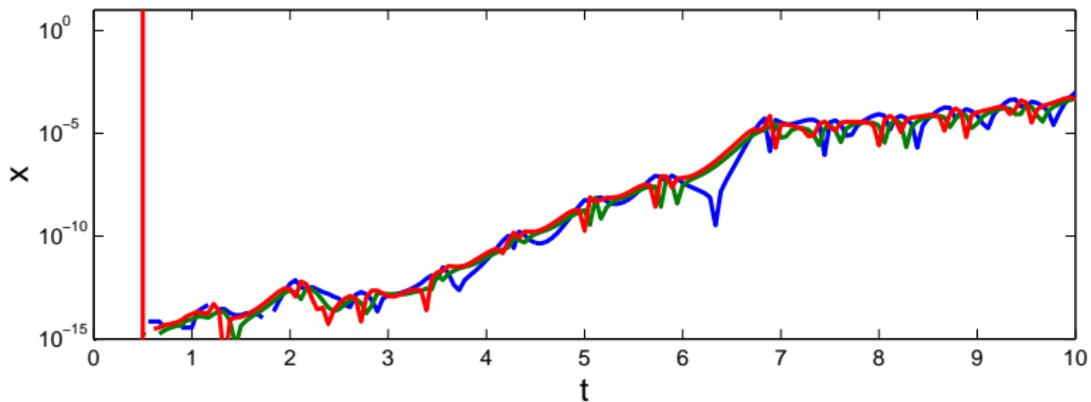
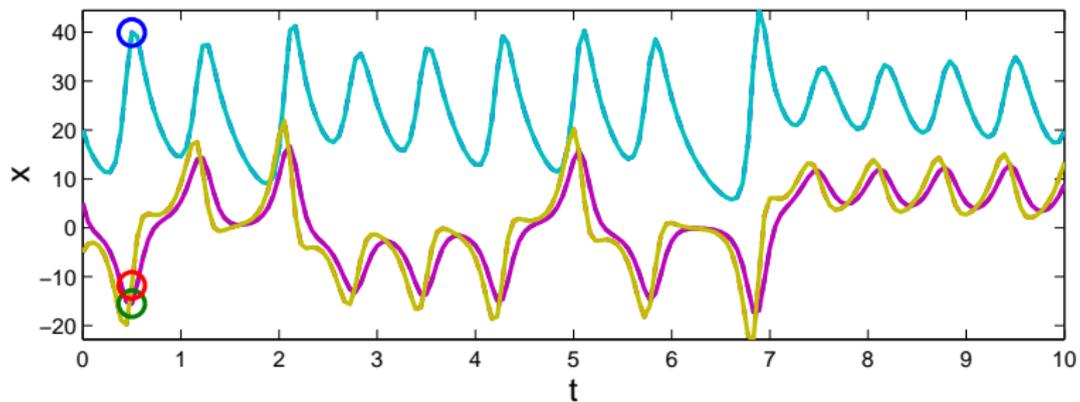


