

A few mathematical tools for biology and medicine

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A few words on the scientific landscape

- <u>D. Barbolosi</u>, A. Iliadis, F. Hubert... (Marseille): tumor growth, chemiotherapy, toxicity
- Th. Colin (Bordeaux) team MONC, start-up Nenuphar: image analysis and prediction of tumor growth
- <u>J. Clairambault</u>, M. Doumic, D. Drasdo & B. Perthame, team MAMBA (Paris): chronotherapy, tumor growth
- GDR MaMoVi, T. Lepoutre
- <u>CEMRACS 2018</u>, Biological and medical applications
- French American Innovation Day : Math & Medicine, Houston, March 2018 with Marc Garbey

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Objectives of the mathematical modeling: why could it be useful ?

- To reproduce the evolution of tumors (with equations and simulations...)
- To anticipate the evolution of tumors
- To anticipate the action of drugs or therapies
- To optimize the action of drugs or therapies

For instance: after resection of a primary tumor, why is a preventative chemiotherapy useful ? Difficulties for imaging techniques to detect micro-metastases : can we predict *in silico* the growth of residual tumors ?

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N(t): number of individuals (cells...) in a population, at time t. Variations due to gain (birth) and loss (death)

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathsf{N}(t)=(\lambda-\mu)\mathsf{N}(t).$$

The solution is $N(t) = e^{(\lambda - \mu)t} N_0$:

exponential growth if $\lambda > \mu$, exponential extinction if $\lambda < \mu$ (cf. Malthus, end of XVIIIth century).

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The discrete viewpoint



We can rewrite $X_n = (1 + (\lambda - \mu)h)^n X_0$; with t = nh, $h \to 0...$ it yields the exponential law.

This is also the approach of **Numerical Schemes** with the "Stability" issue :

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the condition $h \leq (\mu - \lambda)^{-1}$ enforces positivity when $\mu > \lambda$.

(Slightly) more complex models

Verhulst: the more important the population, the stronger the loss rate/the lesser the gain rate

$$\frac{\mathrm{d}}{\mathrm{d}t}X(t) = a\left(1-\frac{X(t)}{K}\right)X(t), \qquad X_{n+1} = X_n + ha\left(1-\frac{X_n}{K}\right)X_n.$$

The population grows when X < K, decays when X > K.

Gompertz: the growth rate itself obeys the exponential law:

$$au(t) = rac{X'(t)}{X(t)} = rac{\mathrm{d}}{\mathrm{d}t} \ln\left(rac{X(t)}{b}
ight)$$

satisfies

$$\frac{\mathrm{d}}{\mathrm{d}t}\tau = -a\tau.$$

It yields $\frac{\mathrm{d}}{\mathrm{d}t}X(t) = aX(t)\ln\left(\frac{b}{X(t)}\right)$ or
 $X_{n+1} = X_n + haX_n\ln(b/X_n).$

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ODE models



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To incorporate the action of drugs



Take into account **toxicity** (pharmaco-kinetic/dynamic)

- c(t) = G(t, u) with u=given dose.
- Admissible set $K = \{u \text{ s. t.} \quad F(t, u) \leq C\}.$

action on healthy cells



keeps the patient alive

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Bolus

$$\frac{\mathrm{d}}{\mathrm{d}t}X(t) = X(t) \big[R(X(t)) - c(t) \big]$$

Question: How can we find $t \mapsto c(t)$ that minimizes X(T) with T=a certain degradation time... under a **constraint** on the dose

$$\int_0^T c(s) \, \mathrm{d} s \leq C_M.$$

• **Bang-bang strategy**: to give C_M at time $T \bullet Gold$

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A multi-dimensional model: Leslie's system, population structured by age

 $x_j = \#$ of individuals with age $j \in \{1, ..., N\}$ $f_j > 0$ fertility rate, $t_{j+1,j}$ transition rate $j \rightarrow j + 1$.



Leslie's system ctn'd

 $\lambda \in \mathbb{C}$ is an **eigenvalue** iff there exists $x \neq 0$ such that $Lx = \lambda x$.

Here *L* has a very specific property: L^n has strictly positive entries (**primitive matrix**).

