The Multiscale Hybrid Mixed method for the Helmholtz equation Advisors: H. Calandra (TOTAL) H. Barucq (INRIA Magique 3D) C. Gout (LMI INSA ROUEN)

T. Chaumont-Frelet (INRIA Magique 3D, LMI), F. Valentin (LNCC)

September 3, 2013





The Helmholtz Equation

- Wave propagation at a given frequency
- Radar, Acoustics, Seismic (Inverse Problems)
- Heterogeneous media (Seismic)

The Helmholtz Equation

- Wave propagation at a given frequency
- Radar, Acoustics, Seismic (Inverse Problems)
- Heterogeneous media (Seismic)

Numerical Method

- Integral Equations (Homogeneous case)
- Finite Differences
- Finite Elements

The Helmholtz Equation

- Wave propagation at a given frequency
- Radar, Acoustics, Seismic (Inverse Problems)
- Heterogeneous media (Seismic)

Numerical Method

- Integral Equations (Homogeneous case)
- Finite Differences
- Finite Elements

Difficulties

- Pollution effect (High frequency)
- Highly heterogeneous media

1 A simple Helmholtz Problem

2 The MHM method

Oiscretization



A simple Helmholtz Problem

The MHM method Discretization Numerical Experiments

Model Problem

Find $u \in H^1(\Omega)$ such that

$$\begin{cases} -k^2 u - \Delta u = f & \text{in } \Omega \\ \nabla u \cdot \mathbf{n} - iku = 0 & \text{on } \partial \Omega, \end{cases}$$

where k is the wave number, and $f \in L^2(\Omega)$ is the source.

Model Problem

Find $u \in H^1(\Omega)$ such that

$$\begin{cases} -k^2 u - \Delta u = f & \text{in } \Omega \\ \nabla u \cdot \mathbf{n} - iku = 0 & \text{on } \partial \Omega, \end{cases}$$

where k is the wave number, and $f \in L^2(\Omega)$ is the source.

Weak Form

Find $u \in H^1(\Omega)$ such that $a(u, v) = (f, v)_{\Omega}$ for all $v \in H^1(\Omega)$, with $a(u, v) = -k^2(u, v)_{\Omega} - ik(u, v)_{\partial\Omega} + (\nabla u, \nabla v)_{\Omega}.$

Introduction

The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

The MHM method

4/30

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

The MHM method

• C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

- C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).
- Elliptic problems (Elasticity, Reaction-Advection-Diffusion).

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

- C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).
- Elliptic problems (Elasticity, Reaction-Advection-Diffusion).
- Primal Hybrid Formulation.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

- C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).
- Elliptic problems (Elasticity, Reaction-Advection-Diffusion).
- Primal Hybrid Formulation.
- Multiscale Method.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

- C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).
- Elliptic problems (Elasticity, Reaction-Advection-Diffusion).
- Primal Hybrid Formulation.
- Multiscale Method.
- Adapted to highly heterogenous media.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

- C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).
- Elliptic problems (Elasticity, Reaction-Advection-Diffusion).
- Primal Hybrid Formulation.
- Multiscale Method.
- Adapted to highly heterogenous media.
- Interpretation as modified basis functions.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

- C. Harder, D. Paredes and F. Valentin 2012 (Darcy Equation).
- Elliptic problems (Elasticity, Reaction-Advection-Diffusion).
- Primal Hybrid Formulation.
- Multiscale Method.
- Adapted to highly heterogenous media.
- Interpretation as modified basis functions.
- Fine scales are localy captured by the basis functions.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Partition

- Let \mathcal{T}_h be a partition of Ω .
- Let \mathcal{E}_h be the set of edges in \mathcal{T}_h .
- Let \mathcal{E}_h^{int} be the set of internal edges.
- Let \mathcal{E}_h^{ext} be the set of external edges.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Partition

- Let \mathcal{T}_h be a partition of Ω .
- Let \mathcal{E}_h be the set of edges in \mathcal{T}_h .
- Let \mathcal{E}_h^{int} be the set of internal edges.
- Let \mathcal{E}_h^{ext} be the set of external edges.

Functional Spaces

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Partition

- Let \mathcal{T}_h be a partition of Ω .
- Let \mathcal{E}_h be the set of edges in \mathcal{T}_h .
- Let \mathcal{E}_h^{int} be the set of internal edges.
- Let \mathcal{E}_h^{ext} be the set of external edges.

Functional Spaces

We define

$$V = H^1(\mathcal{T}_h) = \left\{ v \in L^2(\Omega) \mid v \in H^1(\mathcal{K}) \quad \forall \mathcal{K} \in \mathcal{T}_h \right\},$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Partition

- Let \mathcal{T}_h be a partition of Ω .
- Let \mathcal{E}_h be the set of edges in \mathcal{T}_h .
- Let \mathcal{E}_h^{int} be the set of internal edges.
- Let \mathcal{E}_h^{ext} be the set of external edges.