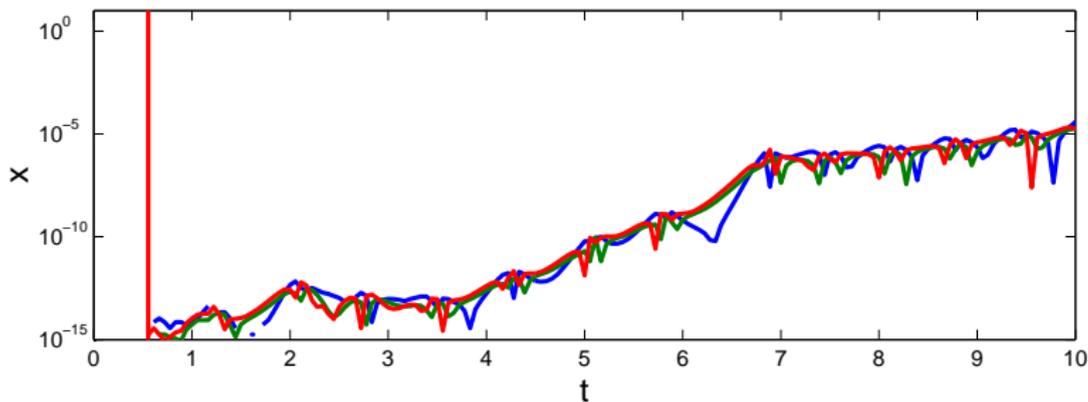
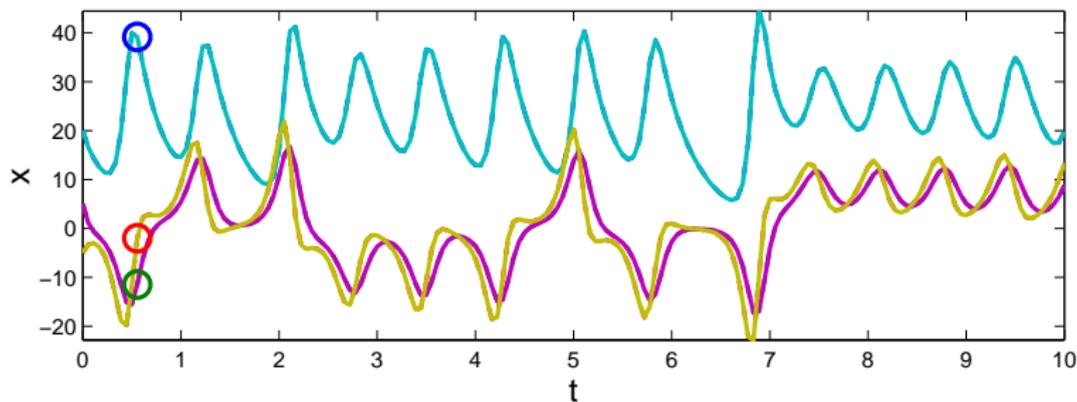


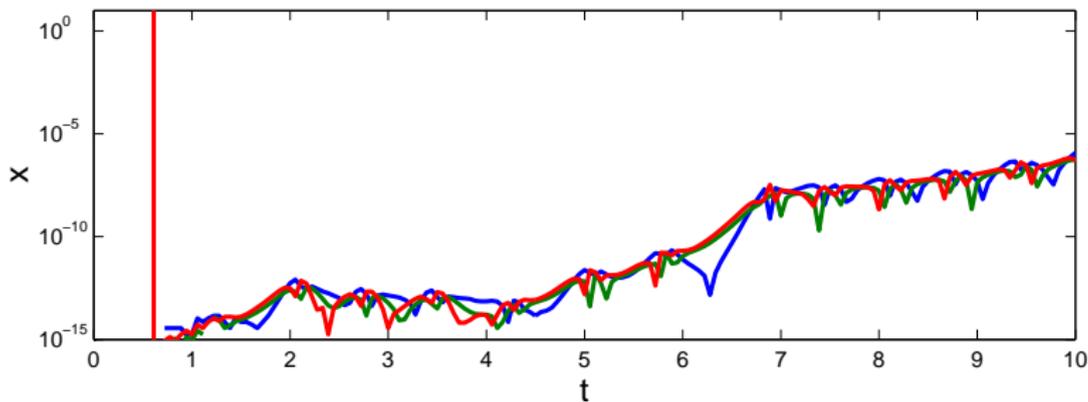
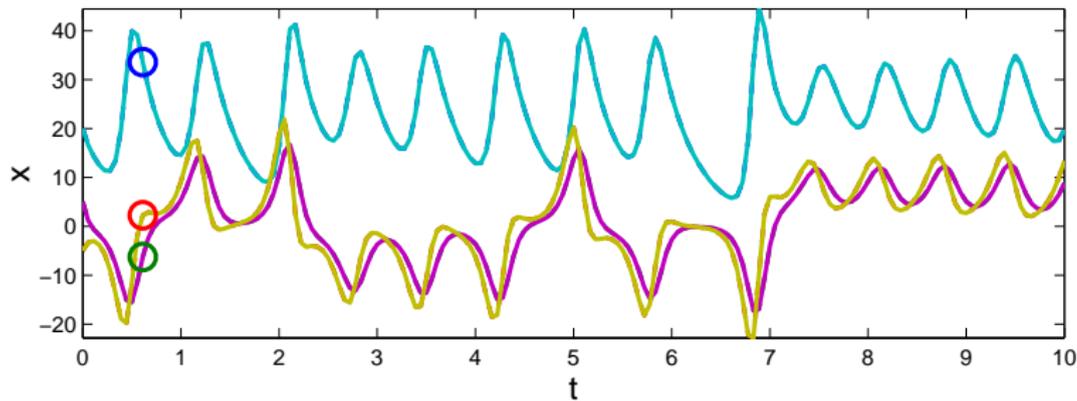