Perron-Frobenius theorem: $\mu = \max\{|\lambda|, \lambda \text{ eigenvalue of } L\}$ is an eigenvalue, it can be associated to a vector \overline{X} with non negative components.

It governs the asymptotic behavior

$$X^{(k)} \underset{k \to \infty}{\sim} \mu^k C(X^{(0)}) \bar{X}.$$

Similar conclusion for the differential system $\frac{\mathrm{d}}{\mathrm{d}t}X = (L - I)X$: blow up of the population if $\mu > 1$, extinction of $\mu < 1$.

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Large time asymptotics often corresponds to **observable** behaviors.

We can try to exhibit **structure properties** (here L is a primitive matrix) that govern this behavior.

We have efficient **numerical procedures** to compute directly the leading eigenpair, and thus to have direct access to the asymptotic state.

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Towards PDE models: diffusion Diffusion $\partial_t u = \partial_{xx}^2 u$



- the solution is spread
- the solution becomes instantaneously smooth

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Towards PDE models: transport Diffusion $\partial_t u + \partial_x(cu) = \overline{0}$ (here c > 0 constant)



- the solution is... transported (finite speed)
- the solution keeps its regularity (singularities)

Note: numerical schemes solve instead $\partial_t u + \partial_x (cu) = \epsilon \partial^2_{xx} u$, $0 < \epsilon \ll 1$,

K. Iwata, K. Kawasaki, N. Shigesada's model for tumor growth

A Dynamical Model for the Growth and Size Distribution of Multiple Metastatic Tumor, Journal of Theoretical Biology, 2000.

- ▶ Population of tumors structured by their size $x \int_{a}^{b} \rho(t, x) dx = \#$ of cells with size in [a, b] at time t.
- The growth rate depends on the size according to governed by Gompertz' law g(x)
- Each tumor produces new (small) cells with a rate $\beta(x)$.
- McKendrick–Von Foerster type equation:

 $\begin{cases} \partial_t \rho + \partial_x (g(x)\rho) = 0, & \text{for } t \ge 0 \text{ and } 1 < x < b, \\ \rho(0, x) = \delta_{x=1}, \\ g(1)\rho(t, 1) = \int_1^b \beta(x)\rho(t, x) \, \mathrm{d}x + \beta(x_p(t)). \end{cases}$ $g(x) = ax \ln\left(\frac{b}{x}\right), \ \beta(x) = mx^{\alpha}, \ (\text{typiquement } \alpha = 2/3) \end{cases}$

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we have a formula for the solution...

By using Laplace's transform

$$\rho(t,x) = \frac{a}{mb^{\alpha} \ln b} \frac{1}{x} \sum_{k=1}^{\infty} e^{\lambda_k t} \left(1 - \frac{\ln x}{\ln b}\right)^{\lambda_k/a - 1} \frac{1}{c(\lambda_k)},$$

where the λ_k 's are the roots of

$$\frac{a}{m}\lambda_k = F(1, \frac{\lambda_k}{a} + 1; \alpha \ln b) = \sum_{n=0}^{\infty} \frac{1}{(\frac{\lambda_k}{a} + 1) \dots (\frac{\lambda_k}{a} + n)} (\alpha \ln b)^n,$$

and

$$c(\lambda_k) = \sum_{n=0}^{\infty} \frac{(-\alpha \ln b)^n}{n!(\frac{\lambda_k}{a} + n)^2}.$$

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Eigenvalue problem

$$\begin{cases} \frac{\partial}{\partial x}(g(x)N(x)) + \lambda N(x) = 0, \\ g(1)N(1) = \int_{1}^{b} \beta(y)N(y)dy, \\ -g(x)\frac{\partial}{\partial x}\Phi(x) + \lambda\Phi(x) = \Phi(1)\beta(x), \\ \lambda > 0, \quad N(x) \ge 0, \quad \Phi(x) \ge 0, \quad \int_{1}^{b} N\Phi = 1, \quad \int_{1}^{b} N = 1. \end{cases}$$

There exists a unique triple (N, λ_0, Φ) solution of this problem, with N(x) > 0, $\Phi(x) > 0$

An infinite-dimensional version of the Perron-Frabenius theorem...