Functional Spaces

We define

$$V = H^1(\mathcal{T}_h) = \left\{ v \in L^2(\Omega) \mid v \in H^1(\mathcal{K}) \quad \forall \mathcal{K} \in \mathcal{T}_h \right\},$$

and

$$\Lambda = \left\{ \mu \in \prod_{K \in \mathcal{T}_h} H^{-1/2}(\partial K) \left| \begin{array}{cc} \mu_+|_F + \mu_-|_F = 0 & \forall F \in \mathcal{E}_h^{int} \\ \mu|_F = 0 & \forall F \in \mathcal{E}_h^{ext} \end{array} \right\}.$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem A simple Helmholtz Problem The MHM method Discretization Numerical Experiments Discretization Discreti

Bilinear form

Define $a: V \times V \to \mathbb{C}$ and $b: V \times \Lambda \to \mathbb{C}$ by

$$\begin{aligned} \mathsf{a}(u,v) &= -k^2(u,v)_{\mathcal{T}_h} - ik(u,v)_{\mathcal{E}_h^{\text{ext}}} + (\nabla u, \nabla v)_{\mathcal{T}_h}, \\ \mathsf{b}(u,\mu) &= -\sum_{K\in\mathcal{T}_h} \langle \mu, u \rangle_{\partial K}. \end{aligned}$$

A simple Helmholtz Problem The MHM method Discretization Numerical Experiments Discretization Discretization Numerical Experiments

Bilinear form

Define $a: V \times V \to \mathbb{C}$ and $b: V \times \Lambda \to \mathbb{C}$ by

$$\begin{aligned} \mathsf{a}(u,v) &= -k^2(u,v)_{\mathcal{T}_h} - ik(u,v)_{\mathcal{E}_h^{\text{ext}}} + (\nabla u, \nabla v)_{\mathcal{T}_h}, \\ \mathsf{b}(u,\mu) &= -\sum_{K\in\mathcal{T}_h} \langle \mu, u \rangle_{\partial K}. \end{aligned}$$

Primal Hybrid Formulation

Find $(u, \lambda) \in V \times \Lambda$ such that

$$\begin{cases} a(u,v) + b(v,\lambda) &= (f,v)_{\mathcal{T}_h} \quad \forall v \in V \\ b(u,\mu) &= 0 \qquad \forall \mu \in \Lambda. \end{cases}$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local problems

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local problems

We can rewrite the first equation as

$$\mathsf{a}(u,v) = (f,v)_{\mathcal{T}_h} - \mathsf{b}(v,\lambda) \quad orall v \in V.$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local problems

We can rewrite the first equation as

$$\mathsf{a}(u,v)=(f,v)_{\mathcal{T}_h}-\mathsf{b}(v,\lambda) \quad \forall v\in V.$$

Now, pick $v_{\mathcal{K}} \in V$ such that supp $v_{\mathcal{K}} \subset \mathcal{K}$, then

$$\begin{aligned} a(u, v_{\mathcal{K}}) &= -k^{2}(u, v_{\mathcal{K}})_{\mathcal{K}} - ik(u, v_{\mathcal{K}})_{\partial \mathcal{K} \cap \partial \Omega} + (\nabla u_{\mathcal{K}}, \nabla v_{\mathcal{K}})_{\mathcal{K}} \\ b(v_{\mathcal{K}}, \lambda) &= -\langle \lambda, v_{\mathcal{K}} \rangle_{\partial \mathcal{K} \setminus \partial \Omega} \\ (f, v_{\mathcal{K}})_{\mathcal{T}_{h}} &= (f, v_{\mathcal{K}})_{\mathcal{K}}. \end{aligned}$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local problems

We can rewrite the first equation as

$$\mathsf{a}(u,v)=(f,v)_{\mathcal{T}_h}-\mathsf{b}(v,\lambda) \quad \forall v\in V.$$

Now, pick $v_{\mathcal{K}} \in V$ such that supp $v_{\mathcal{K}} \subset \mathcal{K}$, then

$$\begin{aligned} a(u, v_{K}) &= -k^{2}(u, v_{K})_{K} - ik(u, v_{K})_{\partial K \cap \partial \Omega} + (\nabla u_{K}, \nabla v_{K})_{K} \\ b(v_{K}, \lambda) &= -\langle \lambda, v_{K} \rangle_{\partial K \setminus \partial \Omega} \\ (f, v_{K})_{\mathcal{T}_{h}} &= (f, v_{K})_{K}. \end{aligned}$$