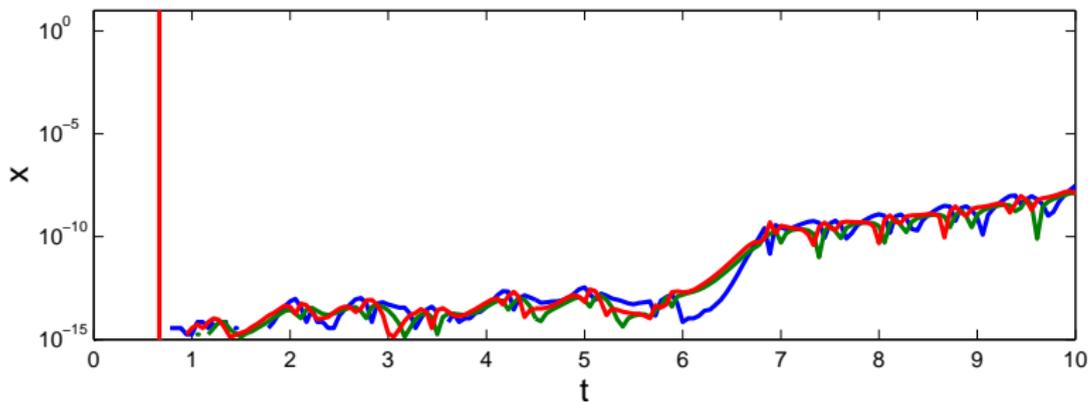
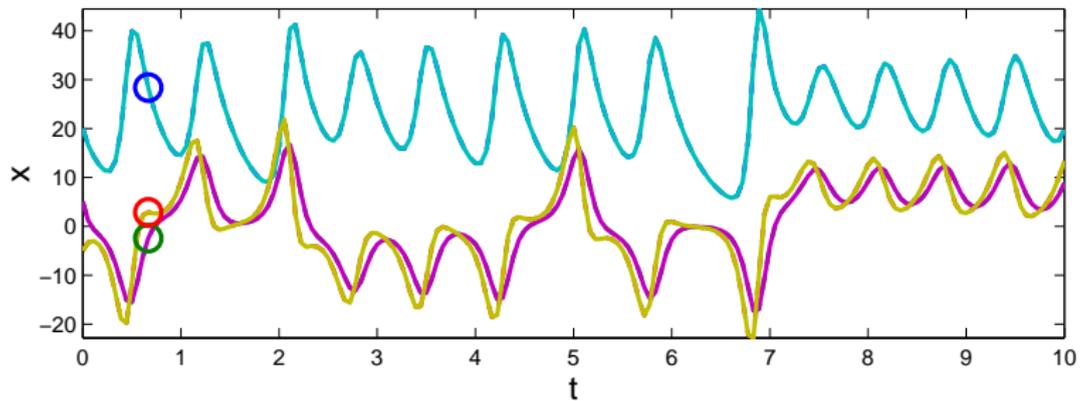