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Consequences

1. A conservation law:

$$\int_1^b \Phi(x)\rho(x,t)e^{-\lambda_0 t}dx = \int_1^b \Phi(x)\rho_0(x)dx$$

- 2. Estimates in L^{∞} : if $cN(x) \le \rho_0(x) \le CN(x)$ then $cN(x) \le \rho(t, x)e^{-\lambda t} \le CN(x)$
- 3. and in $L^p(\Phi(x)N(x)dx)$.
- 4. The asymptotic behavior is driven by the leading eigenpair

$$\rho(t,x) \underset{t\to\infty}{\sim} C e^{\lambda t} N(x).$$

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Simulation (A. Devys' Phd thesis)



Difficulties:

- the growth rate g vanishes at x = b,
- ▶ b ≫ 1,
- the source at x = 1 is "large".

For $t \ge 2000$ jours $\simeq 5.5$ years the asymptotic profile is a fair approximation of the solution.

Comparison with Tubbiana's experimental data and incoroporation

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of the action of treatments by the research group in Marseille

AABG model for tumor growth

Interacting populations, structuration in size and space

size-structured tumor concentration (t, z) → T(t, z).
 The mass of the tumor changes due to

natural growth + cell division.

▶ a bath of passive cells that become either immune $((t,x) \mapsto E(t,x))$ or "collaborative" $((t,x) \mapsto B(t,x))$.

The model involves **two** distinct length scales; it assumes "scale separation: $z \ll x$ ".

It also means that we neglect some fine scale phenomena...

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AABG model: evolution of tumor cells

$$\partial_t T + \underbrace{\partial_z (VT)}_{\text{"transport"=growth}} = \underbrace{Q(T) - m(E, T)}_{\text{cell division - destruction by immune cells}}$$

Q(T)(t,z) = 4K(2z)T(2z) - K(z)T(z) (binary division) The operator *Q* increases the **number of tumoral cells** $\mu_0(t) = \int_0^\infty f(t,z) \, \mathrm{d}z$, but does not change the **total mass** of the tumor $\mu_1(t) = \int_0^\infty z f(t, z) dz$. If m(E, T) = 0, we get $\frac{\mathrm{d}}{\mathrm{d}t}\mu_0(t) = \int_0^\infty K(z)T(t,z)\,\mathrm{d}z \ge 0,$ $\frac{\mathrm{d}}{\mathrm{d}t}\mu_1(t)=V\mu_0(t)\geq 0.$

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AABG model: motion of the immune cells

- a natural space diffusion,
- a natural death,
- a convection guiding the cells towards the tumor,
- conversion rates from the background passive cells, which are activated by the presence of the tumor.

We are thus led to the following PDE

 $\partial_t E - d\Delta_x E - \nabla_x \cdot (E\nabla_x \Phi_f) = -\gamma_f E + p_f \mu_0 S,$

 $\partial_t B - d\Delta_x B - \nabla_x \cdot (B\nabla_x \Phi_{\text{coll}}) = -\gamma_{\text{coll}} B + p_{\text{coll}} \mu_0 S,$

coupled with a "chemotactic-like" effect

$$\Delta_{\rm x} \Phi_{\rm f} = \mu_0 \sigma_{\rm f}, \qquad \Delta \Phi_{\rm coll} = \mu_0 \sigma_{\rm coll}.$$

AABG model: interaction terms

► Coupling term m(E, T): immune response modeled by Michaelis-Menten kinetics

$$m(E, T)(t) = \int_{\Omega} a(y)E(t, y) \,\mathrm{d}y \times \frac{T}{\alpha + T}$$

 The collaborative cells promote cell divisions: multiply the division operator by

$$1+\int_{\Omega}b(y)\frac{B(t,y)}{1+B(t,y)}\,\mathrm{d}y.$$

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Qualitative features

Numerically, we observe "extinction" or "explosion" depending on the parameters.



An important difficulty (in any application in biology): how should be fixed the parameters of the model ?

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Qualitative features: residual distribution

In fact, one may reach a stable state with residual tumors... again characterized by means of an eigenvalue problem.



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Possible mechanism...