So that $u_K = u|_K$ is solution of

$$\begin{cases} -k^2 u_K - \Delta u_K = f & \text{in } K \\ \nabla u_K \cdot \mathbf{n}_K = \lambda & \text{on } \partial K \setminus \partial \Omega \\ \nabla u_K \cdot \mathbf{n} - iku_K = 0 & \text{on } \partial K \cap \partial \Omega. \end{cases}$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Well posedness of the local problems

We only need check for unicity

$$\begin{cases}
-k^2 u - \Delta u = 0 & \text{in } K \\
\nabla u \cdot \mathbf{n} = 0 & \text{on } \partial K \setminus \partial \Omega \\
u = 0 & \text{on } \partial K \cap \partial \Omega.
\end{cases}$$

The problem is well-posed when k^2 is not an eigenvalue of $-\Delta$.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Well posedness of the local problems

We only need check for unicity

$$\begin{cases}
-k^2 u - \Delta u = 0 & \text{in } K \\
\nabla u \cdot \mathbf{n} = 0 & \text{on } \partial K \setminus \partial \Omega \\
u = 0 & \text{on } \partial K \cap \partial \Omega.
\end{cases}$$

The problem is well-posed when k^2 is not an eigenvalue of $-\Delta$.

Example

For an interior square $K = (0, h) \times (0, h) \subset \mathbb{R}^2$.

$$\lambda_0 = 0, \quad \lambda_1 = h^2 \pi^2.$$

If $kh < \pi$, then $\lambda_0 < k^2 < \lambda_1$, and we have well posedness.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local operators

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local operators

We define $T_{\mathcal{K}}: H^{-1/2}(\partial \mathcal{K}) \to H^1(\mathcal{K})$ for all $\mu \in H^{-1/2}(\partial \mathcal{K})$ with

$$a(T_{\mathcal{K}}\mu, v_{\mathcal{K}}) = -\langle \mu, v_{\mathcal{K}} \rangle_{\partial \mathcal{K}} \quad \forall v_{\mathcal{K}} \in H^1(\mathcal{K}).$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local operators

We define $T_{\kappa} : H^{-1/2}(\partial K) \to H^{1}(K)$ for all $\mu \in H^{-1/2}(\partial K)$ with $a(T_{\kappa}\mu, v_{\kappa}) = -\langle \mu, v_{\kappa} \rangle_{\partial K} \quad \forall v_{\kappa} \in H^{1}(K).$ We define $\hat{T}_{\kappa} : L^{2}(K) \to H^{1}(K)$ for all $g \in L^{2}(K)$ with $a(\hat{T}_{\kappa}g, v_{\kappa}) = (g, v_{\kappa})_{\kappa} \quad \forall v_{\kappa} \in H^{1}(K).$
Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Local operators

We define
$$T_{\mathcal{K}}: H^{-1/2}(\partial \mathcal{K}) \to H^1(\mathcal{K})$$
 for all $\mu \in H^{-1/2}(\partial \mathcal{K})$ with

$$\mathsf{a}(\mathsf{T}_{\mathsf{K}}\mu,\mathsf{v}_{\mathsf{K}})=-\langle\mu,\mathsf{v}_{\mathsf{K}}\rangle_{\partial\mathsf{K}}\quad\forall\mathsf{v}_{\mathsf{K}}\in\mathsf{H}^{1}(\mathsf{K}).$$

We define $\hat{T}_{K}: L^{2}(K) \rightarrow H^{1}(K)$ for all $g \in L^{2}(K)$ with

$$a(\hat{T}_{K}g,v_{K})=(g,v_{K})_{K}\quad\forall v_{K}\in H^{1}(K).$$

Local expression of u_K

$$u_K = T_K \lambda + \hat{T}_K f$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Gobal expression of u

Grouping up the pieces we define $T : \Lambda \to V$ and $\hat{T} : L^2(\Omega) \to V$. Then

$$u=T\lambda+\hat{T}f.$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Gobal expression of u

Grouping up the pieces we define $T : \Lambda \to V$ and $\hat{T} : L^2(\Omega) \to V$. Then

$$u=T\lambda+\hat{T}f.$$

Global Problem

To obtain the global problem, we simply substitute $u = T\lambda + \hat{T}f$ in the the second equation $(b(u, \mu) = 0)$.

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem **The MHM method** Interpretation of the global problem

The MHM method

Introduction The Primal Hybrid Formulation Local Problems Global Problem **The MHM method** Interpretation of the global problem

The MHM method

• First, localy define the operator T and \hat{T} such that

$$egin{array}{rcl} m{a}(T\mu,m{v}) &=& -b(\mu,m{v}) \ m{a}(\hat{T}g,m{v}) &=& (g,m{v})_\Omega \end{array}$$

for all $\mu \in \Lambda$, $g \in L^2(\Omega)$ and $v \in V$.