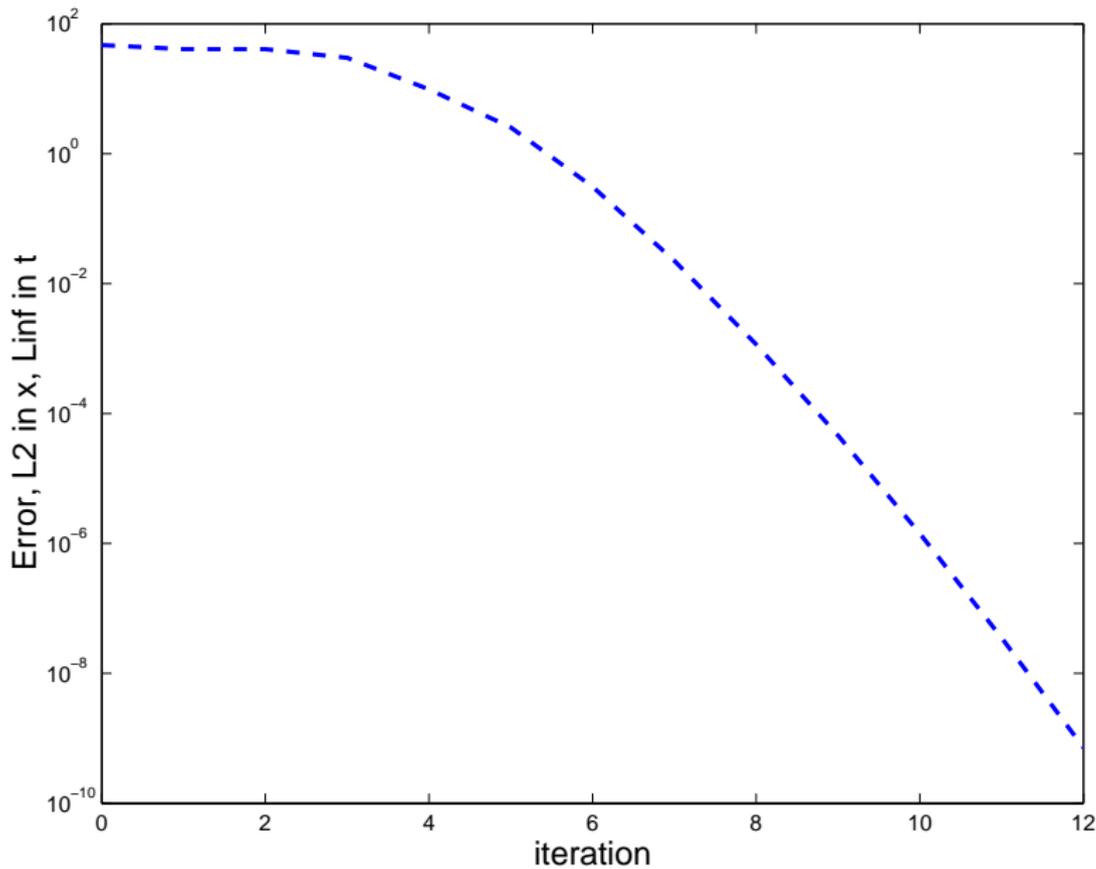












## Numerical experiments for PDEs: Burgers equation

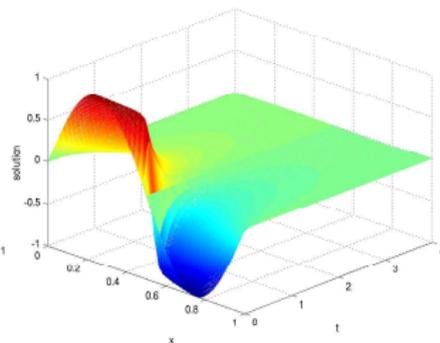
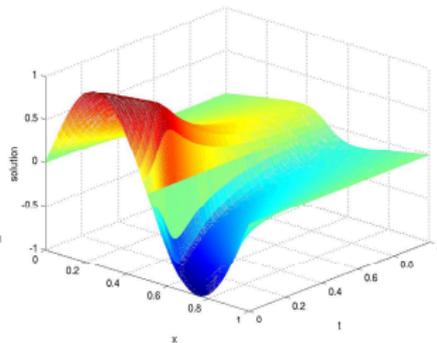
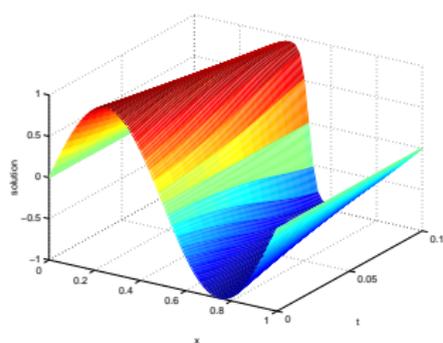
$$u_t + uu_x = \nu u_{xx} \quad \text{in } \Omega = [0, 1]$$