Simplified problem

$$\begin{split} \partial_t T + V \partial_z T &= Q(T) - T \int_{\Omega} a(x) E(t, x) \, \mathrm{d}x, \qquad T(t, 0) = 0, \\ Q(T)(z) &= 4\ell T(2z) - \ell T(z), \\ \partial_t E - d\Delta_x E - \nabla_x \cdot (E\nabla_x \Phi_\mathrm{f}) &= -\gamma_\mathrm{f} E + p_\mathrm{f} \mu_0 S, \\ \Delta_x \Phi_\mathrm{f} &= \mu_0 \sigma_\mathrm{f}. \end{split}$$

It yields¹

$$\frac{\mathrm{d}}{\mathrm{d}t}\mu_0 = \mu_0\left(\boldsymbol{\ell} - \int_{\Omega} \boldsymbol{a}(\boldsymbol{x})\boldsymbol{E}(\boldsymbol{t},\boldsymbol{x})\,\mathrm{d}\boldsymbol{x}\right),\,$$

 and

$$\frac{\mathrm{d}}{\mathrm{d}t}\mu_1 = V\mu_0 - \mu_1 \int_{\Omega} a(x)E(t,x)\,\mathrm{d}x.$$

$$\frac{1}{\mu_0(t)} = \int T(t,z)\,\mathrm{d}z \text{ and } \mu_1(t) = \int z \ T(t,z)\,\mathrm{d}z$$

Possible mechanism...: stationary solutions

With $\partial_t \rightarrow 0$, we get the constraints

$$\boldsymbol{\ell} = \int_{\Omega} \boldsymbol{a}(x) \boldsymbol{E}(t, x) \, \mathrm{d}x, \qquad \boldsymbol{V} \mu_0 = \mu_1 \boldsymbol{\ell}$$

and the stationary equations

$$V \partial_z T = Q(T) - \ell T = 0$$
 eigenvalue pb. !
 $-d\Delta_x E - \mu_0 \nabla_x \cdot (E\nabla_x \Phi_0) = -\gamma_f E + p_f \mu_0 S,$
 $\Delta_x \Phi_0 = \sigma_f.$

It defines μ_0 as a function of ℓ_{\dots}

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Workplan

- Are the numerically observed behaviors qualitatively relevant... or do we miss some important features ?
- Can we understand theoretically these behaviors: parameters thresholds, large time asymptotics... ? What are the hidden structure properties of the PDE system ? Can we relate them to biological interpretation ?
- Analyse and improve the numerical tool (multi-D, parameter investigation, stability issues...).
- Identify the parameters: some are known/observable from experiments, reduce as far as possible the blind parameters. What are the relevant units ? What should be produced by numerics ? Compare to experimental data.
- Incorporate actions of drugs and therapy...

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First
$$X(t) \leq \overline{X}(t)$$
, for $0 \leq t < T$:
Set $U = \ln(X)$, $G(U) = R(X) = R(e^U)$, so that

$$(U-\overline{U})'(t) = G(U(t))-G(\overline{U}(t))-c(t) \leq \int_{\overline{U}(t)}^{U(t)} G'(\sigma) \,\mathrm{d}\sigma \leq C|U-\overline{U}|(t).$$

It implies $([U - \bar{U}]^2_+)'(t) \le 2C[U - \bar{U}]^2_+$.

▶ Second since $R' \leq 0$, and thus $G' \leq 0$, we have

$$(U-ar{U})'(t)\geq -c(t)$$

which yields $U(T) - \overline{U}(T^-) \ge 0 - \int_0^T c(s) \, \mathrm{d}s \ge -C_M$; thus

 $U(T) - \bar{U}(T^+) = U(T) - \bar{U}(T^-) + \bar{U}(T^-) - \bar{U}(T^+) \ge -C_M + C_M = 0.$

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$$X_{1}^{(k+1)} = f_{1}X_{1}^{(k)} + f_{2}X_{2}^{(k)} + \dots + f_{n}X_{n}^{(k)},$$

$$X_{2}^{(k+1)} = t_{2,1} X_{1}^{(k)},$$

$$X_{3}^{(k+1)} = t_{3,2} X_{2}^{(k)},$$

$$\vdots$$

$$X_{n}^{(k+1)} = t_{n,n-1} X_{n-1}^{(k)}$$

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