Introduction The Primal Hybrid Formulation Local Problems Global Problem **The MHM method** Interpretation of the global problem

The MHM method

• First, localy define the operator T and \hat{T} such that

$$egin{array}{rcl} a(T\mu,v)&=&-b(\mu,v)\ a(\hat{T}g,v)&=&(g,v)_{\Omega} \end{array}$$

for all $\mu \in \Lambda$, $g \in L^2(\Omega)$ and $v \in V$.

• Find $\lambda \in \Lambda$ solution of

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda.$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem **The MHM method** Interpretation of the global problem

The MHM method

• First, localy define the operator ${\cal T}$ and $\hat{{\cal T}}$ such that

$$egin{array}{rcl} m{a}(T\mu,m{v}) &=& -b(\mu,m{v})\ m{a}(\hat{T}g,m{v}) &=& (g,m{v})_\Omega \end{array}$$

for all $\mu \in \Lambda$, $g \in L^2(\Omega)$ and $v \in V$.

• Find $\lambda \in \Lambda$ solution of

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda.$$

• Then $u = T\lambda + \hat{T}f$.

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Interpretation of the global problem

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Interpretation of the global problem

• Recall that the global problem reads as

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda.$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Interpretation of the global problem

• Recall that the global problem reads as

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda.$$

• By definition of T with $v = T\lambda$

$$a(T\lambda, T\mu) = a(T\mu, T\lambda) = -b(T\lambda, \mu).$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Interpretation of the global problem

• Recall that the global problem reads as

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda.$$

• By definition of T with $v = T\lambda$

$$a(T\lambda, T\mu) = a(T\mu, T\lambda) = -b(T\lambda, \mu).$$

• By definition of T with $v = \hat{T}f$ and \hat{T} with $v = T\mu$

$$-b(\hat{T}f,\mu) = a(T\mu,\hat{T}f) = a(\hat{T}f,T\mu) = (f,T\mu)_{\mathcal{T}_h}.$$

Introduction The Primal Hybrid Formulation Local Problems Global Problem The MHM method Interpretation of the global problem

Interpretation of the global problem

• Recall that the global problem reads as

$$b(T\lambda,\mu) = -b(\hat{T}f,\mu) \quad \forall \mu \in \Lambda.$$

• By definition of T with $v = T\lambda$

$$a(T\lambda, T\mu) = a(T\mu, T\lambda) = -b(T\lambda, \mu).$$

• By definition of T with $v = \hat{T}f$ and \hat{T} with $v = T\mu$

$$-b(\hat{T}f,\mu) = a(T\mu,\hat{T}f) = a(\hat{T}f,T\mu) = (f,T\mu)_{\mathcal{T}_h}.$$

• We can rewrite the global problem as

$$a(T\lambda, T\mu) = (f, T\mu)_{\mathcal{T}_h} \quad \forall \mu \in \Lambda.$$

Introduction

Notations Local problems Global Problem

Introduction Notations Local problems Global Problem

Introduction

• The operators T_K and \hat{T}_K are computed with a second level method (e.g. FEM).

Introduction Notations Local problems Global Problem

- The operators T_K and \hat{T}_K are computed with a second level method (e.g. FEM).
- We only need to discretise $\Lambda_h \subset \Lambda$.

Introduction Notations Local problems Global Problem

- The operators T_K and \hat{T}_K are computed with a second level method (e.g. FEM).
- We only need to discretise $\Lambda_h \subset \Lambda$.
- The quality of the solution only depends upon Λ_h .

Introduction Notations Local problems Global Problem

- The operators T_K and \hat{T}_K are computed with a second level method (e.g. FEM).
- We only need to discretise $\Lambda_h \subset \Lambda$.
- The quality of the solution only depends upon Λ_h .
- We have to approximate $\lambda = \nabla u \cdot \mathbf{n}$.

Introduction Notations Local problems Global Problem

- The operators T_K and \hat{T}_K are computed with a second level method (e.g. FEM).
- We only need to discretise $\Lambda_h \subset \Lambda$.
- The quality of the solution only depends upon Λ_h .
- We have to approximate $\lambda = \nabla u \cdot \mathbf{n}$.
- In 2d, we need basis functions living on one dimensional edges.

Introduction Notations Local problems Global Problem

Some notations

- We note $(\varphi_j)_j$ a global basis of Λ_h .
- In each cell K, we note $(\psi_m^K)_m$ a local basis in each K.
- If we consider a φ_j , it lives on an edge $F = \partial K_+ \cap \partial K_-$.
- There are corresponding functions $\psi_{m_+}^{K_+}$ and $\psi_{m_-}^{K_-}$.

