### Non-fitting meshes for Maxwell's equations

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# **Borehole Logging**





#### Figure: Borehole Logging



### **Objective**

Given a conductivity model, compute a quantity of interest.



Electromagnetic fields  $\mathbf{E}, \mathbf{H}$ 

Applications: Magnetotellurics, Logging-While-Drilling...





FEMs are mesh-based methods.

Fitting meshes are usually used with FEMs.

The physical interfaces are fitted by mesh faces (edges).

The conductivity parameter is constant (smooth) inside each cell.





### Fitting vs non-fitting meshes



Figure: Example of conductivity model

There are two physical interfaces.





## Fitting vs non-fitting meshes



Figure: Example of fitting mesh

The physical interfaces aligned with mesh edges.





# Fitting vs non-fitting meshes



Figure: Example of non-fitting mesh

The physical interfaces are inside mesh cells.





Fitting meshes are good because:

- matrix assembly is easy
- optimal convergence rates

However, they require to:

- use unstructured meshes
- re-generate a mesh for each conductivity model





### Sequence of problems



Electromagnetic fields  $\mathbf{E}^{(j)}, \mathbf{H}^{(j)}$ 































Figure: Position 1: fitting mesh







#### Figure: Position 2: the mesh becomes non-fitting







Figure: Position 2: remeshing required







Figure: Non-fitting meshes avoid remeshing







Figure: Non-fitting meshes avoid remeshing





### **Non-fitting meshes**

The advantages of non-fitting meshes are:

- mesh generation is easier
- tensorial products are possible
- the same mesh can be used in a sequence of problems

The drawbacks of non-fitting meshes are:

- special quadrature schemes need to be used
- convergence rates might be decreased





We focus on first-order Nedelec's edge elements.

Determine when non-fitting meshes can be used.

Analytical study: error estimates for non-fitting meshes.

Numerical study: 2D experiments.





We solve the Maxwell's equation with a constant permeability  $\mu = \mu_0$ .

We assume that  $i\omega\varepsilon + \boldsymbol{\sigma} \simeq \boldsymbol{\sigma}$ , so that we set  $\varepsilon = 0$ .

To simplify, the conductivity  $\sigma$  is scalar and piecewise constant.





### Maxwell's equations: the E-formulation

**E** is solution to

$$i\omega\mu_0\sigma\mathbf{E} + \nabla\times\nabla\times\mathbf{E} = i\omega\mu_0\mathbf{J}.$$

 ${\ensuremath{\mathsf{H}}}$  is obtained by

$$\mathbf{H} = \nabla \times \mathbf{E}.$$





### Maxwell's equations: the H-formulation

 ${\ensuremath{\mathsf{H}}}$  is solution to

$$i\omega\mu_{0}\mathbf{H} + \nabla \times \left(\boldsymbol{\sigma}^{-1}\nabla \times \mathbf{H}\right) = \nabla \times \left(\boldsymbol{\sigma}^{-1}\mathbf{J}\right).$$

**E** is obtained by

$$\mathbf{E} = \boldsymbol{\sigma}^{-1} \left( \mathbf{J} - \nabla \times \mathbf{H} \right).$$





### Standard error-estimates for fitting meshes

In a layered medium, the solution is piecewise smooth.

Hence, if a fitting mesh is used, we have

$$\|\mathbf{E} - \mathbf{E}_{\mathbf{h}}\|_{L^2(\Omega)} = \mathcal{O}(\mathbf{h}),$$

and

$$\|\mathbf{H} - \mathbf{H}_{\mathbf{h}}\|_{L^2(\Omega)} = \mathcal{O}(\mathbf{h}).$$





### Standard error-estimates for non-fitting meshes

With a non-fitting mesh, the solution can jump inside mesh cells.

The solution being less regular, the standard error-estimates give

$$\|\mathbf{E} - \mathbf{E}_{\mathbf{h}}\|_{L^2(\Omega)} = \mathcal{O}(\mathbf{h}^{1/2})$$

and

$$\|\mathbf{H}-\mathbf{H}_{\mathbf{h}}\|_{L^{2}(\Omega)}=\mathcal{O}(\mathbf{h}^{1/2}).$$

At first sight, convergence rates are bad for E and H.

Actually, these error-estimates are pessismistic.





### New error-estimates for non-fitting meshes

Our new result is that actually, we have

$$\|\mathbf{E} - \mathbf{E}_{\mathbf{h}}\|_{L^{2}(\Omega)} = \mathcal{O}(\mathbf{h}^{1/2}),$$

and

$$\|\mathbf{H}-\mathbf{H}_{\mathbf{h}}\|_{L^{2}(\Omega)}=\mathcal{O}(\mathbf{h}).$$

#### $H_h$ converges lineary for fitting and non-fitting meshes.





We have analyzed the general case, where the solution can have singularities.

Our new error-estimate have different convergence rates for  $\mathsf{E}_{\mathsf{h}}$  and  $\mathsf{H}_{\mathsf{h}}.$ 

In general  $H_h$  has a better convergence rate than "expected".

The convergence rate for  $\mathbf{H}_{\mathbf{h}}$  is the same for fitting and non-fitting meshes.







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#### Figure: Layered medium





### Numerical experiments: Layered medium settings



Figure: Layered medium: fitting mesh





### Numerical experiments: Layered medium settings





#### Figure: Layered medium: non-fitting mesh



The current density is polarized as  $\mathbf{J} = (0, \mathbf{J}_y, 0)$ .

The electric field is a scalar  $E_y$ .

The magnetic field is a vector  $\mathbf{H} = (\mathbf{H}_x, \mathbf{H}_z)$ .

H is approximated by edge finite elements.

We compute  $\mathbf{E} = \boldsymbol{\sigma}^{-1} \left( \mathbf{J} - \nabla \times \mathbf{H} \right)$  by post-processing.



# **TE-polarization**



Figure: Numerical errors in TE-polarization





The current density is polarized as  $\mathbf{J} = (0, 0, \mathbf{J}_z)$ .

The electric field is a vector  $\mathbf{E} = (\mathbf{E}_x, \mathbf{E}_z)$ .

The magnetic field is a scalar  $H_{y}$ .

We approximate **E** with edge finite elements.

We obtain  $\mathbf{H}_{y} = \nabla \times \mathbf{E}$  by post-processing.



# **TM-polarization**



Figure: Numerical errors in TM-polarization





The predicted convergence rates are observed numerically.

The convergence rate of  $E_h$  is decreased for non-fitting meshes.

The convergence rate of  $\mathbf{H}_{\mathbf{h}}$  is same for both types of meshes.





For geophysical applications ( $\mu = \mu_0$ ), non-fitting meshes can be used to approximate **H**.

The accuracy loss due to non-fitting meshes is "reasonable".

In our numerical experiments, the error is at most multiplied by 2.





Sharper error estimates for layered media.

More realistic simulations (borehole logging, MT...).

2.5D and 3D Maxwell's equations.

Comparison between the  ${\bf E}$  and  ${\bf H}$  formulations.