$$u(x, 0) = \sin(2\pi x)$$

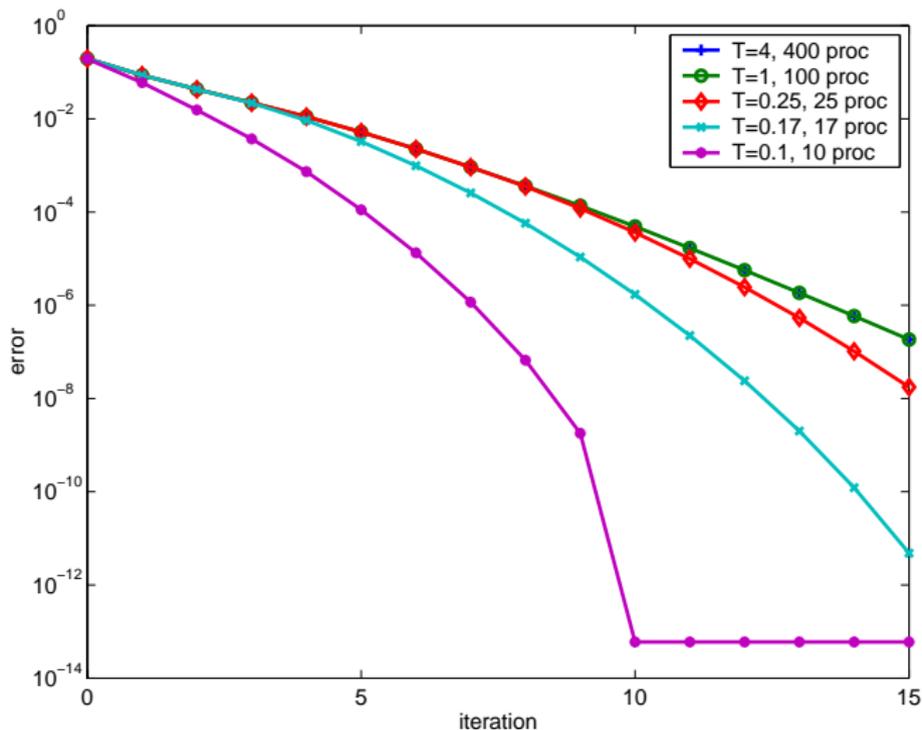
Viscosity  $\nu = \frac{1}{50}$ , homogeneous boundary conditions

Centered finite difference discretization,  $\Delta x = \frac{1}{50}$

Backward Euler in time,  $\Delta T = \frac{1}{10}$ ,  $\Delta t = \frac{1}{100}$ .



# Burgers equation: convergence behavior



## Further Applications and Results on Parareal

- ▶ **Oscillatory Problems:** Cortial, Farhat, Chandesris (2003, 2006)
- ▶ **Control of Quantum Systems:** Maday, Salomon, Turinici (2002, 2006)
- ▶ **Reservoir Simulation:** Garrido, Espedal, Fladmark (2003, 2005)
- ▶ **Navier-Stokes:** Fischer, Hecht, Maday (2003)
- ▶ **Stability Analysis:** Staff and Rønquist (2003)
- ▶ **Molecular Dynamics:** Baffico, Bernard, Maday, Turinici, Zerah (2002)
- ▶ **Finance:** Bal, Maday (2002)

**Google hits for parareal algorithm (6.5.2007): 433**

# Conclusions

Parallel speedup in time is possible, but the speedup is more modest than in space.

## Further results:

- ▶ Two multilevel versions of the algorithm.
- ▶ Understand hyperbolic case with boundary conditions, and extension by Farhat for general hyperbolic problems.

## Future work:

- ▶ Analysis of Parareal for DAEs.
- ▶ Preservation of symplectic structure in Parareal.