Introduction Notations Local problems Global Problem

Local Problems

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}\rangle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$a(\hat{f},v_{K})=(f,v_{K})_{K},\quad \forall v_{K}\in H^{1}(K).$$

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}
angle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$\mathsf{a}(\hat{f},\mathsf{v}_{\mathcal{K}})=(f,\mathsf{v}_{\mathcal{K}})_{\mathcal{K}},\quad \forall \mathsf{v}_{\mathcal{K}}\in H^1(\mathcal{K}).$$

• Then its clear that $\eta_m^K = T_K \psi_m^K$ and $\hat{f}^K = \hat{T}_K f$.

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}
angle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$a(\hat{f},v_{\mathcal{K}})=(f,v_{\mathcal{K}})_{\mathcal{K}},\quad \forall v_{\mathcal{K}}\in H^1(\mathcal{K}).$$

Then its clear that η^K_m = T_Kψ^K_m and f^K = T^ˆ_Kf.
By regrouping, we have η_j = Tφ_j and f̂ = T̂f.

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}
angle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$\mathsf{a}(\widehat{f},\mathsf{v}_{\mathcal{K}})=(f,\mathsf{v}_{\mathcal{K}})_{\mathcal{K}},\quad \forall \mathsf{v}_{\mathcal{K}}\in H^1(\mathcal{K}).$$

Then its clear that η^K_m = T_Kψ^K_m and f^K = T^ˆ_Kf.
By regrouping, we have η_j = Tφ_j and f̂ = T̂f.

Storage

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}
angle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$\mathsf{a}(\widehat{f},\mathsf{v}_{\mathcal{K}})=(f,\mathsf{v}_{\mathcal{K}})_{\mathcal{K}},\quad orall \mathsf{v}_{\mathcal{K}}\in H^1(\mathcal{K}).$$

Then its clear that η^K_m = T_Kψ^K_m and f^K = T[˜]_Kf.
By regrouping, we have η_j = Tφ_j and f = T[˜]f.

Storage

• The η_i will be used as basis functions in the global problem.

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}
angle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$\mathsf{a}(\widehat{f},\mathsf{v}_{\mathcal{K}})=(f,\mathsf{v}_{\mathcal{K}})_{\mathcal{K}},\quad orall \mathsf{v}_{\mathcal{K}}\in H^1(\mathcal{K}).$$

Then its clear that η^K_m = T_Kψ^K_m and f^K = T[˜]_Kf.
By regrouping, we have η_j = Tφ_j and f = T[˜]f.

Storage

- The η_i will be used as basis functions in the global problem.
- \hat{f} will be used in the second member of the global problem.

Introduction Notations Local problems Global Problem

Local Problems

• Pick a cell $K \in \mathcal{T}_h$. For all ψ_m^K , we solve

$$\mathsf{P}(\eta_m^{\mathsf{K}},\mathsf{v}_{\mathsf{K}}) = -\langle \mathsf{v}_{\mathsf{K}},\psi_m^{\mathsf{K}}
angle_{\partial \mathsf{K}} \quad \forall \mathsf{v}_{\mathsf{K}} \in H^1(\mathsf{K}),$$

as well as

$$\mathsf{a}(\widehat{f},\mathsf{v}_{\mathcal{K}})=(f,\mathsf{v}_{\mathcal{K}})_{\mathcal{K}},\quad orall \mathsf{v}_{\mathcal{K}}\in H^1(\mathcal{K}).$$

Then its clear that η^K_m = T_Kψ^K_m and f^K = T_Kf.
By regrouping, we have η_i = Tφ_i and f̂ = Tf.

Storage

- The η_j will be used as basis functions in the global problem.
- \hat{f} will be used in the second member of the global problem.
- We only need to store their value on the edge of the mesh.

Introduction Notations Local problems Global Problem

Global Problem

Introduction Notations Local problems Global Problem

Global Problem

• We seek the discrete solution $\lambda_h \in \Lambda_h$ as $\lambda_h = \sum_j \alpha_j \varphi_j$.

