Nonequilibrium dynamics on complex topologies: models for epidemics and opinions

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Outline

- Introduction
- Models for epidemics: SIS and SIR
- Disordered contact process: rare region effects
- Opinion dynamics: voter model on networks
- Conclusions

Introduction

- Statistical physics approach to interdisciplinary research
- Complex topologies are the natural substrate
- Highly nontrivial interplay between structure and dynamics

Susceptible-Infected-Susceptible (SIS) model

- Two possible states: susceptible and infected
- Two possible events for infected nodes:
- Recovery (rate 1) Infection to neighbors (rate λ) 0.4 0.3 0.2 ho0.1 0.0 0.2 0.4 0.6 0.0 λ

Heterogeneous Mean-Field theory for SIS Pastor-Satorras and Vespignani (2001)

- Degree distribution $P(k) \sim k^{-\gamma}$
- ρ_k = density of infected nodes of degree k



Susceptible-Infected-Removed (SIR) model

- Three possible states: susceptible, infected and removed.
- Two possible events for infected nodes:
- Death/recovery (rate 1)
- Infection to neighbors (rate λ)



Transition between healthy and infected

HMF for SIR

• HMF theory

$$\lambda_c = \frac{\langle k \rangle}{\langle k^2 \rangle - \langle k \rangle}$$

- Zero epidemic threshold for scale-free networks
- Finite epidemic threshold for scale-rich networks

Beyond HMF for SIS

• Wang et al., 2003 Λ_N = largest eigenvalue of adjacency matrix



• Chung et al., 2005

$$\Lambda_N = \begin{cases} c_1 \sqrt{k_c} & \sqrt{k_c} > \frac{\langle k^2 \rangle}{\langle k \rangle} \ln^2(N) \\ c_2 \frac{\langle k^2 \rangle}{\langle k \rangle} & \frac{\langle k^2 \rangle}{\langle k \rangle} > \sqrt{k_c} \ln(N) \end{cases}$$

 k_c = largest degree in the network

Beyond HMF for SIS

• Summing up

$$\lambda_c \simeq \begin{cases} 1/\sqrt{k_c} & \gamma > 5/2\\ \frac{\langle k \rangle}{\langle k^2 \rangle} & 2 < \gamma < 5/2 \end{cases}$$

- In any uncorrelated quenched random network with power-law distributed connectivities, the epidemic threshold goes to zero as the system size goes to infinity.
- This has nothing to do with the scale-free nature of the degree distribution.

SIS $\gamma = 4.5$



Finite Size Scaling SIS $\gamma = 4.5$



SIR $\gamma = 4.5$



Mathematical origin of HMF failure for SIS

 HMF is equivalent to using <u>annealed</u> networks with adjacency matrix

$$a_{ij} = \frac{k_i k_j}{\langle k \rangle N}$$

• This matrix has a unique nonzero eigenvalue

$$\Lambda_N = \frac{\langle k^2 \rangle}{\langle k \rangle}$$

Physical origin of HMF failure for SIS

• Star graph with nodes

 $\rho_{max} \propto (\lambda^2 k_{max} - 1)$ $\rho_1 \propto (\lambda^2 k_{max} - 1)$



• For $\lambda > 1/\sqrt{k_{max}}$ the hub and its neighbors are a self-sustained core of infected nodes, which spread the activity to the rest of the system.

Star vs full graph same k_{max}



Fluctuations of k_{max}



Summary on epidemics

- Zero epidemic threshold for SIS on scalerich networks.
- Finite epidemic threshold for SIR on scalerich networks.
- Conjecture: zero threshold for all models with steady state.
- Caveat: annealed networks are important for real epidemics.

Contact Process (CP)

• Two possible states: susceptible and infected

λ/k

- Two possible events for infected nodes:
- Recovery (rate 1)
- Infection to neighbors (rate λ/k)

Phase-transition with finite threshold

Quenched (disordered) contact process

- A fraction q of nodes has reduced infection rate: λr $0 \le r \le 1$
- A fraction I-q of nodes has normal infection rate: λ

$$\lambda_c(q,r) = \frac{\langle k \rangle}{\langle k \rangle - 1} \ \frac{1}{1 - q(1-r)}.$$

 For q > q_{perc} nodes with normal infection rate form only small clusters

$$q_{perc} = 1 - 1/\langle k \rangle$$







Rare regions effect

• For the "dirty" system, the threshold is larger than the threshold for the pure system.

 $\lambda_c^{dirty} > \lambda_c^{pure}$

- For $\lambda_c^{pure} < \lambda < \lambda_c^{dirty}$, there are rare local clusters of "pure" nodes, which are above the threshold, i.e. in the active phase.
- Activity in pure clusters lives until a coherent fluctuation destroys it. This occurs in a time $\tau(s) \simeq t_0 \exp[A(\lambda)s]$

Rare regions effect

• Size distribution of "pure" clusters

$$P(s) \sim \frac{1}{\sqrt{2\pi p}} s^{-3/2} e^{-s(p-1-\ln(p))} \qquad p = \langle k \rangle (1-q)$$

• Overall activity decay

$$\rho(t) \sim \int ds \ s \ P(s) \exp\left[-t/(t_0 e^{A(\lambda)s})\right] \sim t^{-\gamma(p,\lambda)}$$
$$\gamma(p,\lambda) = -(p-1-\ln(p))/A(\lambda)$$

Generic power-law decay with continuously varying exponents



Conclusions

- Models for epidemics have zero threshold also on scale-rich networks if they have a steady state. HMF may fail.
- Quenched disorder may yield generic power law decays.
- Voter dynamics is strongly affected by scale-free nature. Heterogeneous pair approximation works.

- C. Castellano and R. Pastor-Satorras, "Thresholds for epidemic spreading in networks" (soon in arXiv)
- M.A. Munoz, R. Juhasz, C. Castellano and G. Odor, "Griffiths phases in networks" (soon in arXiv)
- E. Pugliese and C. Castellano, "Heterogeneous pair approximation for voter models on networks", EPL, 88, 58004 (2009)