Introduction Notations Local problems Global Problem

Global Problem

- We seek the discrete solution $\lambda_h \in \Lambda_h$ as $\lambda_h = \sum_j \alpha_j \varphi_j$.
- Then the global problem reads

$$\sum_{j} b(\varphi_i, T\varphi_j) \alpha_j = -b(\hat{T}f, \varphi_i) \quad \forall i.$$

Introduction Notations Local problems Global Problem

Global Problem

- We seek the discrete solution $\lambda_h \in \Lambda_h$ as $\lambda_h = \sum_j \alpha_j \varphi_j$.
- Then the global problem reads

$$\sum_{j} b(\varphi_i, T\varphi_j) \alpha_j = -b(\hat{T}f, \varphi_i) \quad \forall i.$$

• Using the previous stage we can transform it into

$$\sum_{j} b(\varphi_i, \eta_j) \alpha_j = -b(\hat{f}, \varphi_i) \quad \forall i.$$

Introduction Notations Local problems Global Problem

Global Problem

- We seek the discrete solution $\lambda_h \in \Lambda_h$ as $\lambda_h = \sum_j \alpha_j \varphi_j$.
- Then the global problem reads

$$\sum_{j} b(\varphi_i, T\varphi_j) \alpha_j = -b(\hat{T}f, \varphi_i) \quad \forall i.$$

• Using the previous stage we can transform it into

$$\sum_{j} b(\varphi_i, \eta_j) \alpha_j = -b(\hat{f}, \varphi_i) \quad \forall i.$$

• So that we can easily assemble the matrix and solve the linear system. Then we have

$$u_h = \sum_j \alpha_j \eta_j + \hat{f}.$$

Introduction Notations Local problems Global Problem

Interpretation of the global problem

Introduction Notations Local problems Global Problem

Interpretation of the global problem

• The global problem may be seen as

 $a(T\lambda, T\mu) = (f, T\mu)_{\mathcal{T}_h}.$

Introduction Notations Local problems Global Problem

Interpretation of the global problem

• The global problem may be seen as

 $a(T\lambda, T\mu) = (f, T\mu)_{\mathcal{T}_h}.$

• In this case, the discrete counterpart is

$$\sum_{j=1}^{N} a(\eta_i, \eta_j) \alpha_j = (f, \eta_i)_{\mathcal{T}_h} \quad \forall i.$$
Introduction Notations Local problems Global Problem

Interpretation of the global problem

• The global problem may be seen as

 $a(T\lambda, T\mu) = (f, T\mu)_{\mathcal{T}_h}.$

• In this case, the discrete counterpart is

$$\sum_{j=1}^{N} a(\eta_i, \eta_j) \alpha_j = (f, \eta_i)_{\mathcal{T}_h} \quad \forall i.$$

• It looks like we have solve the classical FEM formulation, with special basis functions η_j .

Introduction Notations Local problems Global Problem

Interpretation of the global problem

• The global problem may be seen as

 $a(T\lambda, T\mu) = (f, T\mu)_{\mathcal{T}_h}.$

• In this case, the discrete counterpart is

$$\sum_{j=1}^{N} a(\eta_i, \eta_j) \alpha_j = (f, \eta_i)_{\mathcal{T}_h} \quad \forall i.$$

- It looks like we have solve the classical FEM formulation, with special basis functions η_j.
- The basis functions η_j being computed localy as solution of local subproblems.

A simple test Anisotropy study Comparison with the classical FEM

A simple test

Let $\Omega = (0,1) \times (0,1)$ and y = (0.5,1.1). We consider the problem

$$\begin{cases} -k^2 u - \Delta u = 0 & \text{in } \Omega \\ \nabla u \cdot \mathbf{n} - iku = \nabla \mathcal{H} \cdot \mathbf{n} - ik\mathcal{H} & \text{on } \partial \Omega, \end{cases}$$

where

$$\mathcal{H}(x) = Y(k|x-y|) + iJ(k|x-y|),$$

for all $x \in \Omega$. We set $k = 2\pi f$, with f = 12.1.

A simple test Anisotropy study Comparison with the classical FEM

Real and Imaginary part of the solution



A simple test Anisotropy study Comparison with the classical FEM

Convergence curve I = 2



A simple test Anisotropy study Comparison with the classical FEM

Convergence curve I = 3



A simple test Anisotropy study Comparison with the classical FEM

Convergence curve I = 4



A simple test Anisotropy study Comparison with the classical FEM

Convergence curve l = 5



A simple test Anisotropy study Comparison with the classical FEM

Anisotropy study

We still consider $\Omega = (0, 1) \times (0, 1)$ and f = 12.1. For every angle $\theta \in [0, \pi/4]$, we consider the plane wave

$$\mathbf{e}_{\theta}(x)=e^{ik\nu\cdot x},$$

for $x \in \Omega$ with $\nu = (\cos \theta, \sin \theta)$. We solve the following problem

$$\begin{cases} -k^2 u - \Delta u = 0 & \text{in } \Omega \\ \nabla u \cdot \mathbf{n} - iku = \nabla \mathbf{e}_{\theta} \cdot \mathbf{n} - ik\mathbf{e}_{\theta} & \text{on } \partial\Omega. \end{cases}$$

A simple test Anisotropy study Comparison with the classical FEM

n = 15, l = 2



A simple test Anisotropy study Comparison with the classical FEM

n = 10, l = 3



A simple test Anisotropy study Comparison with the classical FEM

n = 10, l = 4



A simple test Anisotropy study Comparison with the classical FEM

n = 10, l = 5



A simple test Anisotropy study Comparison with the classical FEM

Exactness of the scheme for $\theta = j\pi/2$

A simple test Anisotropy study Comparison with the classical FEM

Exactness of the scheme for $\theta = j\pi/2$

• Consider an edge F. We can parametrize it as

$$F = \{x_0 + t\mathbf{v}_F \mid t \in (0, h)\}.$$

A simple test Anisotropy study Comparison with the classical FEM

Exactness of the scheme for $\theta = j\pi/2$

• Consider an edge F. We can parametrize it as

$$F = \{x_0 + t\mathbf{v}_F \mid t \in (0, h)\}.$$

$$(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_{F})(x) = ik\mathbf{n}_{F} \cdot \nu e^{ik\nu \cdot x}$$

A simple test Anisotropy study Comparison with the classical FEM

Exactness of the scheme for $\theta = j\pi/2$

• Consider an edge F. We can parametrize it as

$$F = \{x_0 + t\mathbf{v}_F \mid t \in (0, h)\}.$$

$$(\nabla \mathbf{e}_{ heta} \cdot \mathbf{n}_F)(x) = ik\mathbf{n}_F \cdot
u e^{ik
u \cdot x}$$

• Taking
$$x = x(t) = x_0 + t\mathbf{v}_F$$
, we get
 $(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_F)|_F(t) = ike^{ik\nu \cdot x_0}\mathbf{n}_F \cdot \nu e^{ik\nu \cdot \mathbf{v}_F t}$

A simple test Anisotropy study Comparison with the classical FEM

Exactness of the scheme for $\theta = j\pi/2$

• Consider an edge F. We can parametrize it as

$$F = \{x_0 + t \mathbf{v}_F \mid t \in (0, h)\}.$$

$$(\nabla \mathbf{e}_{ heta} \cdot \mathbf{n}_F)(x) = ik\mathbf{n}_F \cdot
u e^{ik
u \cdot x}$$

• Taking
$$x = x(t) = x_0 + t\mathbf{v}_F$$
, we get
 $(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_F)|_F(t) = ike^{ik\nu \cdot x_0}\mathbf{n}_F \cdot \nu e^{ik\nu \cdot \mathbf{v}_F t}$

• If
$$\mathbf{n}_F \cdot \nu = 0$$
, then $(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_F)|_F = 0 \in \mathbb{C}$

A simple test Anisotropy study Comparison with the classical FEM

Exactness of the scheme for $\theta = j\pi/2$

• Consider an edge F. We can parametrize it as

$$F = \{x_0 + t \mathbf{v}_F \mid t \in (0, h)\}.$$

$$(\nabla \mathbf{e}_{ heta} \cdot \mathbf{n}_F)(x) = ik\mathbf{n}_F \cdot \nu e^{ik\nu \cdot x}$$

• Taking
$$x = x(t) = x_0 + t\mathbf{v}_F$$
, we get
 $(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_F)|_F(t) = ike^{ik\nu \cdot x_0}\mathbf{n}_F \cdot \nu e^{ik\nu \cdot \mathbf{v}_F t}$

• If
$$\mathbf{n}_F \cdot \nu = 0$$
, then $(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_F)|_F = 0 \in \mathbb{C}$
• If $\mathbf{v}_F \cdot \nu = 0$, then $(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_F)|_F = ike^{ik\nu \cdot x_0} = C \in \mathbb{C}$.

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• We can make the scheme exact for plane waves with angle

$$\theta = \pi/4, \ 3\pi/4, \ -\pi/4, \ -3\pi/4.$$

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• We can make the scheme exact for plane waves with angle

$$\theta = \pi/4, \ 3\pi/4, \ -\pi/4, \ -3\pi/4.$$

• To do that, remember that

$$(\nabla \mathbf{e}_{\theta} \cdot \mathbf{n}_{F})|_{F}(t) = ike^{ik\nu\cdot\mathbf{x}_{0}}\mathbf{n}_{F}\cdot\nu e^{ik\nu\cdot\mathbf{v}_{F}t} = Ce^{ik\nu\cdot\mathbf{v}_{F}t}$$

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• We can make the scheme exact for plane waves with angle

$$\theta = \pi/4, \ 3\pi/4, \ -\pi/4, \ -3\pi/4.$$

To do that, remember that

$$(\nabla \mathbf{e}_{ heta} \cdot \mathbf{n}_F)|_F(t) = ike^{ik\nu \cdot \mathbf{x}_0}\mathbf{n}_F \cdot \nu e^{ik\nu \cdot \mathbf{v}_F t} = Ce^{ik\nu \cdot \mathbf{v}_F t}$$

• On a cartesian mesh, we have either u = (1,0) or (0,1) and

$$\nu \cdot \mathbf{v}_F = \pm \frac{\sqrt{2}}{2}.$$

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• In each edge, consider the following basis functions

$$arphi_0(t)=1, \quad arphi_+(t)=e^{ik\sqrt{2}/2t}, \quad arphi_-(t)=e^{-ik\sqrt{2}/2t}.$$

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• In each edge, consider the following basis functions

$$arphi_0(t)=1, \quad arphi_+(t)=e^{ik\sqrt{2}/2t}, \quad arphi_-(t)=e^{-ik\sqrt{2}/2t}.$$

• Then the corresponding scheme is exact for plane waves with angle $\theta = j\pi/4$.

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• In each edge, consider the following basis functions

$$arphi_0(t)=1, \quad arphi_+(t)=e^{ik\sqrt{2}/2t}, \quad arphi_-(t)=e^{-ik\sqrt{2}/2t}.$$

- Then the corresponding scheme is exact for plane waves with angle $\theta = j\pi/4$.
- We can generalize make the scheme exact for plane waves with angles $\theta = j\pi/2n$,

A simple test Anisotropy study Comparison with the classical FEM

Plane Wave-Adapted Elements

• In each edge, consider the following basis functions

$$arphi_0(t) = 1, \quad arphi_+(t) = e^{ik\sqrt{2}/2t}, \quad arphi_-(t) = e^{-ik\sqrt{2}/2t}.$$

- Then the corresponding scheme is exact for plane waves with angle $\theta = j\pi/4$.
- We can generalize make the scheme exact for plane waves with angles $\theta = j\pi/2n$,

$$\phi_0(t)=1, \quad \phi^m_+(t)=e^{iklpha_m t}, \quad \phi^m_-(t)=e^{iklpha_m t},$$

with $\alpha_m = \cos(m\pi/2n)$, and m = 1, n - 1.

A simple test Anisotropy study Comparison with the classical FEM

I=15, Λ_2 vs Λ_0^1



T. Chaumont-Frelet, F. Valentin The MHM method for the Helmholtz equation

A simple test Anisotropy study Comparison with the classical FEM

I=10, Λ_3 vs Λ_1^1



A simple test Anisotropy study Comparison with the classical FEM

I=10, Λ_4 vs Λ_0^2



A simple test Anisotropy study Comparison with the classical FEM

I=10, Λ_5 vs Λ_1^2



A simple test Anisotropy study Comparison with the classical FEM

A more complex test

We solve

$$\begin{cases} -k^2 u - \Delta u = 0 & \text{in } \Omega \\ \nabla u \cdot \mathbf{n} - iku = g & \text{on } \partial\Omega \setminus \Gamma \\ \nabla u \cdot \mathbf{n} = 0 & \text{on } \Gamma, \end{cases}$$

where g is composed of 3 ponctual sources with different locations and amplitudes.

A simple test Anisotropy study Comparison with the classical FEM

Real and Imaginary part of the solution



A simple test Anisotropy study Comparison with the classical FEM

\mathcal{P}_3 , Λ_2 vs Λ_0^1



A simple test Anisotropy study Comparison with the classical FEM




A simple test Anisotropy study Comparison with the classical FEM

\mathcal{P}_4 , Λ_3 vs Λ_1^1



A simple test Anisotropy study Comparison with the classical FEM





A simple test Anisotropy study Comparison with the classical FEM

Conclusion and perspectives

- The MHM methodology works on the Helmholtz equation.
- The results are good compared to the classical FEM.
- Up to 7 times less degrees of freedom.

A simple test Anisotropy study Comparison with the classical FEM

Conclusion and perspectives

- The MHM methodology works on the Helmholtz equation.
- The results are good compared to the classical FEM.
- Up to 7 times less degrees of freedom.
- We need to investigate resonance in subproblems.
- We could also try to derive conditions on the mesh.
- There is a solution in the MHM framework.

A simple test Anisotropy study Comparison with the classical FEM

Conclusion and perspectives

- The MHM methodology works on the Helmholtz equation.
- The results are good compared to the classical FEM.
- Up to 7 times less degrees of freedom.
- We need to investigate resonance in subproblems.
- We could also try to derive conditions on the mesh.
- There is a solution in the MHM framework.
- We want to tackle (highly) heterogeneous problems.
- Condition on the mesh to avoid resonance?
- Choice of basis functions?